

Overview and frontiers in the hunt for Supersymmetry at the LHC



Till Eifert (CERN),
Colloquium, DESY, May 2018



Outline

- **Motivation**

Why are we looking for physics beyond the standard model?
What is supersymmetry and why is it interesting?

- **Searching for Supersymmetry at the LHC**

Brief introduction to the experimental setup.
Description of a typical supersymmetry analysis.

- **Overview of search program and results**

How robust are the exclusions?

- **New frontiers**

complex final states and searches for rare & challenging signatures.

- **Outlook**

Next steps in the hunt for supersymmetry.

THE STANDARD MODEL OF FUNDAMENTAL PARTICLES AND INTERACTIONS

The Standard Model is a quantum theory that summarizes our current knowledge of the physics of fundamental particles and fundamental interactions (interactions are manifested by forces and by decay rates of unstable particles).

FERMIONS

matter constituents
spin = 1/2, 3/2, 5/2, ...

| Leptons spin = 1/2 | | | Quarks spin = 1/2 | | |
|----------------------------|----------------------------|-----------------|-------------------|---------------------------------|-----------------|
| Flavor | Mass GeV/c ² | Electric charge | Flavor | Approx. Mass GeV/c ² | Electric charge |
| ν_L lightest neutrino* | $(0-2) \times 10^{-9}$ | 0 | u up | 0.002 | 2/3 |
| e electron | 0.000511 | -1 | d down | 0.005 | -1/3 |
| ν_M middle neutrino* | $(0.009-2) \times 10^{-9}$ | 0 | c charm | 1.3 | 2/3 |
| μ muon | 0.106 | -1 | s strange | 0.1 | -1/3 |
| ν_H heaviest neutrino* | $(0.05-2) \times 10^{-9}$ | 0 | t top | 173 | 2/3 |
| τ tau | 1.777 | -1 | b bottom | 4.2 | -1/3 |

*See the neutrino paragraph below.

Spin is the intrinsic angular momentum of particles. Spin is given in units of \hbar , which is the quantum unit of angular momentum where $\hbar = h/2\pi = 6.58 \times 10^{-25} \text{ GeV s} = 1.05 \times 10^{-34} \text{ J s}$.

Electric charges are given in units of the proton's charge. In SI units the electric charge of the proton is 1.60×10^{-19} coulombs.

The energy unit of particle physics is the electronvolt (eV), the energy gained by one electron in crossing a potential difference of one volt. Masses are given in GeV/c² (remember $E = mc^2$) where $1 \text{ GeV} = 10^9 \text{ eV} = 1.60 \times 10^{-10} \text{ joule}$. The mass of the proton is 0.938 GeV/c² = $1.67 \times 10^{-27} \text{ kg}$.

Neutrinos

Neutrinos are produced in the sun, supernovae, reactors, accelerator collisions, and many other processes. Any produced neutrino can be described as one of three neutrino flavor states ν_e , ν_μ , or ν_τ , labelled by the type of charged lepton associated with its production. Each is a defined quantum mixture of the three definite-mass neutrinos ν_L , ν_M , and ν_H for which currently allowed mass ranges are shown in the table. Further exploration of the properties of neutrinos may yield powerful clues to puzzles about matter and antimatter and the evolution of stars and galaxy structures.

Matter and Antimatter

For every particle type there is a corresponding antiparticle type, denoted by a bar over the particle symbol (unless + or - charge is shown). Particle and antiparticle have identical mass and spin but opposite charges. Some electrically neutral bosons (e.g., Z^0 , γ , and $\eta_c = c\bar{c}$ but not $K^0 = d\bar{s}$) are their own antiparticles.

Particle Processes

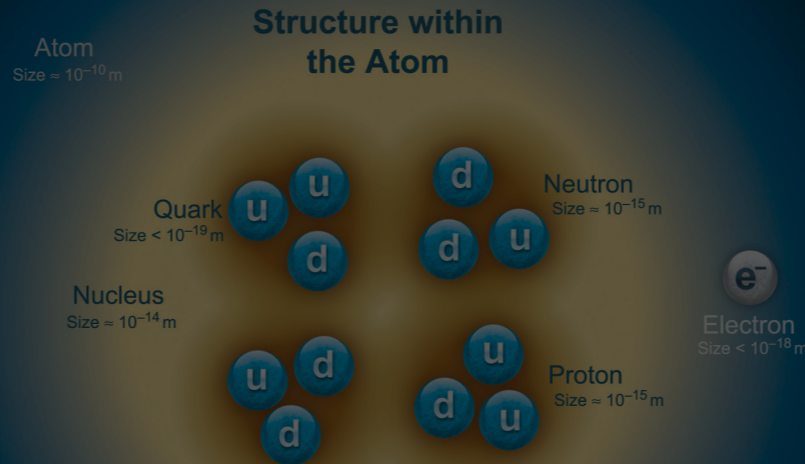
These diagrams are an artist's conception. Orange shaded areas represent the cloud of gluons.

$n \rightarrow p e^- \bar{\nu}_e$

A free neutron (udd) decays to a proton (uud), an electron, and an antineutrino via a virtual (mediating) W boson. This is neutron β (beta) decay.

$e^+ e^- \rightarrow B^0 \bar{B}^0$

An electron and positron (antielectron) colliding at high energy can annihilate to produce B^0 and \bar{B}^0 mesons via a virtual Z boson or a virtual photon.



If the proton and neutrons in this picture were

10 cm across, then the quarks and electrons

would be less than 0.1 mm in size and the

quarks would be about 10 km apart.

Motivation

Properties of interactions

The strengths of the interactions (forces) are shown relative to the strength of the electromagnetic force for two u quarks separated by the specified distances.

| Property | Gravitational Interaction | Weak Interaction (Electroweak) | Electromagnetic Interaction | Strong Interaction |
|-------------------------------|-----------------------------|--------------------------------|-----------------------------|--------------------|
| Acts on: | Mass – Energy | Flavor | Electric Charge | Color Charge |
| Particles experiencing: | All | Quarks, Leptons | Electrically Charged | Quarks, Gluons |
| Particles mediating: | Graviton (not yet observed) | W^+ W^- Z^0 | γ | Gluons |
| Strength at | | | | |
| 10^{-18} m | 10^{-41} | 0.8 | 1 | 25 |
| $3 \times 10^{-17} \text{ m}$ | 10^{-41} | 10^{-4} | 1 | 60 |

BOSONS

force carriers
spin = 0, 1, 2, ...

| Unified Electroweak spin = 1 | | |
|------------------------------|-------------------------|-----------------|
| Name | Mass GeV/c ² | Electric charge |
| γ photon | 0 | 0 |
| W^- | 80.39 | -1 |
| W^+ | 80.39 | +1 |
| W bosons | | |
| Z^0 | 91.188 | 0 |
| Z boson | | |

| Strong (color) spin = 1 | | |
|-------------------------|-------------------------|-----------------|
| Name | Mass GeV/c ² | Electric charge |
| g gluon | 0 | 0 |

| Higgs Boson spin = 0 | | |
|----------------------|-------------------------|-----------------|
| Name | Mass GeV/c ² | Electric charge |
| H Higgs | 126 | 0 |

Higgs Boson

The Higgs boson is a critical component of the Standard Model. Its discovery helps confirm the mechanism by which fundamental particles get mass.

Color Charge

Only quarks and gluons carry "strong charge" (also called "color charge") and can have strong interactions. Each quark carries three types of color charge. These charges have nothing to do with the colors of visible light. Just as electrically-charged particles interact by exchanging photons, in strong interactions, color-charged particles interact by exchanging gluons.

Quarks Confined in Mesons and Baryons

Quarks and gluons cannot be isolated – they are confined in color-neutral particles called hadrons. This confinement (binding) results from multiple exchanges of gluons among the color-charged constituents. As

color-charged particles (quarks and gluons) move apart, the energy in the color-force field between them increases. This energy eventually is converted into additional quark-antiquark pairs. The quarks and antiquarks then combine into hadrons; these are the particles seen to emerge.

Two types of hadrons have been observed in nature **mesons** $q\bar{q}$ and **baryons** qqq . Among the many types of baryons observed are the proton (uud), antiproton ($\bar{u}\bar{u}\bar{d}$), and neutron (udd). Quark charges add in such a way as to make the proton have charge 1 and the neutron charge 0. Among the many types of mesons are the pion π^+ ($u\bar{d}$), kaon K^- ($s\bar{u}$), and B^0 ($d\bar{b}$).

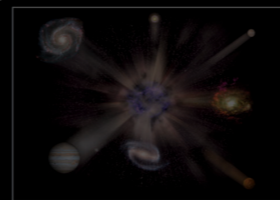
Learn more at ParticleAdventure.org



Unsolved Mysteries

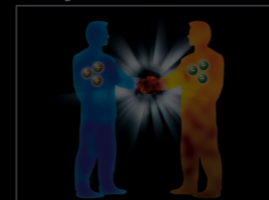
Driven by new puzzles in our understanding of the physical world, particle physicists are following paths to new wonders and startling discoveries. Experiments may even find extra dimensions of space, microscopic black holes, and/or evidence of string theory.

Why is the Universe Accelerating?



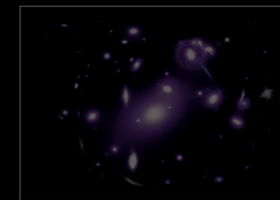
The expansion of the universe appears to be accelerating. Is this due to Einstein's Cosmological Constant? If not, will experiments reveal a new force of nature or even extra (hidden) dimensions of space?

Why No Antimatter?



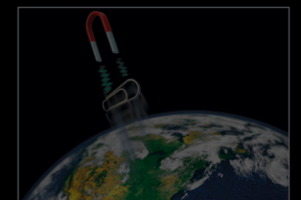
Matter and antimatter were created in the Big Bang. Why do we now see only matter except for the tiny amounts of antimatter that we make in the lab and observe in cosmic rays?

What is Dark Matter?



Invisible forms of matter make up much of the mass observed in galaxies and clusters of galaxies. Does this dark matter consist of new types of particles that interact very weakly with ordinary matter?

Are there Extra Dimensions?



An indication for extra dimensions may be the extreme weakness of gravity compared with the other three fundamental forces (gravity is so weak that a small magnet can pick up a paper clip overwhelming Earth's gravity).

Standard model of particle physics

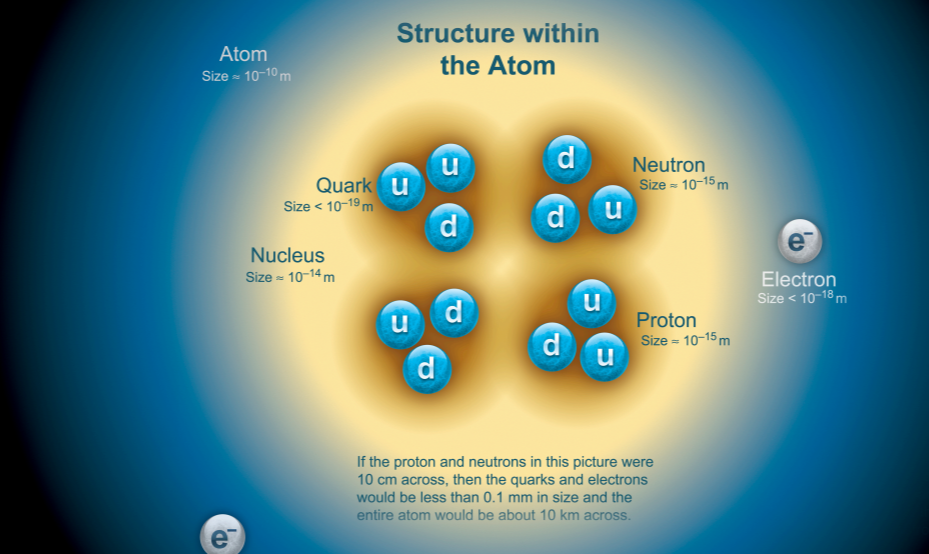
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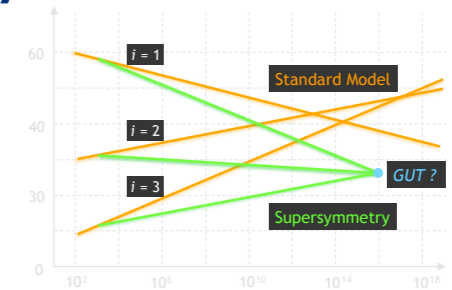
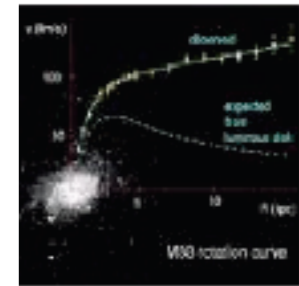
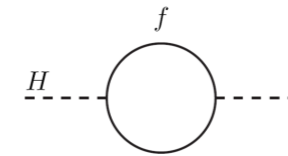
“We are at a very exciting and puzzling time for particle physics.”

(F. Gianotti, 2018 CERN New Year presentation.)

Exciting because the LHC discovered the Higgs boson (2012), the Standard Model is complete and it works beautifully!

Puzzling because the standard model is not a complete theory of particle physics. Several fundamental questions require physics beyond the standard model.

- Why is the Higgs boson so light?
- What is dark matter?
- What is the origin of the matter-antimatter asymmetry?
- Is there a grand unification of forces?
- How to include neutrino masses and oscillations?
- ...

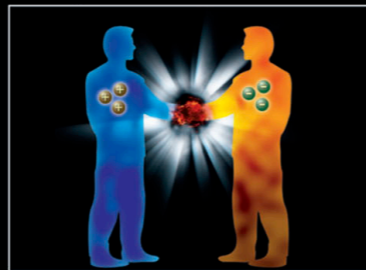


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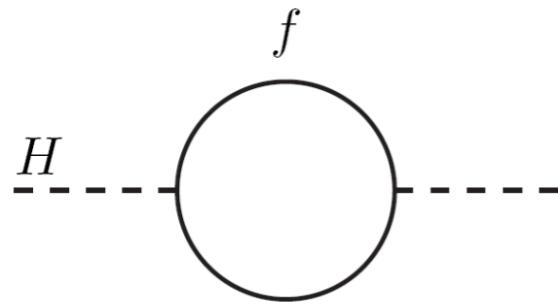


An indication for extra dimensions may be the extreme weakness of gravity compared with the other three fundamental forces (gravity is so weak that a small magnet can pick up a paper clip overwhelming Earth's gravity).

Strong motivation to search for new physics!

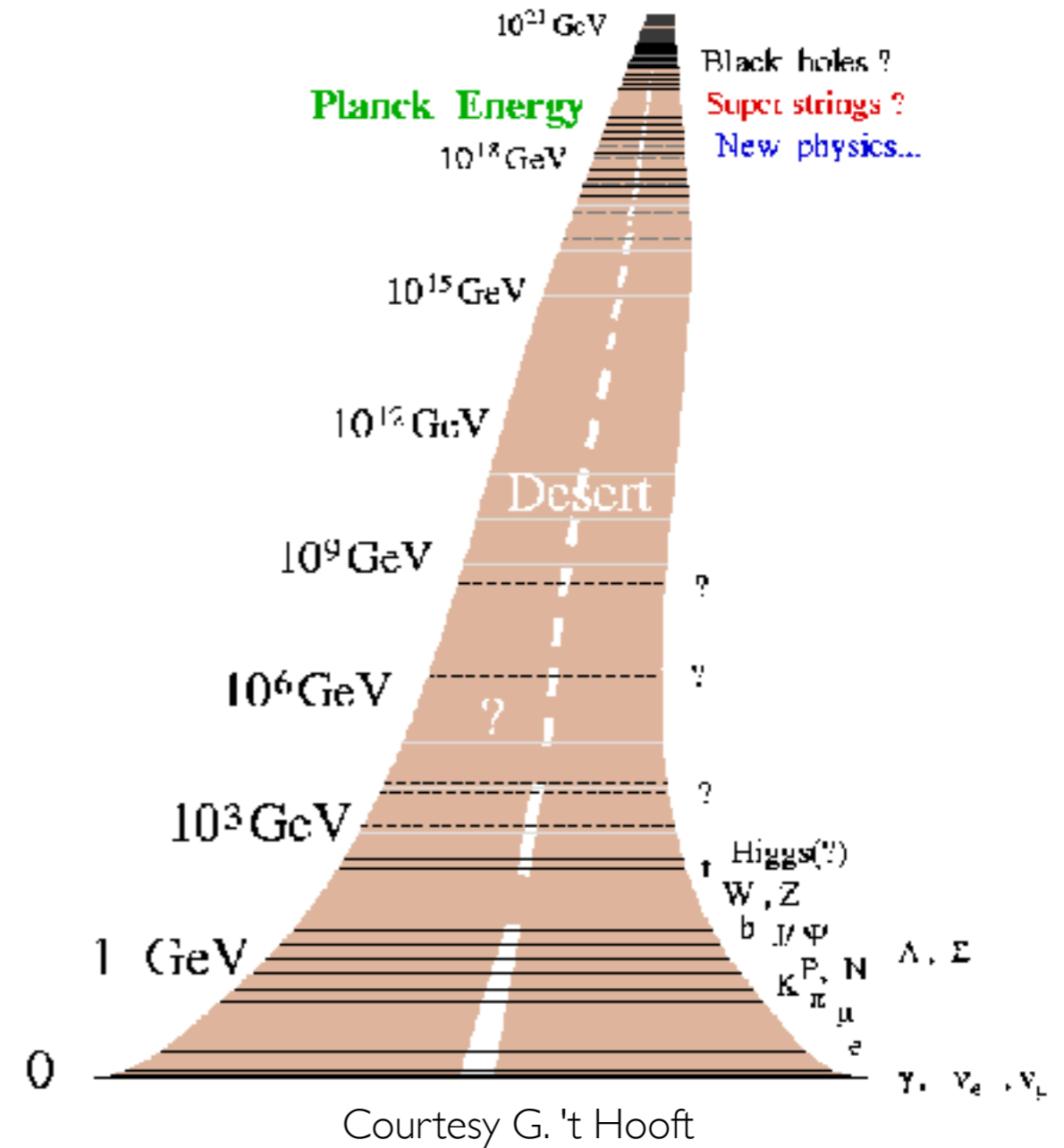
Why is the Higgs boson so light?

Very large quantum loop corrections to Higgs mass in any extension of the Standard Model.



$$\Delta m_H^2 = -\frac{|\lambda_f|^2}{8\pi^2} \Lambda_{UV}^2$$

New physics cut-off



Why is the Higgs boson so light?

$$\text{measured Higgs mass} \quad m_h^2 \approx \overset{\text{bare mass}}{m_{h0}^2} - \overset{\text{loop corrections}}{\Delta m_H^2}$$
$$(125 \text{ GeV})^2 \stackrel{?}{=} (10^{19} \text{ GeV})^2 - (10^{19} \text{ GeV})^2$$

Requires extremely precise cancelation of bare mass with correction term
“**fine-tuning**”

listening to your favorite radio needs the tuned frequency to match that of the radio channel:

radio freq. = 59.05871852091501091981287962349857612 kHz
tuned freq. = 59.05871852091501091981287962349857987 kHz



Why is the Higgs boson so much lighter than the Planck mass (or GUT energy, or heavy neutrino mass scale)?

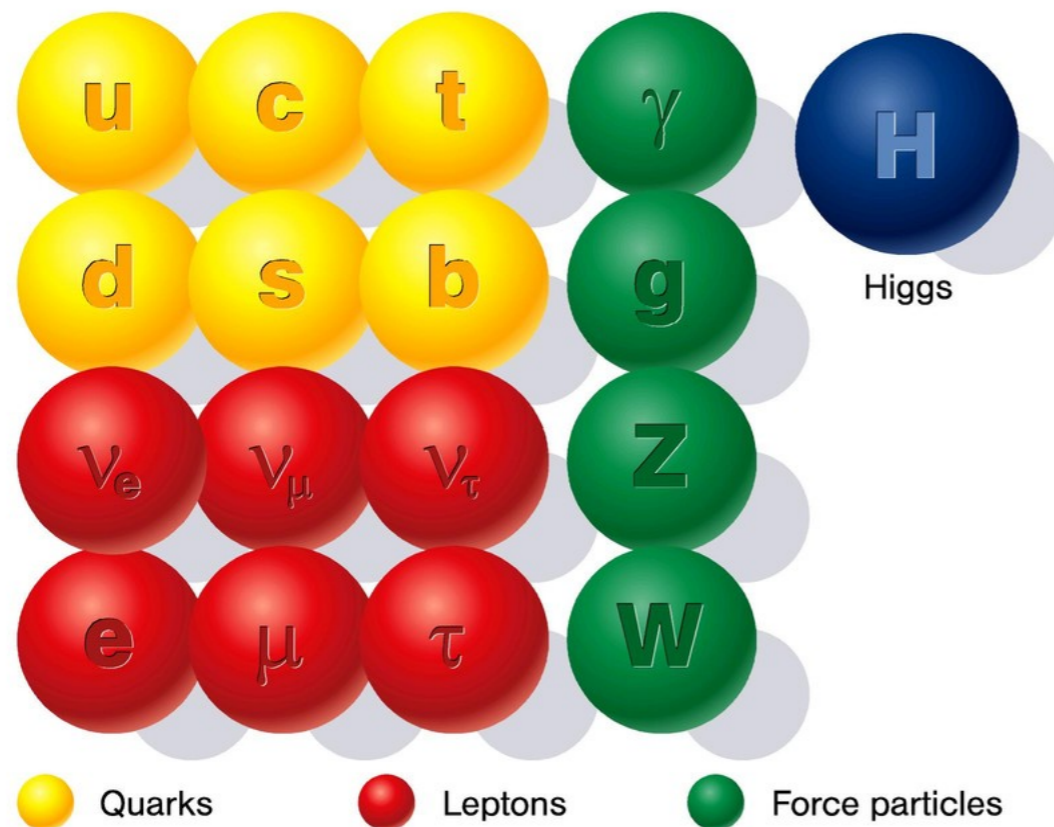
Supersymmetry

Supersymmetry (SUSY):

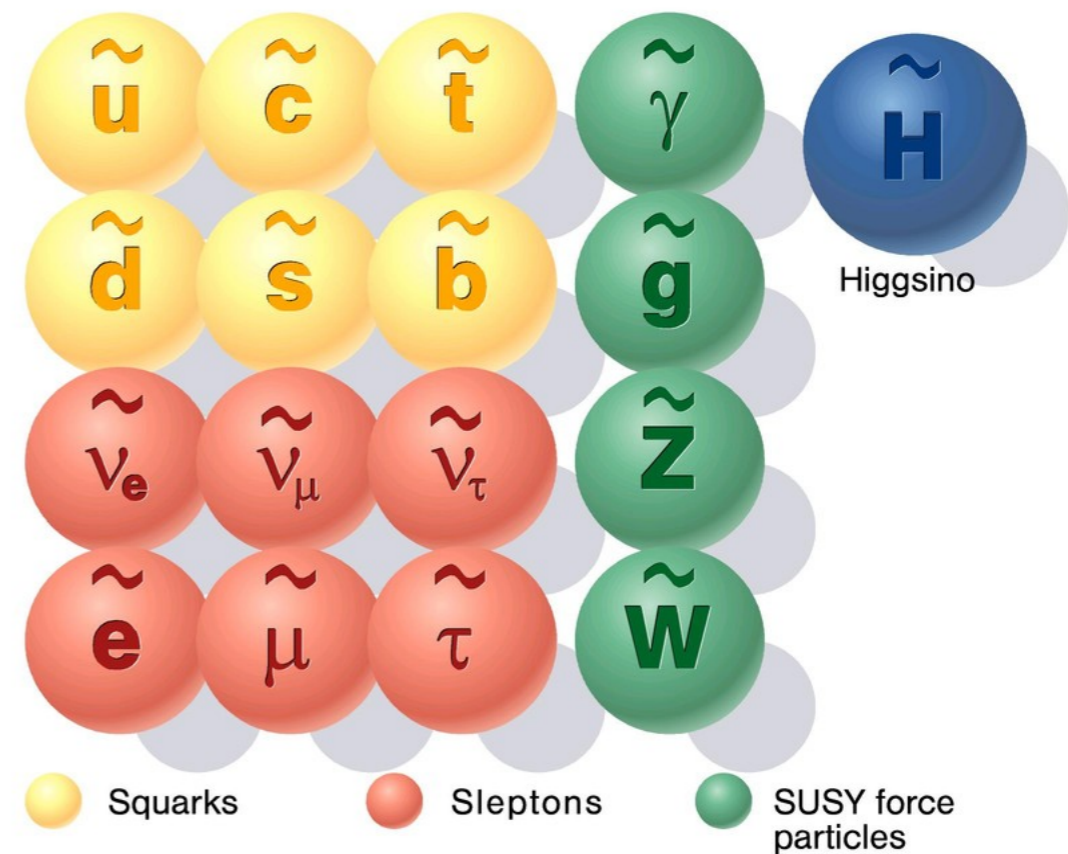
New symmetry between bosons and fermions.

For every SM particle introduce a supersymmetric partner with $\Delta\text{spin}=1/2$.

Standard particles



SUSY particles



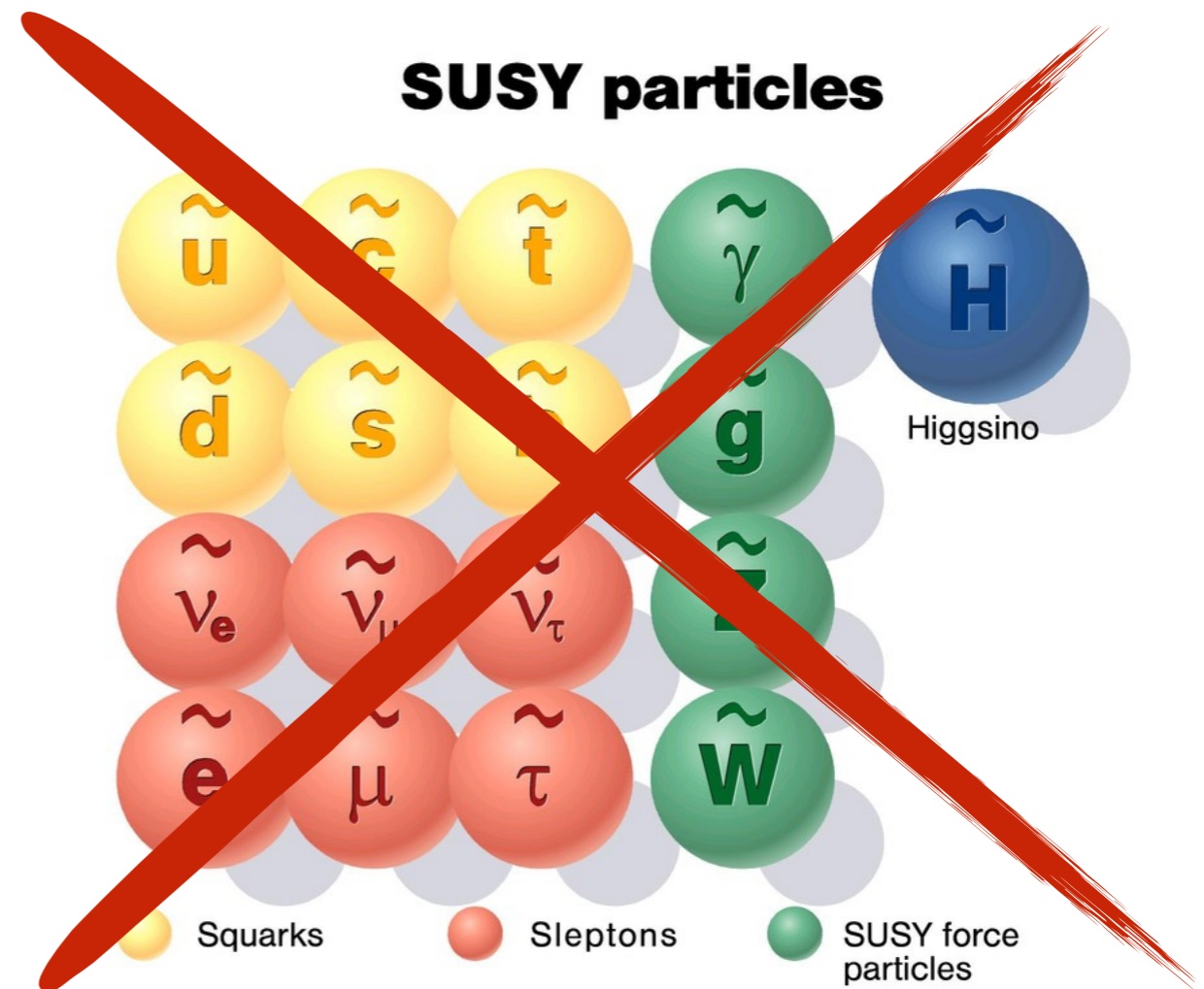
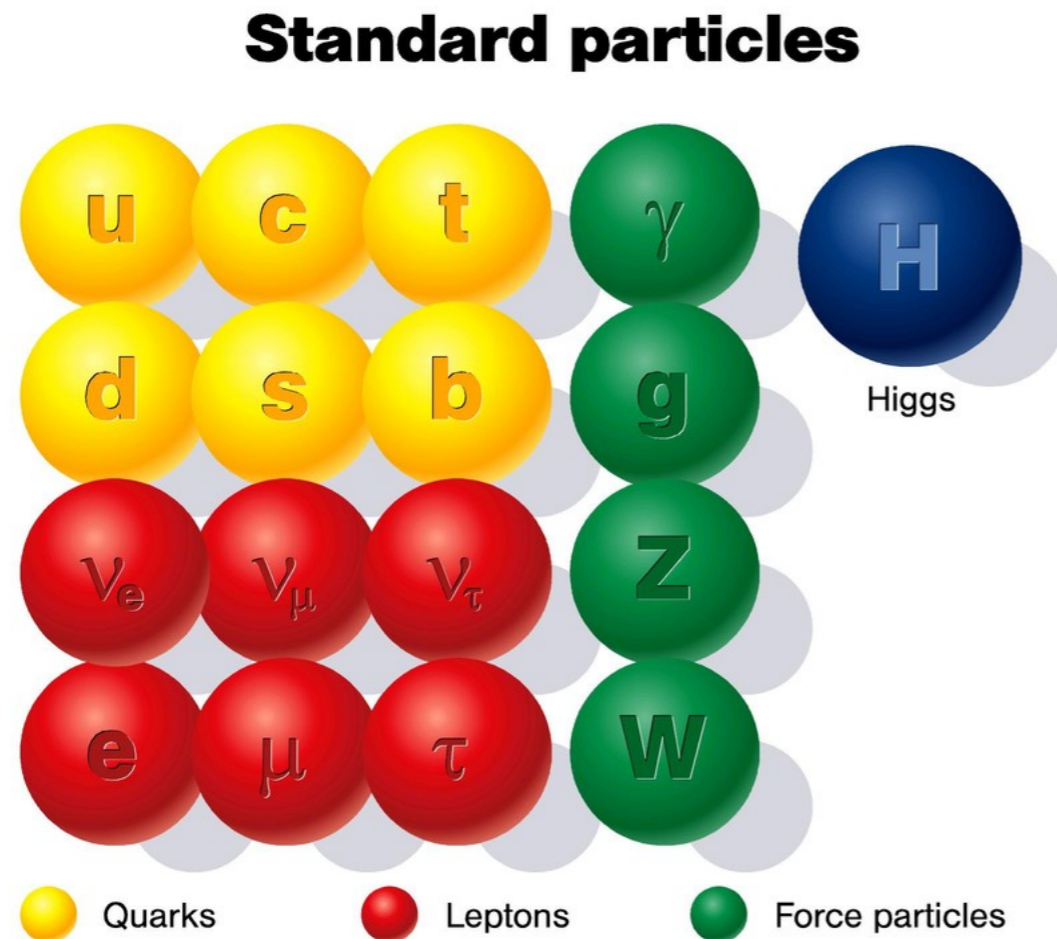
Supersymmetry

Supersymmetry (SUSY):

New symmetry between bosons and fermions.

~~For every SM particle introduce a supersymmetric partner with $\Delta\text{spin}=1/2$.~~

Equal number of bosonic and fermionic states



Supersymmetry

The undiscovered particles in the Minimal Supersymmetric Standard Model (MSSM) [SUSY primer, S. Martin]

| Names | Spin | P_R | Gauge Eigenstates | Mass Eigenstates |
|--------------------------|--------------|-------|---|---|
| Higgs bosons | 0 | +1 | H_u^0 H_d^0 H_u^+ H_d^- | h^0 H^0 A^0 H^\pm |
| squarks | 0 | -1 | \tilde{u}_L \tilde{u}_R \tilde{d}_L \tilde{d}_R | (same) |
| | | | \tilde{s}_L \tilde{s}_R \tilde{c}_L \tilde{c}_R | (same) |
| | | | \tilde{t}_L \tilde{t}_R \tilde{b}_L \tilde{b}_R | \tilde{t}_1 \tilde{t}_2 \tilde{b}_1 \tilde{b}_2 |
| sleptons | 0 | -1 | \tilde{e}_L \tilde{e}_R $\tilde{\nu}_e$ | (same) |
| | | | $\tilde{\mu}_L$ $\tilde{\mu}_R$ $\tilde{\nu}_\mu$ | (same) |
| | | | $\tilde{\tau}_L$ $\tilde{\tau}_R$ $\tilde{\nu}_\tau$ | $\tilde{\tau}_1$ $\tilde{\tau}_2$ $\tilde{\nu}_\tau$ |
| neutralinos | 1/2 | -1 | \tilde{B}^0 \tilde{W}^0 \tilde{H}_u^0 \tilde{H}_d^0 | \tilde{N}_1 \tilde{N}_2 \tilde{N}_3 \tilde{N}_4 |
| charginos | 1/2 | -1 | \tilde{W}^\pm \tilde{H}_u^\pm \tilde{H}_d^\pm | \tilde{C}_1^\pm \tilde{C}_2^\pm |
| gluino | 1/2 | -1 | \tilde{g} | (same) |
| goldstino (gravitino) | 1/2 (3/2) | -1 | \tilde{G} | (same) |

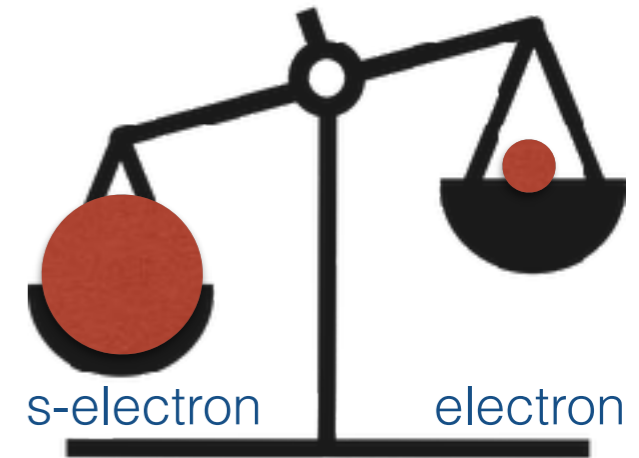
bino, neutral wino & neutral higgsinos mix to form four neutralinos.

charged wino & charged higgsinos mix to form four charginos.

Supersymmetry

Broken symmetry

- Otherwise, $m(\text{s-electron}) = m(\text{electron})$, etc.
- Several theories for supersymmetry breaking (gauge mediated, gravity mediated, etc.).
- Minimal supersymmetric standard model (MSSM):
 1 new parameter
 soft supersymmetry breaking (our ignorance!): 124 parameters
 (SM comes with 19 parameters)



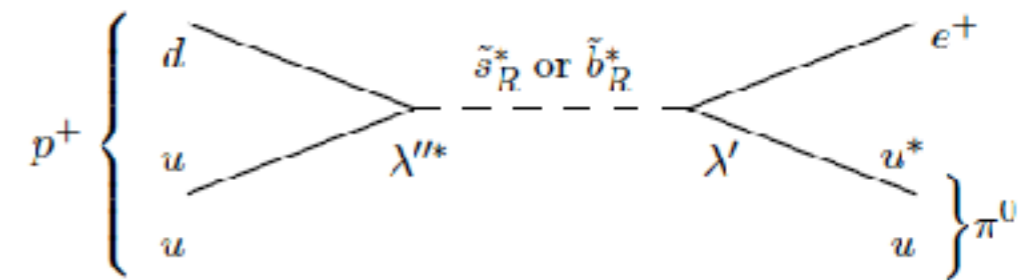
SUSY: Lepton (L) and baryon (B) number violation allowed, proton decay

R-parity = +1 (-1) SM (SUSY) particles

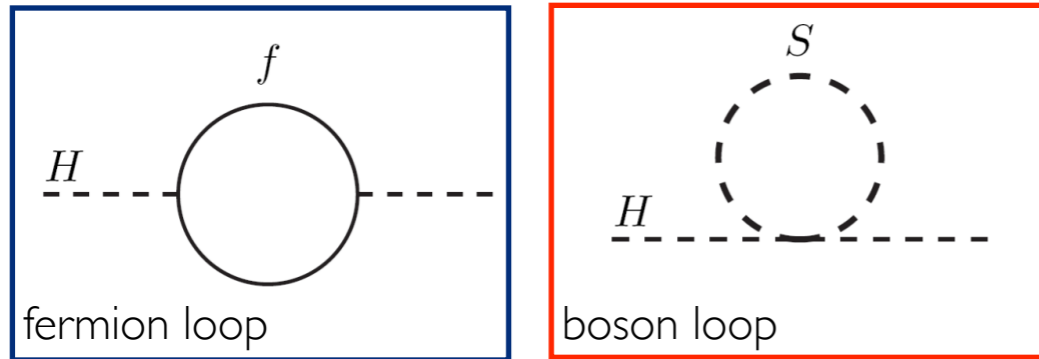
$$P_R = (-1)^{3(B-L)+2s}$$

R-parity conservation (RPC) implies:

- ▶ eliminate B and L number violating terms,
- ▶ SUSY production only in even numbers (typically 2),
- ▶ lightest supersymmetric particle (LSP) is stable
 if electrically and color neutral then **DM candidate**, and LSP escapes detection
 → missing (transverse) momentum.



SUSY solution to “Why is the Higgs boson so



fermion and boson loops contribute with different signs to the Higgs radiative corrections; fermion-boson symmetry protects the scalar Higgs.

With SUSY:

$$m_h^2 \approx m_{h0}^2 + \frac{\lambda_f^2}{8\pi^2} N_c^f (m_{\tilde{f}}^2 - m_f^2) \ln(\Lambda^2/m_{\tilde{f}}^2)$$

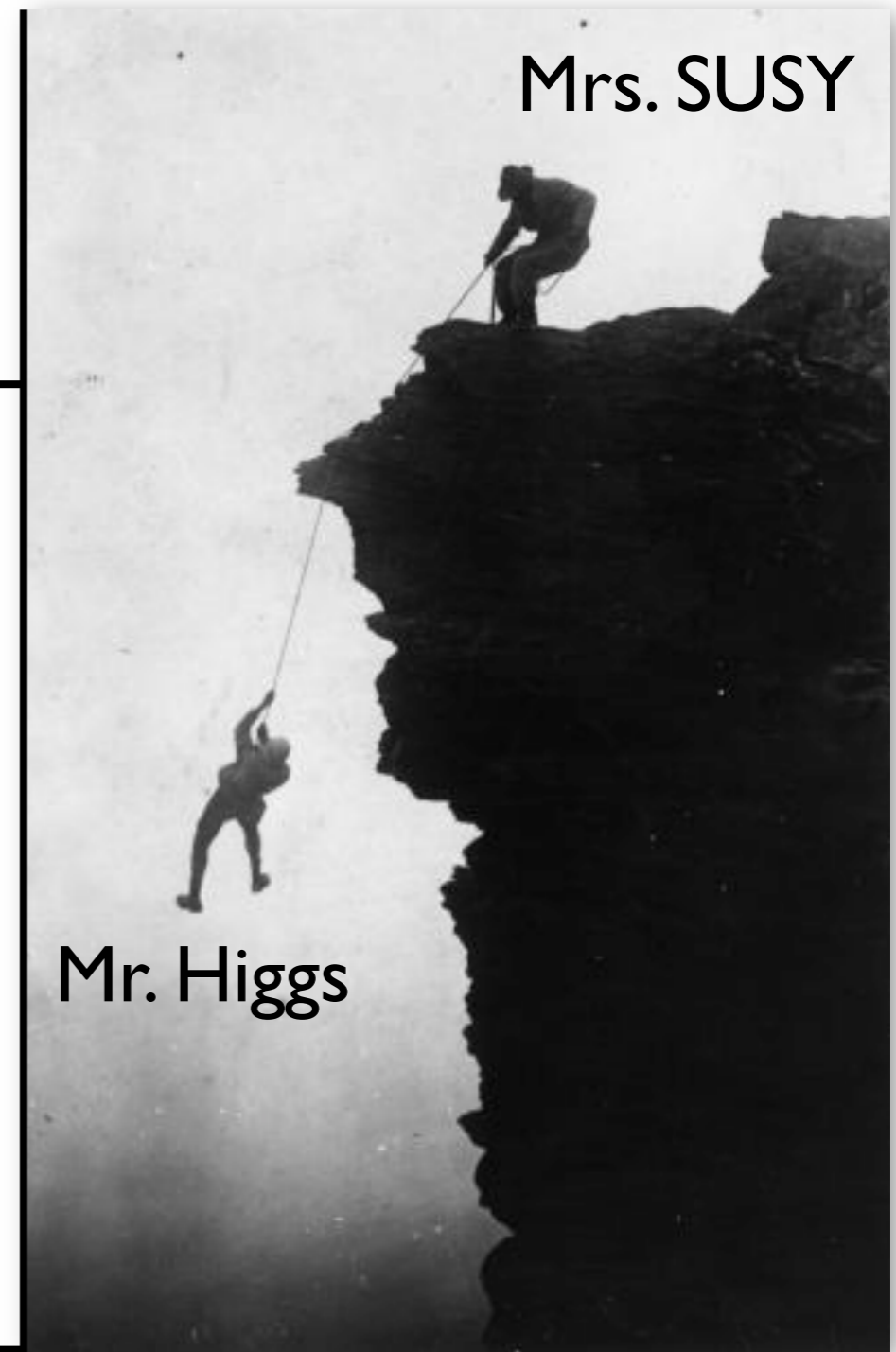


Ideally, small mass difference btw. SM and SUSY partner particles.

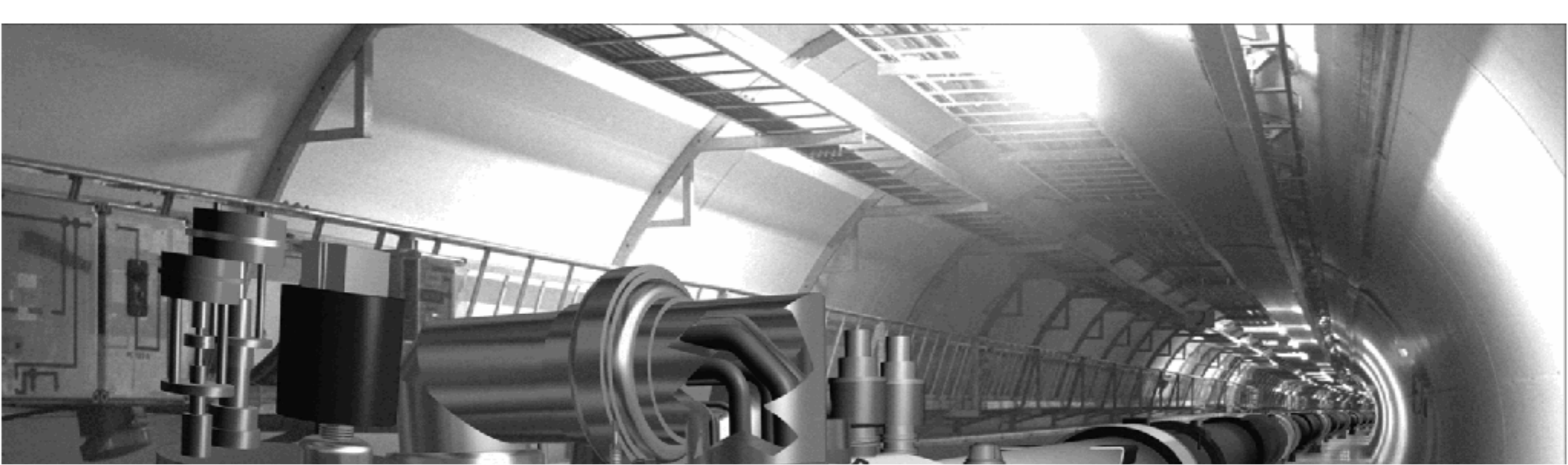
→ Expect SUSY particles close to weak scale.

weak scale

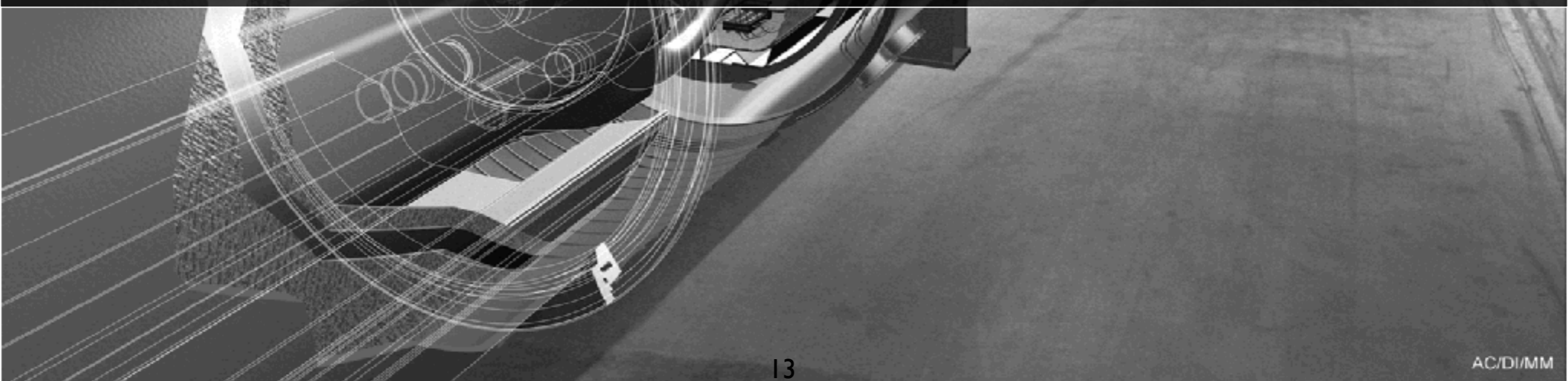
GUT scale
Planck scale



fundamental scalar energy (cliff)



Experimental Searches at the Large Hadron Collider



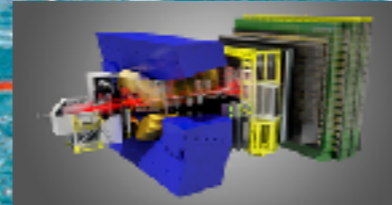
Large Hadron Collider (LHC)

proton-proton (also HI) collisions

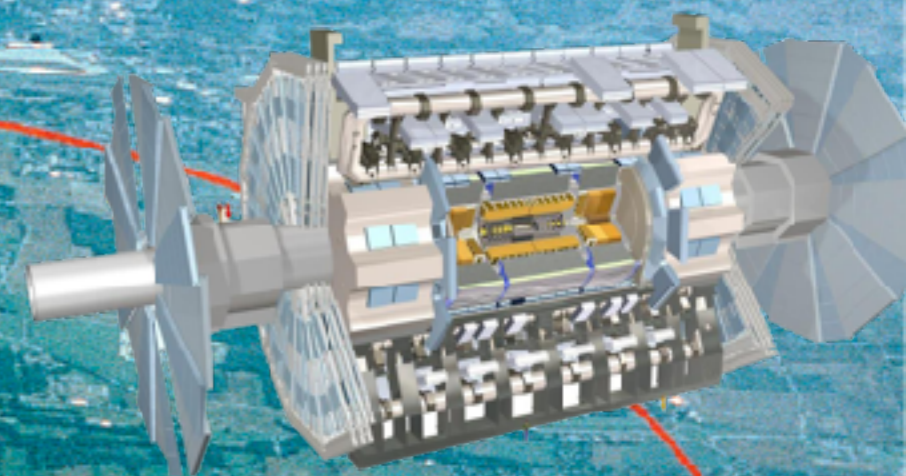
Run1: 2010-2012, centre-of-mass energy 7 then 8 TeV

Run2: 2015-2018, centre-of-mass energy 13 TeV

LHCb



ATLAS



CMS



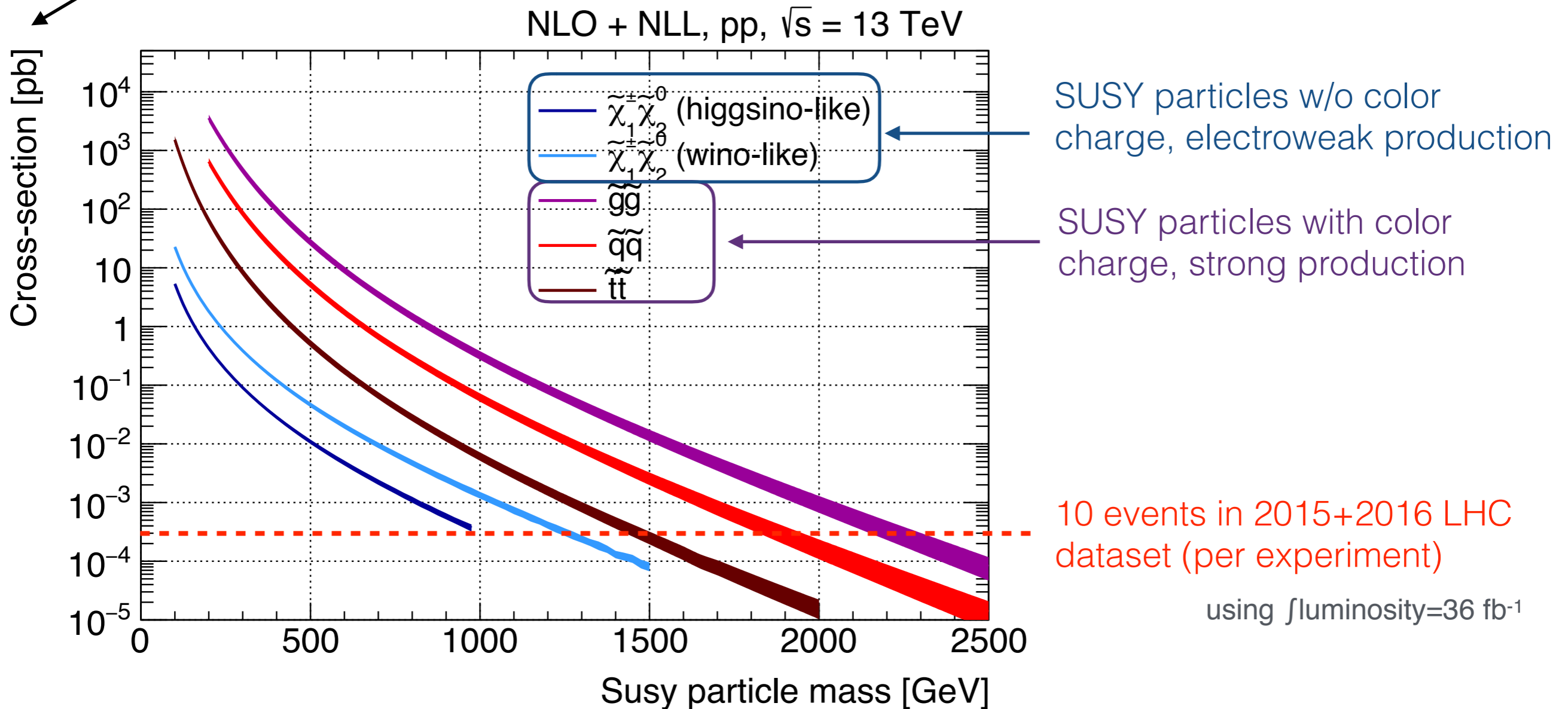
2 general purpose detectors



ALICE

SUSY particle production

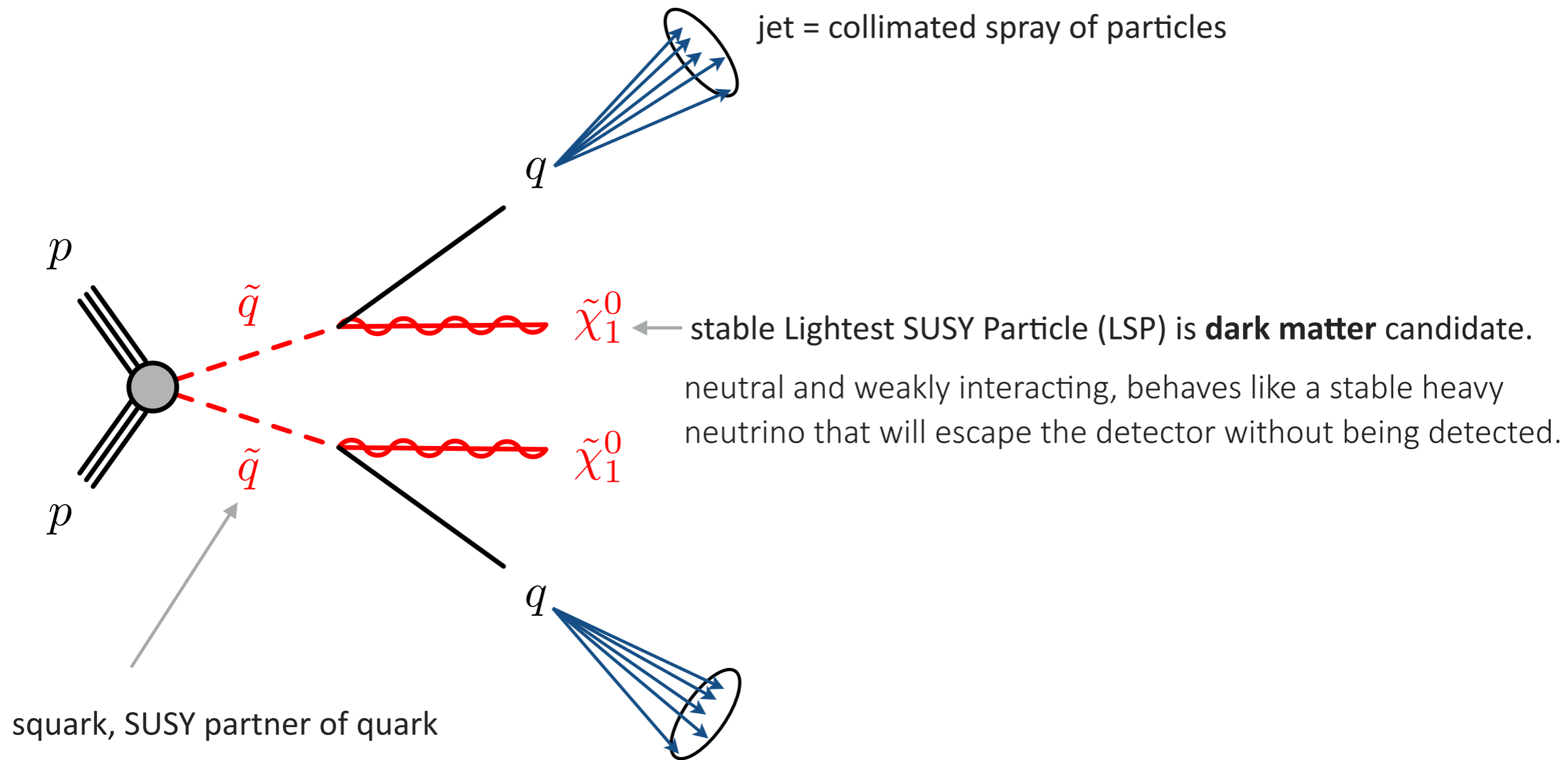
Number of collisions events with SUSY particles
 = **cross-section** (physics) \times \int **luminosity** (LHC machine, beam)



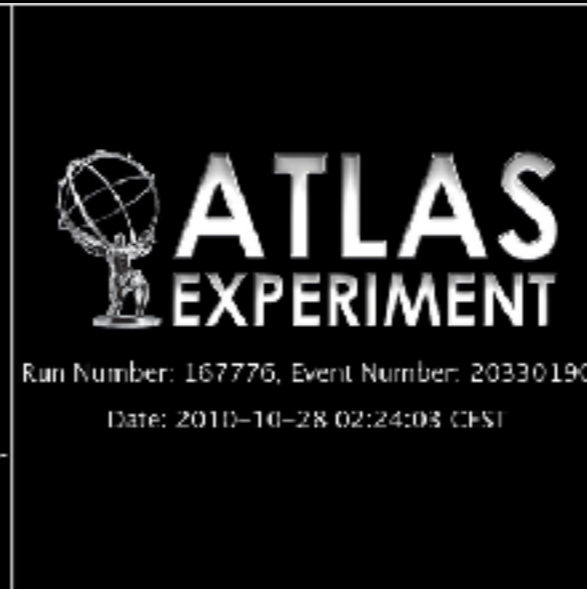
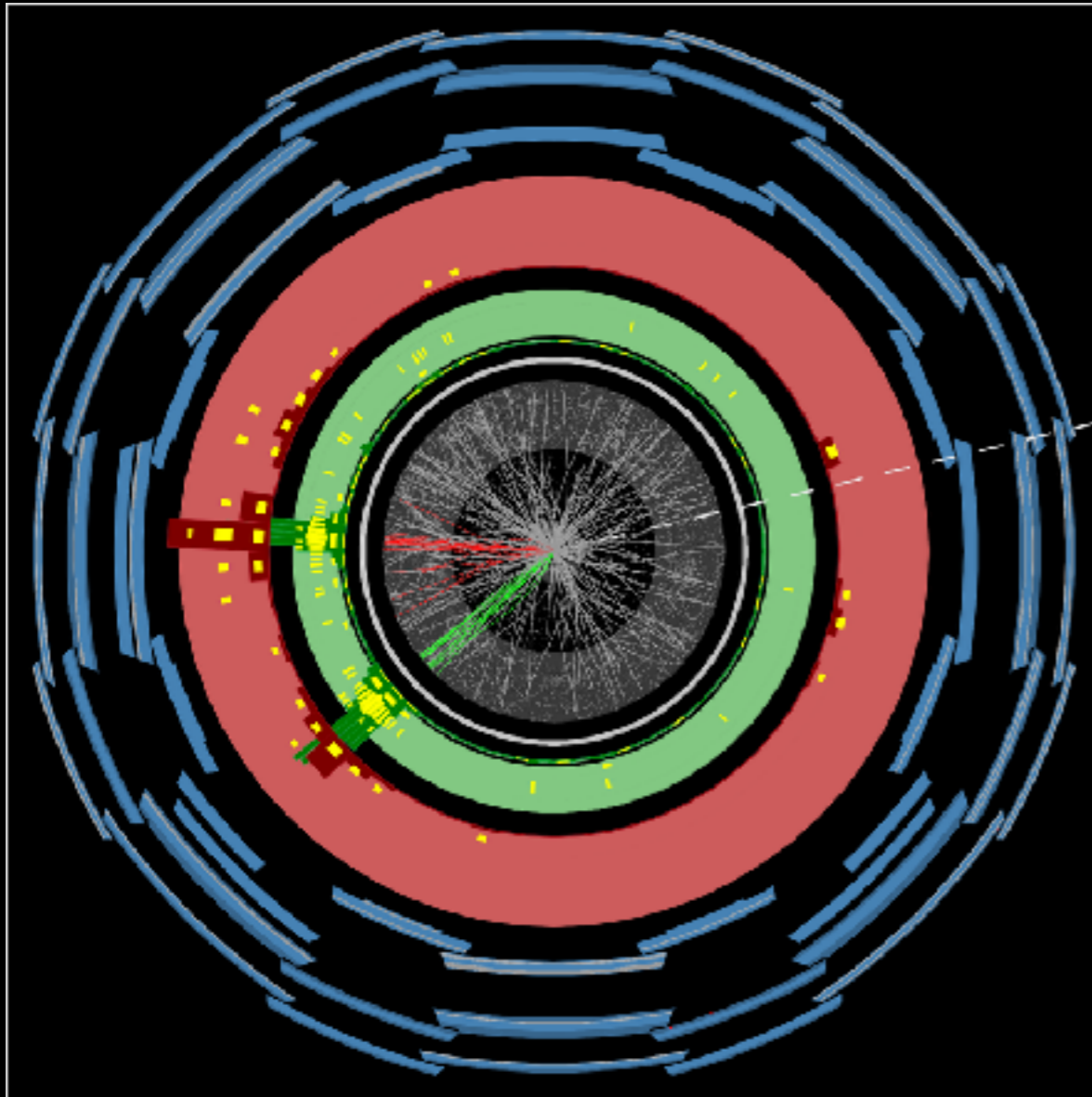
SUSY particle cross sections

[NLO + NLL Tool, C. Borschensky et al]

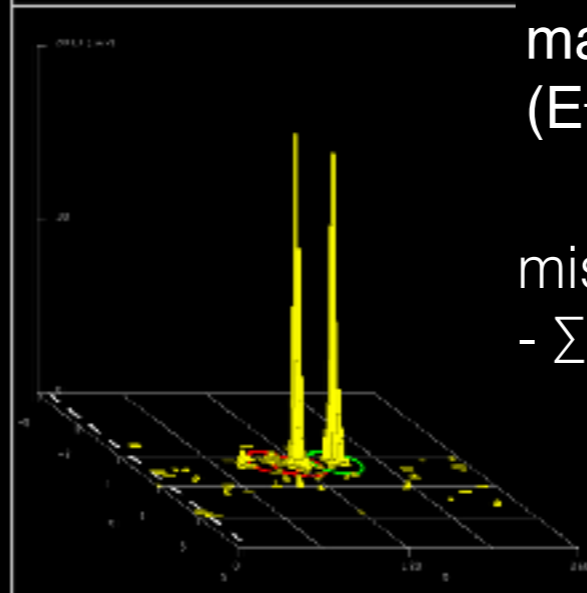
Example squark production and decay



Candidate event picked up in 2010 dataset



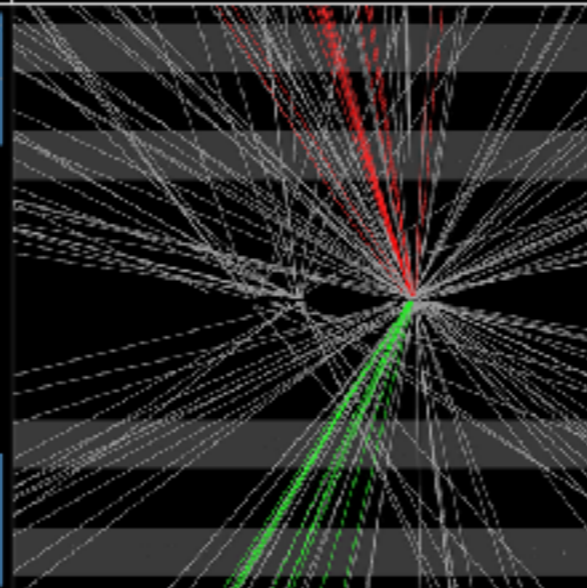
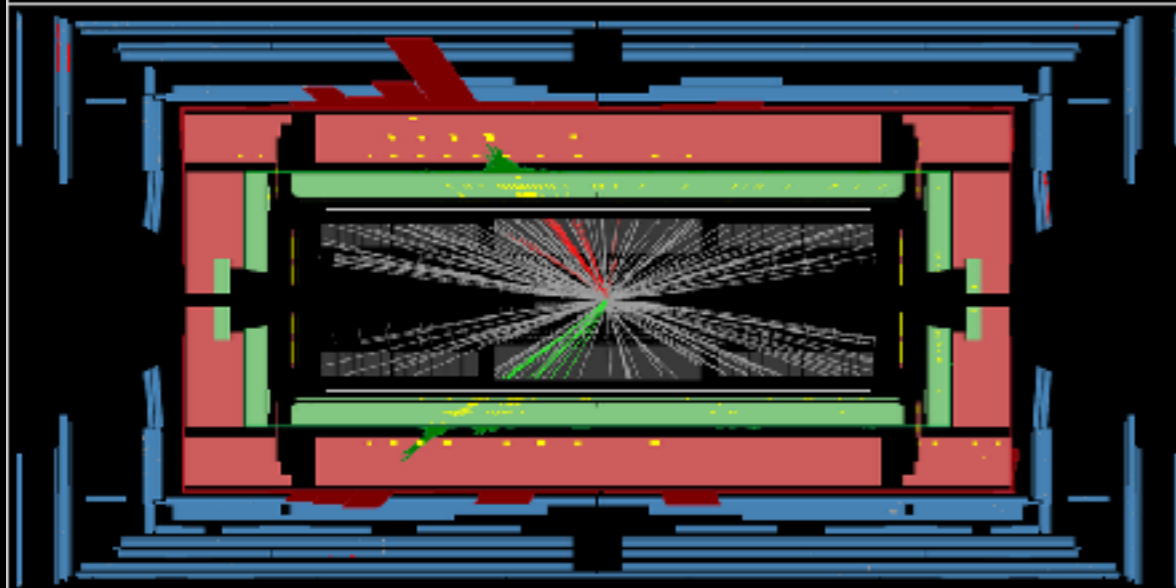
Two jets each with transverse momentum (p_T) of ~ 170 GeV



$\text{mag}(\text{missing transverse momentum})$ (E_T^{miss}) ~ 330 GeV.

missing transverse momentum =
 $-\sum p_T(\text{all visible particles in detector})$

transverse: momentum along beam direction unknown in initial state in pp collisions.



All-hadronic search for squarks and gluinos

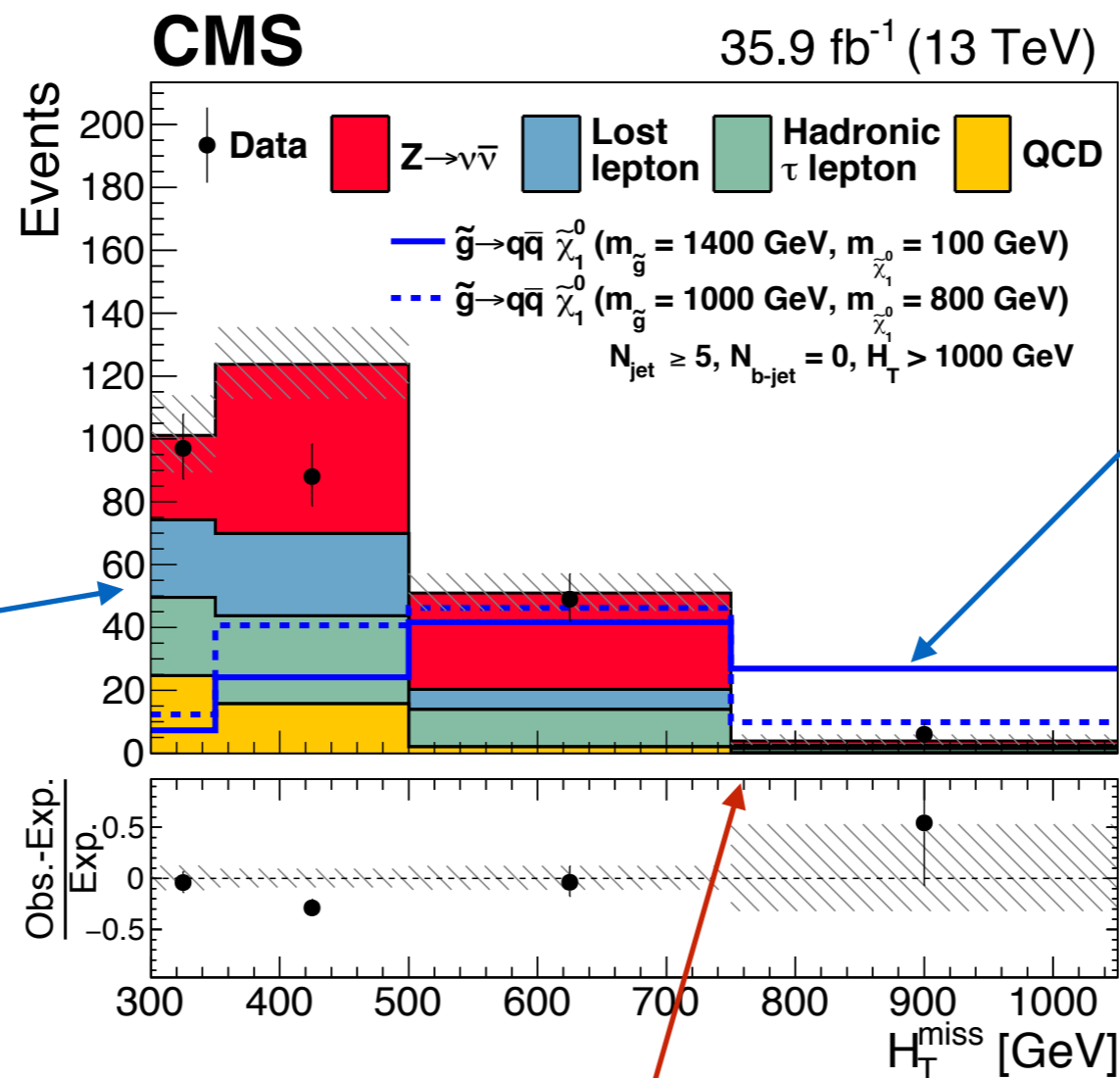
Trigger on missing transverse momentum ($E_T^{\text{miss}}, H_T^{\text{miss}}$)

online selection,
record data on tape

Veto events with identified electrons and muons

Require several jets with high momentum, large E_T^{miss} ,
large effective mass $m_{\text{eff}} = \sum p_T(\text{jets}) + E_T^{\text{miss}}$

separate signal
from background

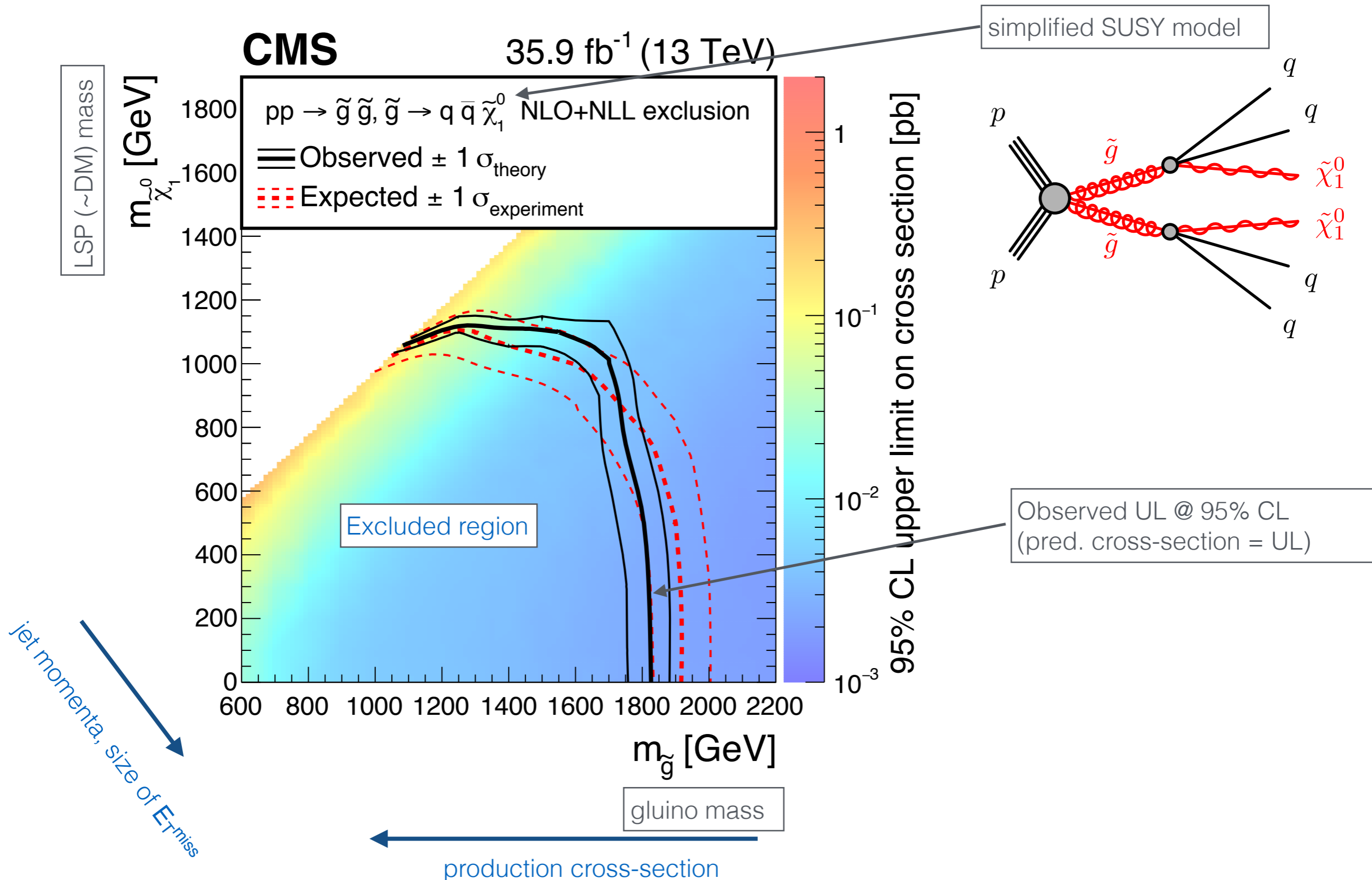


Critical to have robust
background estimation

Expectation from a signal model

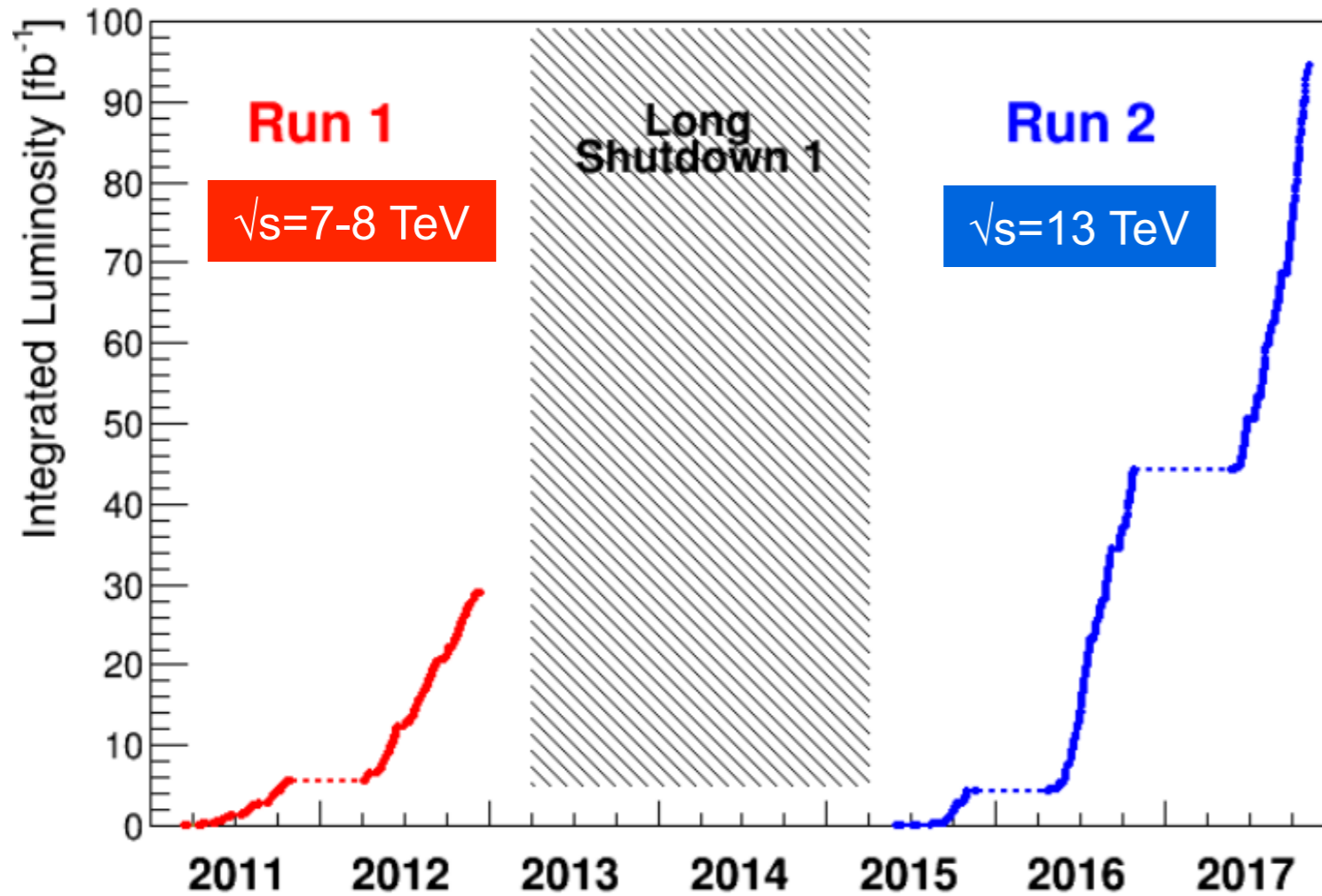
Select events with characteristic SUSY signature and unlikely due to background.
Compare event count with background prediction.

Statistical Interpretation



Overview

Huge LHC sample of proton-proton collisions



LHC peak luminosity $\sim 2 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
(design 1×10^{34})

$\sqrt{s}=8 \text{ TeV}$ dataset (2012): 5M top quark pairs, 0.4M Higgs bosons, and would expect 2.1M (500) gluino pairs @ 300 GeV (1 TeV) mass,

$\sqrt{s}=13 \text{ TeV}$ dataset (2015-2017): 75M top quark pairs, 4.5M Higgs bosons, and would expect 30k (90) gluino pairs @ 1 TeV (2 TeV) mass

Large SUSY search program @ LHC

ATLAS SUSY Searches* - 95% CL Lower Limits

December 2017

ATLAS Preliminary

$\sqrt{s} = 7, 8, 13$ TeV

| Model | e, μ, τ, γ | Jets | E_T^{miss} | $\int \mathcal{L} dt [\text{fb}^{-1}]$ | Mass limit | | Reference | |
|--|---|-------------------------------|---------------------|--|----------------------------|---|---|------------------------------------|
| | | | | | $\sqrt{s} = 7, 8$ TeV | $\sqrt{s} = 13$ TeV | | |
| Inclusive Searches | $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 36.1 | \tilde{q} 1.57 TeV | $m(\tilde{\chi}_1^0) < 200$ GeV, $m(1^{\text{st}} \text{ gen. } \tilde{q}) = m(2^{\text{nd}} \text{ gen. } \tilde{q})$ | 1712.02332 |
| | $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (compressed) | mono-jet | 1-3 jets | Yes | 36.1 | \tilde{q} 710 GeV | $m(\tilde{q}) - m(\tilde{\chi}_1^0) < 5$ GeV | 1711.03301 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 36.1 | \tilde{g} 2.02 TeV | $m(\tilde{\chi}_1^0) < 200$ GeV | 1712.02332 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^{\pm} \rightarrow qqW^{\pm}\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 36.1 | \tilde{g} 2.01 TeV | $m(\tilde{\chi}_1^0) < 200$ GeV, $m(\tilde{\chi}^{\pm}) = 0.5(m(\tilde{\chi}_1^0) + m(\tilde{g}))$ | 1712.02332 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$ | $ee, \mu\mu$ | 2 jets | Yes | 14.7 | \tilde{g} 1.7 TeV | $m(\tilde{\chi}_1^0) < 300$ GeV, | 1611.05791 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq(\ell\ell/\nu\nu)\tilde{\chi}_1^0$ | $3 e, \mu$ | 4 jets | - | 36.1 | \tilde{g} 1.87 TeV | $m(\tilde{\chi}_1^0) = 0$ GeV | 1706.03731 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$ | 0 | 7-11 jets | Yes | 36.1 | \tilde{g} 1.8 TeV | $m(\tilde{\chi}_1^0) < 400$ GeV | 1708.02794 |
| | GMSB ($\tilde{\ell}$ NLSP) | $1-2 \tau + 0-1 \ell$ | 0-2 jets | Yes | 3.2 | \tilde{g} 2.0 TeV | $c\tau(\text{NLSP}) < 0.1$ mm | 1607.05979 |
| | GGM (bino NLSP) | 2γ | - | Yes | 36.1 | \tilde{g} 2.15 TeV | $c\tau(\text{NLSP}) < 0.1$ mm, | ATLAS-CONF-2017-080 |
| | GGM (higgsino-bino NLSP) | γ | 2 jets | Yes | 36.1 | \tilde{g} 2.05 TeV | $m(\tilde{\chi}_1^0) = 1700$ GeV, $c\tau(\text{NLSP}) < 0.1$ mm, $\mu > 0$ | ATLAS-CONF-2017-080 |
| Gravitino LSP | 0 | mono-jet | Yes | 20.3 | $F^{1/2}$ scale 865 GeV | $m(\tilde{G}) > 1.8 \times 10^{-4}$ eV, $m(\tilde{g}) = m(\tilde{q}) = 1.5$ TeV | 1502.01518 | |
| 3 rd gen. \tilde{g} med. | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$ | 0 | 3 b | Yes | 36.1 | \tilde{g} 1.92 TeV | $m(\tilde{\chi}_1^0) < 600$ GeV | 1711.01901 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$ | $0-1 e, \mu$ | 3 b | Yes | 36.1 | \tilde{g} 1.97 TeV | $m(\tilde{\chi}_1^0) < 200$ GeV | 1711.01901 |
| 3 rd gen. squarks direct production | $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ | 0 | 2 b | Yes | 36.1 | \tilde{b}_1 950 GeV | $m(\tilde{\chi}_1^0) < 420$ GeV | 1708.09266 |
| | $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^{\pm}$ | $2 e, \mu$ (SS) | 1 b | Yes | 36.1 | \tilde{b}_1 275-700 GeV | $m(\tilde{\chi}_1^0) < 200$ GeV, $m(\tilde{\chi}_1^{\pm}) = m(\tilde{\chi}_1^0) + 100$ GeV | 1706.03731 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}_1^{\pm}$ | $0-2 e, \mu$ | 1-2 b | Yes | 4.7/13.3 | \tilde{t}_1 117-170 GeV 200-720 GeV | $m(\tilde{\chi}_1^{\pm}) = 2m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0) = 55$ GeV | 1209.2102, ATLAS-CONF-2016-077 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$ | $0-2 e, \mu$ | 0-2 jets/1-2 b | Yes | 20.3/36.1 | \tilde{t}_1 90-198 GeV 0.195-1.0 TeV | $m(\tilde{\chi}_1^0) = 1$ GeV | 1506.08616, 1709.04183, 1711.11520 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$ | 0 | mono-jet | Yes | 36.1 | \tilde{t}_1 90-430 GeV | $m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 5$ GeV | 1711.03301 |
| | $\tilde{t}_1\tilde{t}_1$ (natural GMSB) | $2 e, \mu$ (Z) | 1 b | Yes | 20.3 | \tilde{t}_1 150-600 GeV | $m(\tilde{\chi}_1^0) > 150$ GeV | 1403.5222 |
| | $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$ | $3 e, \mu$ (Z) | 1 b | Yes | 36.1 | \tilde{t}_2 290-790 GeV | $m(\tilde{\chi}_1^0) = 0$ GeV | 1706.03986 |
| | $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$ | $1-2 e, \mu$ | 4 b | Yes | 36.1 | \tilde{t}_2 320-880 GeV | $m(\tilde{\chi}_1^0) = 0$ GeV | 1706.03986 |
| EW direct | $\tilde{\ell}_{L,R}, \tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$ | $2 e, \mu$ | 0 | Yes | 36.1 | $\tilde{\ell}$ 90-500 GeV | $m(\tilde{\chi}_1^0) = 0$ | ATLAS-CONF-2017-039 |
| | $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\mp}, \tilde{\chi}_1^{\pm} \rightarrow \tilde{\ell}\nu(\ell\bar{\nu})$ | $2 e, \mu$ | 0 | Yes | 36.1 | $\tilde{\chi}_1^{\pm}$ 750 GeV | $m(\tilde{\chi}_1^0) = 0, m(\tilde{\ell}, \bar{\nu}) = 0.5(m(\tilde{\chi}_1^{\pm}) + m(\tilde{\chi}_1^0))$ | ATLAS-CONF-2017-039 |
| | $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\mp}/\tilde{\chi}_2^0, \tilde{\chi}_1^{\pm} \rightarrow \tilde{\tau}\nu(\tau\bar{\nu}), \tilde{\chi}_2^0 \rightarrow \tilde{\tau}\tau(\nu\bar{\nu})$ | 2τ | - | Yes | 36.1 | $\tilde{\chi}_1^{\pm}$ 760 GeV | $m(\tilde{\chi}_1^0) = 0, m(\tilde{\tau}, \bar{\nu}) = 0.5(m(\tilde{\chi}_1^{\pm}) + m(\tilde{\chi}_1^0))$ | 1708.07875 |
| | $\tilde{\chi}_1^{\pm}\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_L\nu\tilde{\ell}_L(\ell\bar{\nu}), \tilde{\nu}\tilde{\ell}_L(\ell\bar{\nu})$ | $3 e, \mu$ | 0 | Yes | 36.1 | $\tilde{\chi}_1^{\pm}, \tilde{\chi}_2^0$ 1.13 TeV | $m(\tilde{\chi}_1^{\pm}) = m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0) = 0, m(\tilde{\ell}, \bar{\nu}) = 0.5(m(\tilde{\chi}_1^{\pm}) + m(\tilde{\chi}_1^0))$ | ATLAS-CONF-2017-039 |
| | $\tilde{\chi}_1^{\pm}\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 Z\tilde{\chi}_1^0$ | $2-3 e, \mu$ | 0-2 jets | Yes | 36.1 | $\tilde{\chi}_1^{\pm}, \tilde{\chi}_2^0$ 580 GeV | $m(\tilde{\chi}_1^{\pm}) = m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0) = 0, \tilde{\ell}$ decoupled | ATLAS-CONF-2017-039 |
| | $\tilde{\chi}_1^{\pm}\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 h\tilde{\chi}_1^0, h \rightarrow b\tilde{b}/WW/\tau\tau/\gamma\gamma$ | e, μ, γ | 0-2 b | Yes | 20.3 | $\tilde{\chi}_1^{\pm}, \tilde{\chi}_2^0$ 270 GeV | $m(\tilde{\chi}_1^{\pm}) = m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0) = 0, \tilde{\ell}$ decoupled | 1501.07110 |
| | $\tilde{\chi}_2^0\tilde{\chi}_3^0, \tilde{\chi}_2^0 \rightarrow \tilde{\ell}_R\ell$ | $4 e, \mu$ | 0 | Yes | 20.3 | $\tilde{\chi}_2^0, \tilde{\chi}_3^0$ 635 GeV | $m(\tilde{\chi}_2^0) = m(\tilde{\chi}_3^0), m(\tilde{\chi}_1^0) = 0, m(\tilde{\ell}, \bar{\nu}) = 0.5(m(\tilde{\chi}_2^0) + m(\tilde{\chi}_1^0))$ | 1405.5086 |
| | GGM (wino NLSP) weak prod., $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ | $1 e, \mu + \gamma$ | - | Yes | 20.3 | \tilde{W} 115-370 GeV | $c\tau < 1$ mm | 1507.05493 |
| GGM (bino NLSP) weak prod., $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ | 2γ | - | Yes | 36.1 | \tilde{W} 1.06 TeV | $c\tau < 1$ mm | ATLAS-CONF-2017-080 | |
| Long-lived particles | Direct $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\mp}$ prod., long-lived $\tilde{\chi}_1^{\pm}$ | Disapp. trk | 1 jet | Yes | 36.1 | $\tilde{\chi}_1^{\pm}$ 460 GeV | $m(\tilde{\chi}_1^{\pm}) - m(\tilde{\chi}_1^0) \sim 160$ MeV, $\tau(\tilde{\chi}_1^{\pm}) = 0.2$ ns | 1712.02118 |
| | Direct $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\mp}$ prod., long-lived $\tilde{\chi}_1^{\pm}$ | dE/dx trk | - | Yes | 18.4 | $\tilde{\chi}_1^{\pm}$ 495 GeV | $m(\tilde{\chi}_1^{\pm}) - m(\tilde{\chi}_1^0) \sim 160$ MeV, $\tau(\tilde{\chi}_1^{\pm}) < 15$ ns | 1506.05332 |
| | Stable, stopped \tilde{g} R-hadron | 0 | 1-5 jets | Yes | 27.9 | \tilde{g} 850 GeV | $m(\tilde{\chi}_1^0) = 100$ GeV, $10 \mu\text{s} < \tau(\tilde{g}) < 1000$ s | 1310.6584 |
| | Stable \tilde{g} R-hadron | trk | - | - | 3.2 | \tilde{g} 1.58 TeV | | 1606.05129 |
| | Metastable \tilde{g} R-hadron | dE/dx trk | - | - | 3.2 | \tilde{g} 1.57 TeV | $m(\tilde{\chi}_1^0) = 100$ GeV, $\tau > 10$ ns | 1604.04520 |
| | Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq\tilde{\chi}_1^0$ | displ. vtx | - | Yes | 32.8 | \tilde{g} 2.37 TeV | $\tau(\tilde{g}) = 0.17$ ns, $m(\tilde{\chi}_1^0) = 100$ GeV | 1710.04901 |
| | GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{\ell}, \bar{\mu}) + \tau(e, \mu)$ | $1-2 \mu$ | - | - | 19.1 | $\tilde{\chi}_1^0$ 537 GeV | $10 < \tan\beta < 50$ | 1411.6795 |
| | GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$ | 2γ | - | Yes | 20.3 | $\tilde{\chi}_1^0$ 440 GeV | $1 < \tau(\tilde{\chi}_1^0) < 3$ ns, SPS8 model | 1409.5542 |
| $\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow e\bar{e}\nu/\mu\bar{\mu}\nu$ | displ. $ee/\mu\mu/\mu\nu$ | - | - | 20.3 | $\tilde{\chi}_1^0$ 1.0 TeV | $7 < c\tau(\tilde{\chi}_1^0) < 740$ mm, $m(\tilde{g}) = 1.3$ TeV | 1504.05162 | |
| RPV | LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\epsilon\tau/\mu\tau$ | $e\mu, \epsilon\tau, \mu\tau$ | - | - | 3.2 | $\tilde{\nu}_\tau$ 1.9 TeV | $\lambda_{511} = 0.11, \lambda_{132/133/233} = 0.07$ | 1607.08079 |
| | Bilinear RPV CMSSM | $2 e, \mu$ (SS) | 0-3 b | Yes | 20.3 | \tilde{q}, \tilde{g} 1.45 TeV | $m(\tilde{q}) = m(\tilde{g}), c\tau_{LSP} < 1$ mm | 1404.2500 |
| | $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\mp}, \tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow e\bar{e}\nu, \mu\bar{\mu}\nu$ | $4 e, \mu$ | - | Yes | 13.3 | $\tilde{\chi}_1^{\pm}$ 1.14 TeV | $m(\tilde{\chi}_1^0) > 400$ GeV, $\lambda_{12k} \neq 0$ ($k = 1, 2$) | ATLAS-CONF-2016-075 |
| | $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\mp}, \tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\nu\tau, e\nu e, \tau\nu\tau$ | $3 e, \mu + \tau$ | - | Yes | 20.3 | $\tilde{\chi}_1^{\pm}$ 450 GeV | $m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^{\pm}), \lambda_{133} \neq 0$ | 1405.5086 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq$ | 0 | 4-5 large-R jets | - | 36.1 | \tilde{g} 1.875 TeV | $m(\tilde{\chi}_1^0) = 1075$ GeV | SUSY-2016-22 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq$ | $1 e, \mu$ | 8-10 jets/0-4 b | - | 36.1 | \tilde{g} 2.1 TeV | $m(\tilde{\chi}_1^0) = 1$ TeV, $\lambda_{112} \neq 0$ | 1704.08493 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$ | $1 e, \mu$ | 8-10 jets/0-4 b | - | 36.1 | \tilde{g} 1.65 TeV | $m(\tilde{t}_1) = 1$ TeV, $\lambda_{323} \neq 0$ | 1704.08493 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$ | 0 | 2 jets + 2 b | - | 36.7 | \tilde{t}_1 100-470 GeV 480-610 GeV | | 1710.07171 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\ell}$ | $2 e, \mu$ | 2 b | - | 36.1 | \tilde{t}_1 0.4-1.45 TeV | $BR(\tilde{t}_1 \rightarrow b\ell/\mu) > 20\%$ | 1710.05544 |
| Other | Scalar charm, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$ | 0 | 2 c | Yes | 20.3 | \tilde{c} 510 GeV | $m(\tilde{\chi}_1^0) < 200$ GeV | 1501.01325 |

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

10⁻¹ 1 Mass scale [TeV]

Large SUSY search program @ LHC

ATLAS SUSY Searches* - 95% CL Lower

December 2017

| Model | e, μ, τ, γ | Jets | E_T^{miss} | $\int \mathcal{L} dt [\text{fb}^{-1}]$ | | |
|---|---|--|---------------------|--|--------------------|--|
| Inclusive Searches | $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 36.1 | \tilde{q} |
| | $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (compressed) | mono-jet | 1-3 jets | Yes | 36.1 | \tilde{q} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^\pm \rightarrow qqW^\pm\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$ | $ee, \mu\mu$ | 2 jets | Yes | 14.7 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq(\ell\ell/\nu\nu)\tilde{\chi}_1^0$ | 3 e, μ | 4 jets | - | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$ | 0 | 7-11 jets | Yes | 36.1 | \tilde{g} |
| | GMSB ($\tilde{\ell}$ NLSP) | 1-2 τ + 0-1 ℓ | 0-2 jets | Yes | 3.2 | \tilde{g} |
| | GGM (bino NLSP) | 2 γ | - | Yes | 36.1 | \tilde{g} |
| | GGM (higgsino-bino NLSP) | γ | 2 jets | Yes | 36.1 | \tilde{g} |
| Gravitino LSP | 0 | mono-jet | Yes | 20.3 | $F^{1/2}$ | |
| 3 rd gen. \tilde{g} med. | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$ | 0 | 3 b | Yes | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$ | 0-1 e, μ | 3 b | Yes | 36.1 | \tilde{g} |
| 3 rd gen. squarks direct production | $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ | 0 | 2 b | Yes | 36.1 | \tilde{b}_1 |
| | $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^\pm$ | 2 e, μ (SS) | 1 b | Yes | 36.1 | \tilde{b}_1 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$ | 0-2 e, μ | 1-2 b | Yes | 4.7/13.3 | \tilde{t}_1 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$ | 0-2 e, μ | 0-2 jets/1-2 b | Yes | 20.3/36.1 | \tilde{t}_1 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$ | 0 | mono-jet | Yes | 36.1 | \tilde{t}_1 |
| | $\tilde{t}_1\tilde{t}_1$ (natural GMSB) | 2 e, μ (Z) | 1 b | Yes | 20.3 | \tilde{t}_1 |
| | $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$ | 3 e, μ (Z) | 1 b | Yes | 36.1 | \tilde{t}_2 |
| $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$ | 1-2 e, μ | 4 b | Yes | 36.1 | \tilde{t}_2 | |
| EW direct | $\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$ | 2 e, μ | 0 | Yes | 36.1 | $\tilde{\ell}$ |
| | $\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^\pm \rightarrow \tilde{\ell}\nu(\tilde{\ell}\bar{\nu})$ | 2 e, μ | 0 | Yes | 36.1 | $\tilde{\chi}_1^\pm$ |
| | $\tilde{\chi}_1^+\tilde{\chi}_1^0/\tilde{\chi}_2^0, \tilde{\chi}_1^\pm \rightarrow \tilde{\nu}\nu(\tilde{\nu}\bar{\nu}), \tilde{\chi}_2^0 \rightarrow \tilde{\tau}\tau(\tilde{\nu}\bar{\nu})$ | 2 τ | - | Yes | 36.1 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ |
| | $\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_L\nu\tilde{\ell}_L\ell(\tilde{\nu}\bar{\nu}), \tilde{\ell}\tilde{\nu}\tilde{\ell}_L\ell(\tilde{\nu}\bar{\nu})$ | 3 e, μ | 0 | Yes | 36.1 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ |
| | $\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 Z\tilde{\chi}_1^0$ | 2-3 e, μ | 0-2 jets | Yes | 36.1 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ |
| | $\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 h\tilde{\chi}_1^0, h \rightarrow b\tilde{b}/WW/\tau\tau/\gamma\gamma$ | e, μ, γ | 0-2 b | Yes | 20.3 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ |
| | $\tilde{\chi}_2^0\tilde{\chi}_3^0, \tilde{\chi}_2^0 \rightarrow \tilde{\ell}_R\ell$ | 4 e, μ | 0 | Yes | 20.3 | $\tilde{\chi}_2^0, \tilde{\chi}_3^0$ |
| | GGM (wino NLSP) weak prod., $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ | 1 $e, \mu + \gamma$ | - | Yes | 20.3 | \tilde{W} |
| | GGM (bino NLSP) weak prod., $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ | 2 γ | - | Yes | 36.1 | \tilde{W} |
| | Long-lived particles | Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$ | Disapp. trk | 1 jet | Yes | 36.1 |
| Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$ | | dE/dx trk | - | Yes | 18.4 | $\tilde{\chi}_1^\pm$ |
| Stable, stopped \tilde{g} R-hadron | | 0 | 1-5 jets | Yes | 27.9 | \tilde{g} |
| Stable \tilde{g} R-hadron | | trk | - | - | 3.2 | \tilde{g} |
| Metastable \tilde{g} R-hadron | | dE/dx trk | - | - | 3.2 | \tilde{g} |
| Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq\tilde{\chi}_1^0$ | | displ. vtx | - | Yes | 32.8 | \tilde{g} |
| GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$ | | 1-2 μ | - | - | 19.1 | $\tilde{\chi}_1^0$ |
| GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$ | 2 γ | - | Yes | 20.3 | $\tilde{\chi}_1^0$ | |
| $\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow ee/\mu\mu/\mu\nu$ | displ. $ee/\mu\mu/\mu\nu$ | - | - | 20.3 | $\tilde{\chi}_1^0$ | |
| RPV | LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\epsilon\tau/\mu\tau$ | $e\mu, \epsilon\tau, \mu\tau$ | - | - | 3.2 | $\tilde{\nu}_\tau, \tilde{q}, \tilde{g}$ |
| | Bilinear RPV CMSSM | 2 e, μ (SS) | 0-3 b | Yes | 20.3 | $\tilde{\chi}_1^\pm$ |
| | $\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow ee\nu, e\mu\nu, \mu\mu\nu$ | 4 e, μ | - | Yes | 13.3 | $\tilde{\chi}_1^\pm$ |
| | $\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\tau\nu_e, e\tau\nu_\tau$ | 3 $e, \mu + \tau$ | - | Yes | 20.3 | $\tilde{\chi}_1^\pm$ |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$ | 0 | 4-5 large- R jets | - | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$ | 1 e, μ | 8-10 jets/0-4 b | - | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$ | 1 e, μ | 8-10 jets/0-4 b | - | 36.1 | \tilde{g} |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$ | 0 | 2 jets + 2 b | - | 36.7 | \tilde{t}_1 | |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\ell$ | 2 e, μ | 2 b | - | 36.1 | \tilde{t}_1 | |
| Other | Scalar charm, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$ | 0 | 2 c | Yes | 20.3 | \tilde{c} |

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

10⁻¹

Searches for squarks and gluinos

jets + E_T^{miss}

cascade decays often yield leptons

jets + el/mu + E_T^{miss}

jets + tau + E_T^{miss}

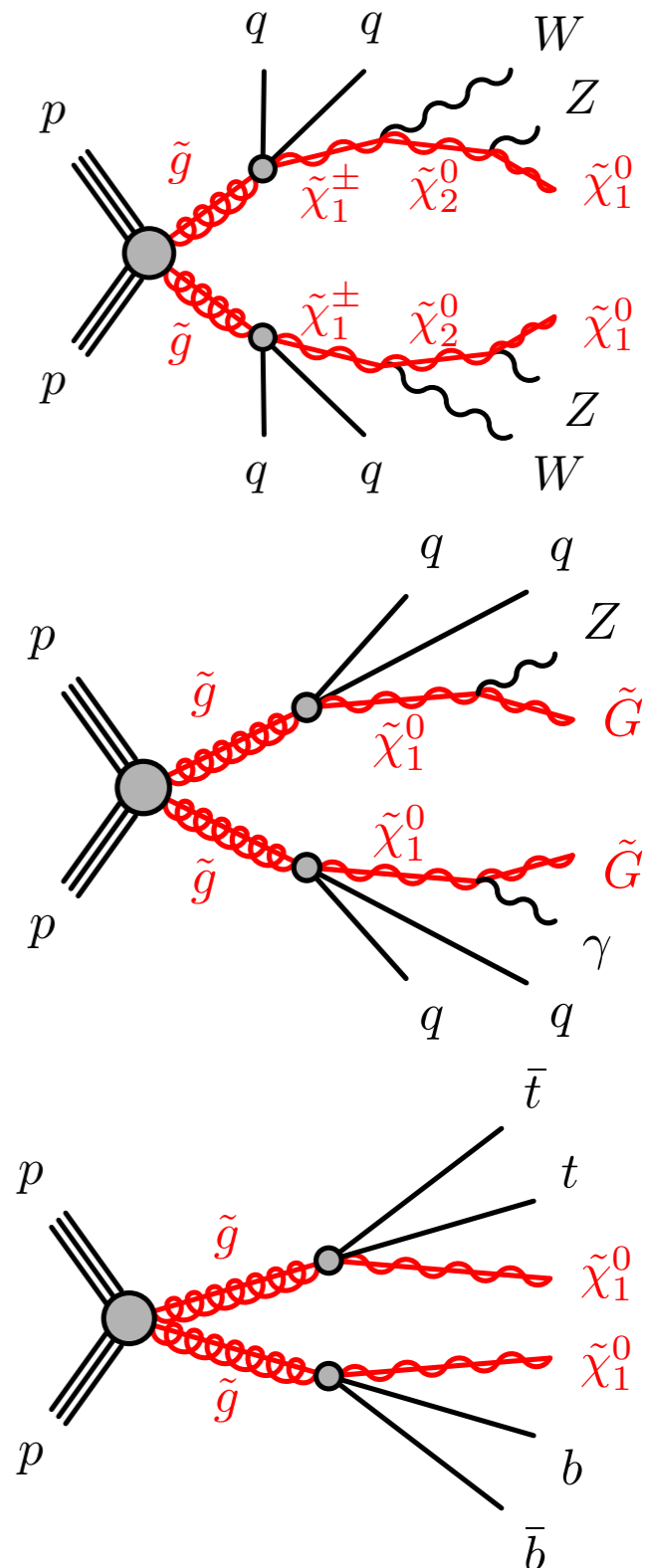
decays to Gravitinos can yield photons

jets + photon(s) + E_T^{miss}

decays can yield top & bottom quarks (natural SUSY)

jets + b-tags + (leptons) + E_T^{miss}

Example diagrams



Large SUSY search program @ LHC

ATLAS SUSY Searches* - 95% CL Lower Limits

December 2017

ATLAS Preliminary

| Model | e, μ, τ, γ | Jets | E_T^{miss} | $[\mathcal{L} dt(\text{fb}^{-1})]$ | | |
|---|---|--|---------------------|------------------------------------|---------------------|--|
| Inclusive Searches | $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 36.1 | \tilde{q} |
| | $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (compressed) | mono-jet | 1-3 jets | Yes | 36.1 | \tilde{q} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^\pm \rightarrow qqW^\pm\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$ | $ee, \mu\mu$ | 2 jets | Yes | 14.7 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq(\ell\ell/\nu\nu)\tilde{\chi}_1^0$ | 3 e, μ | 4 jets | - | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$ | 0 | 7-11 jets | Yes | 36.1 | \tilde{g} |
| | GMSB ($\tilde{\ell}$ NLSP) | 1-2 τ + 0-1 ℓ | 0-2 jets | Yes | 3.2 | \tilde{g} |
| | GGM (bino NLSP) | 2 γ | - | Yes | 36.1 | \tilde{g} |
| | GGM (higgsino-bino NLSP) | γ | 2 jets | Yes | 36.1 | \tilde{g} |
| Gravitino LSP | 0 | mono-jet | Yes | 20.3 | $F^{1/2} \tilde{g}$ | |
| 2 nd gen. \tilde{g} med. | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$ | 0 | 3 b | Yes | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$ | 0-1 e, μ | 3 b | Yes | 36.1 | \tilde{g} |
| 3 rd gen. squarks direct production | $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ | 0 | 2 b | Yes | 36.1 | \tilde{b}_1 |
| | $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^\pm$ | 2 e, μ (SS) | 1 b | Yes | 36.1 | \tilde{b}_1 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$ | 0-2 e, μ | 1-2 b | Yes | 4.7/13.3 | \tilde{t}_1 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$ | 0-2 e, μ | 0-2 jets/1-2 b | Yes | 20.3/36.1 | \tilde{t}_1 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$ | 0 | mono-jet | Yes | 36.1 | \tilde{t}_1 |
| | $\tilde{t}_1\tilde{t}_1$ (natural GMSB) | 2 e, μ (Z) | 1 b | Yes | 20.3 | \tilde{t}_1 |
| | $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$ | 3 e, μ (Z) | 1 b | Yes | 36.1 | \tilde{t}_2 |
| | $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$ | 1-2 e, μ | 4 b | Yes | 36.1 | \tilde{t}_2 |
| EW direct | $\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$ | 2 e, μ | 0 | Yes | 36.1 | $\tilde{\ell}$ |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp, \tilde{\chi}_1^\pm \rightarrow \ell\nu(\ell\bar{\nu})$ | 2 e, μ | 0 | Yes | 36.1 | $\tilde{\chi}_1^\pm$ |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0, \tilde{\chi}_1^\pm \rightarrow \tau\nu(\tau\bar{\nu}), \tilde{\chi}_2^0 \rightarrow \tau\tau(\nu\bar{\nu})$ | 2 τ | - | Yes | 36.1 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_L\nu\tilde{\ell}_L(\ell\nu), \tilde{\ell}\tilde{\nu}\tilde{\ell}_L(\ell\nu)$ | 3 e, μ | 0 | Yes | 36.1 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0Z\tilde{\chi}_1^0$ | 2-3 e, μ | 0-2 jets | Yes | 36.1 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0h\tilde{\chi}_1^0, h \rightarrow b\tilde{b}/WW/\tau\tau/\gamma\gamma$ | e, μ, γ | 0-2 b | Yes | 20.3 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ |
| | $\tilde{\chi}_2^0\tilde{\chi}_3^0, \tilde{\chi}_2^0 \rightarrow \tilde{\ell}_R\ell$ | 4 e, μ | 0 | Yes | 20.3 | $\tilde{\chi}_2^0, \tilde{\chi}_3^0$ |
| | GGM (wino NLSP) weak prod., $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ | 1 $e, \mu + \gamma$ | - | Yes | 20.3 | \tilde{W} |
| | GGM (bino NLSP) weak prod., $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ | 2 γ | - | Yes | 36.1 | \tilde{W} |
| | Long-lived particles | Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$ | Disapp. trk | 1 jet | Yes | 36.1 |
| Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$ | | dE/dx trk | - | Yes | 18.4 | $\tilde{\chi}_1^\pm$ |
| Stable, stopped \tilde{g} R-hadron | | 0 | 1-5 jets | Yes | 27.9 | \tilde{g} |
| Stable \tilde{g} R-hadron | | trk | - | - | 3.2 | \tilde{g} |
| Metastable \tilde{g} R-hadron | | dE/dx trk | - | - | 3.2 | \tilde{g} |
| Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq\tilde{\chi}_1^0$ | | displ. vtx | - | Yes | 32.8 | \tilde{g} |
| GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tau(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$ | | 1-2 μ | - | - | 19.1 | $\tilde{\chi}_1^0$ |
| GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$ | | 2 γ | - | Yes | 20.3 | $\tilde{\chi}_1^0$ |
| $\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow e\bar{\nu}/e\nu/\mu\bar{\nu}/\mu\nu$ | displ. $ee/e\mu/\mu\mu$ | - | - | 20.3 | $\tilde{\chi}_1^0$ | |
| RPV | LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\tau\mu$ | $e\mu, \tau\mu$ | - | - | 3.2 | $\tilde{\nu}_\tau$ |
| | Bilinear RPV CMSSM | 2 e, μ (SS) | 0-3 b | Yes | 20.3 | \tilde{q}, \tilde{g} |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow e\bar{\nu}, e\nu, \mu\bar{\nu}$ | 4 e, μ | - | Yes | 13.3 | $\tilde{\chi}_1^\pm$ |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\nu_e, e\nu_\tau$ | 3 $e, \mu + \tau$ | - | Yes | 20.3 | $\tilde{\chi}_1^\pm$ |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq$ | 0 | 4-5 large- R jets | - | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq$ | 1 e, μ | 8-10 jets/0-4 b | - | 36.1 | \tilde{g} |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$ | 1 e, μ | 8-10 jets/0-4 b | - | 36.1 | \tilde{g} |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$ | 0 | 2 jets + 2 b | - | 36.7 | \tilde{t}_1 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\ell$ | 2 e, μ | 2 b | - | 36.1 | \tilde{t}_1 | |
| Other | Scalar charm, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$ | 0 | 2 c | Yes | 20.3 | \tilde{c} |

Dedicated search program for stop & sbottom production

top quark pair + E_T^{miss}

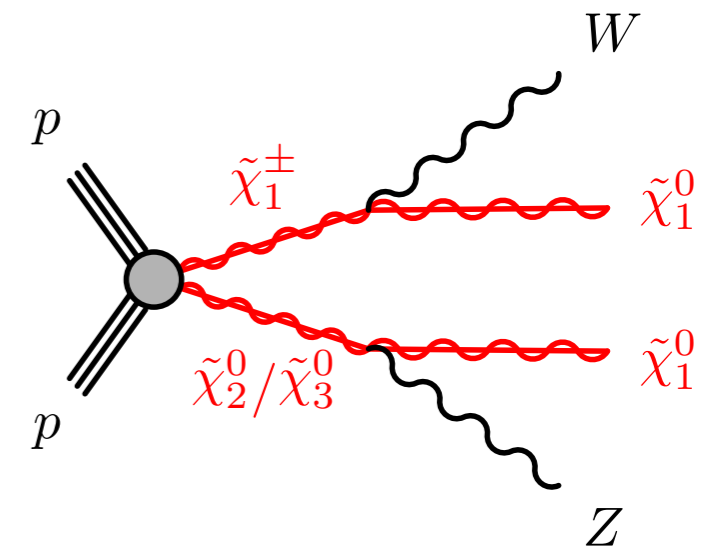
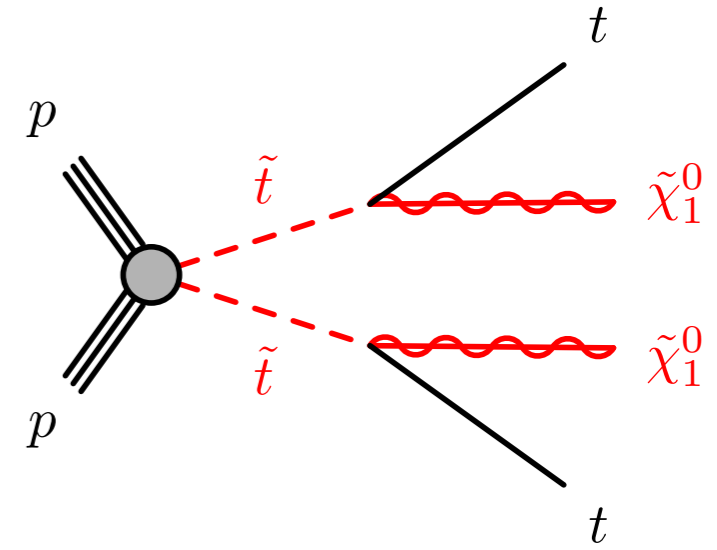
and more complex signatures

Dedicated search program for electroweak production: charginos, neutralinos, sleptons

di-boson + E_T^{miss}

and more complex signatures

Example diagrams



*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

10⁻¹

Large SUSY search program @ LHC

ATLAS SUSY Searches* - 95

December 2017

| Model | e, μ, τ, γ | Jets |
|---|-------------------------|-------------------|
| Inclusive Searches | | |
| $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ | 0 | 2-6 jets |
| $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (compressed) | mono-jet | 1-3 jets |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{\chi}_1^0$ | 0 | 2-6 jets |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^\pm \rightarrow qqW^\pm\tilde{\chi}_1^0$ | 0 | 2-6 jets |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$ | $ee, \mu\mu$ | 2 jets |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq(\ell\ell/\nu\nu)\tilde{\chi}_1^0$ | 3 e, μ | 4 jets |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$ | 0 | 7-11 jets |
| GMSB ($\tilde{\ell}$ NLSP) | 1-2 τ + 0-1 ℓ | 0-2 jets |
| GGM (bino NLSP) | 2 γ | - |
| GGM (higgsino-bino NLSP) | γ | 2 jets |
| Gravitino LSP | 0 | mono-jet |
| 3rd gen. \tilde{g} med. | | |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$ | 0 | 3 jets |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$ | 0-1 e, μ | 3 jets |
| 3rd gen. squarks direct production | | |
| $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ | 0 | 2 jets |
| $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^\pm$ | 2 e, μ (SS) | 1 jet |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$ | 0-2 e, μ | 1-2 jets |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$ | 0-2 e, μ | 0-2 jets/mono-jet |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$ | 0 | mono-jet |
| $\tilde{t}_1\tilde{t}_1$ (natural GMSB) | 2 e, μ (Z) | 1 jet |
| $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$ | 3 e, μ (Z) | 1 jet |
| $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$ | 1-2 e, μ | 4 jets |
| EW direct | | |
| $\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$ | 2 e, μ | 0 |
| $\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow \ell\nu(\ell\bar{\nu})$ | 2 e, μ | 0 |
| $\tilde{\chi}_1^+\tilde{\chi}_1^0/\tilde{\chi}_2^0, \tilde{\chi}_1^+ \rightarrow \tau\nu(\tau\bar{\nu}), \tilde{\chi}_2^0 \rightarrow \tau\tau(\nu\bar{\nu})$ | 2 τ | - |
| $\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_L\nu\tilde{\ell}_L(\ell\bar{\nu}), \tilde{\ell}\tilde{\nu}\tilde{\ell}_L(\ell\bar{\nu})$ | 3 e, μ | 0 |
| $\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0Z\tilde{\chi}_1^0$ | 2-3 e, μ | 0-2 jets |
| $\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0h\tilde{\chi}_1^0, h \rightarrow b\bar{b}/WW/\tau\tau/\gamma\gamma$ | e, μ, γ | 0-2 jets |
| $\tilde{\chi}_2^0\tilde{\chi}_3^0, \tilde{\chi}_2^0 \rightarrow \tilde{\ell}_R\ell$ | 4 e, μ | 0 |
| GGM (wino NLSP) weak prod., $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ | 1 e, μ + γ | - |
| GGM (bino NLSP) weak prod., $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ | 2 γ | - |
| Long-lived particles | | |
| Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$ | Disapp. trk | 1 jet |
| Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$ | dE/dx trk | - |
| Stable, stopped \tilde{g} R-hadron | 0 | 1-5 jets |
| Stable \tilde{g} R-hadron | trk | - |
| Metastable \tilde{g} R-hadron | dE/dx trk | - |
| Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq\tilde{\chi}_1^0$ | displ. vtx | - |
| GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$ | 1-2 μ | - |
| GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$ | 2 γ | - |
| $\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow ee\nu/e\mu/\mu\mu\nu$ | displ. $ee/e\mu/\mu\mu$ | - |
| RPV | | |
| LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/e\tau/\mu\tau$ | $e\mu, e\tau, \mu\tau$ | - |
| Bilinear RPV CMSSM | 2 e, μ (SS) | 0-3 jets |
| $\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow ee\nu, e\mu\nu, \mu\mu\nu$ | 4 e, μ | - |
| $\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\nu e, e\tau\nu$ | 3 $e, \mu + \tau$ | - |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$ | 0 | 4-5 large-jets |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$ | 1 e, μ | 8-10 jets |
| $\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$ | 1 e, μ | 8-10 jets |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$ | 0 | 2 jets + |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\ell$ | 2 e, μ | 2 jets |
| Other | | |
| Scalar charm, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$ | 0 | 2 jets |

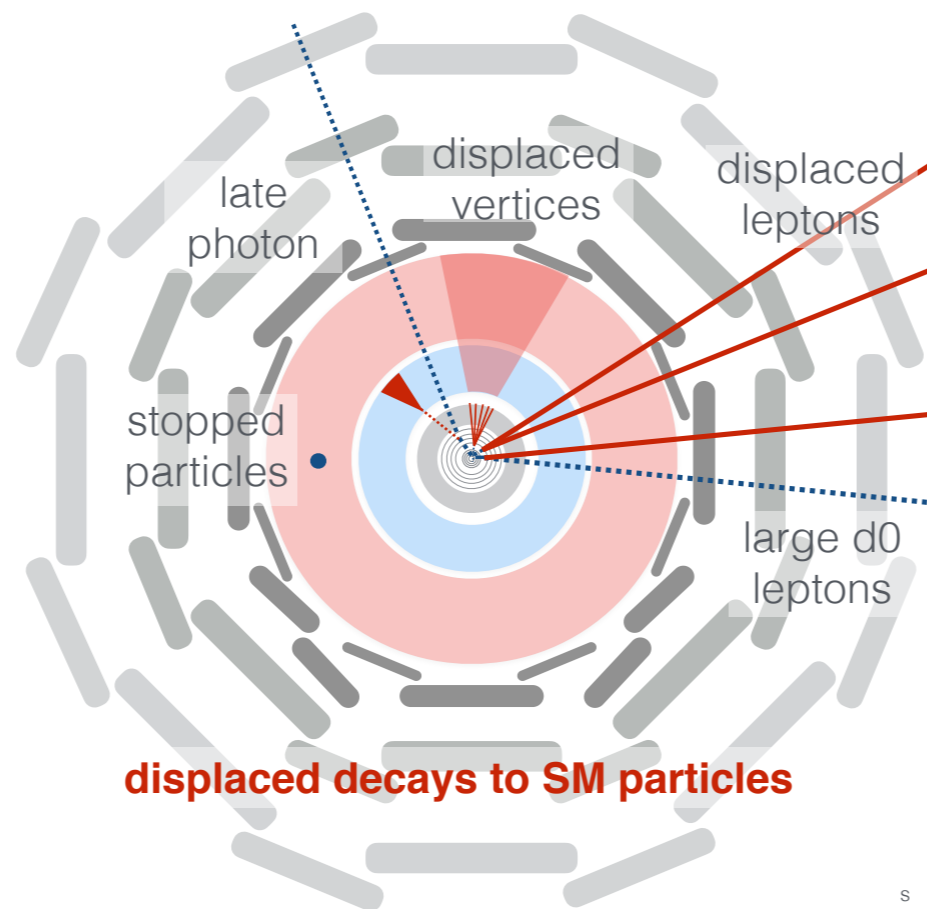
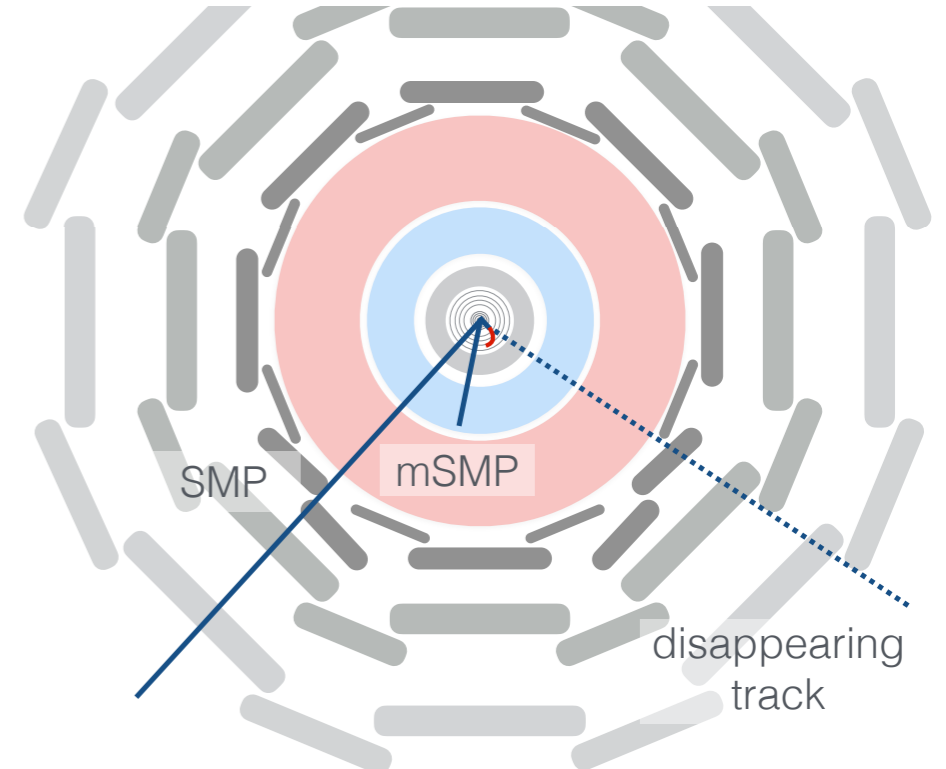
*Only a selection of the available mass limits on new states and phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

Long-lived particles

SUSY mechanisms:

- small couplings,
- off-shell decays,
- phase-space suppression

new particles interact with detector



displaced decays to SM particles

Sketches by Michael Adersberger

Large SUSY search program @ LHC

ATLAS SUSY
December 2017
Model

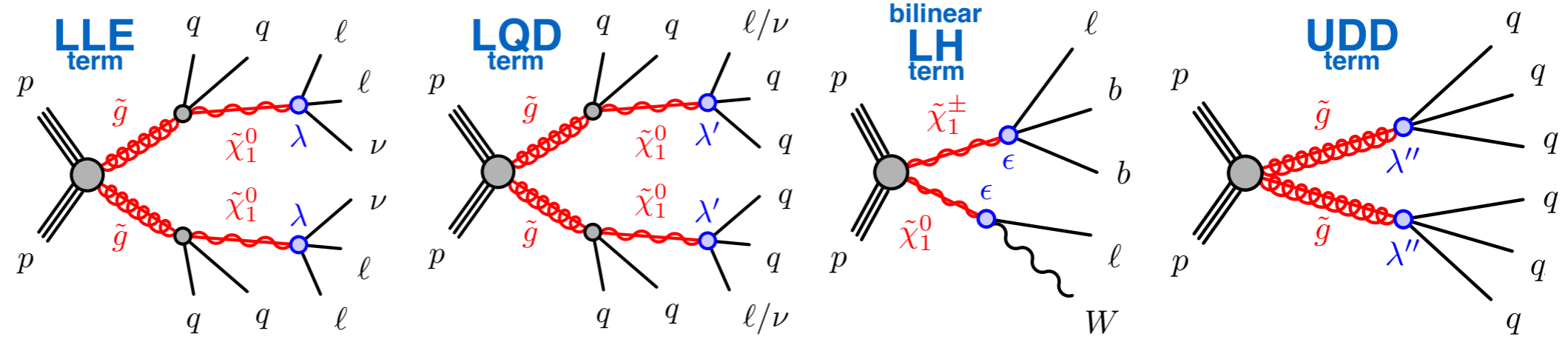
| | |
|--|---|
| Inclusive Searches | $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (com) $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^{\pm} \rightarrow qq\tilde{\chi}_1^0$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq(\ell\ell/\nu\nu)$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ$ GMSB ($\tilde{\ell}$ NLSP) GGM (bino NLSP) GGM (higgsino NLSP) Gravitino LSP |
| 3 rd gen. \tilde{g} med. | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$ |
| 3 rd gen. squarks direct production | $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^{\pm}$ $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\tilde{\chi}_1^{\pm}$ $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$ $\tilde{t}_1\tilde{t}_1$ (natural GMSB) $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$ $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$ |
| EW direct | $\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$ $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\pm}, \tilde{\chi}_1^{\pm} \rightarrow \ell\nu(\ell\bar{\nu})$ $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\pm}/\tilde{\chi}_2^0, \tilde{\chi}_1^{\pm} \rightarrow \tau\nu$ $\tilde{\chi}_1^{\pm}\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_L\nu\tilde{\ell}_L(\nu\bar{\nu})$ $\tilde{\chi}_1^{\pm}\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0Z\tilde{\chi}_1^0$ $\tilde{\chi}_1^{\pm}\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0h\tilde{\chi}_1^0$ $\tilde{\chi}_1^{\pm}\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0\tilde{\chi}_1^0$ $\tilde{\chi}_1^{\pm}\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_R\nu\tilde{\ell}_R$ GGM (wino NLSP) GGM (bino NLSP) |
| Long-lived particles | Direct $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\pm}$ prod Direct $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\pm}$ prod Stable, stopped \tilde{g} Stable \tilde{g} R-hadron Metastable \tilde{g} R-hadron Metastable \tilde{g} R-hadron GMSB, stable $\tilde{\tau}$ GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ $\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow e\bar{\nu}/e\nu/\mu\bar{\nu}/\mu\nu$ |

R-parity violation scenarios

Tree-level terms in SUSY that violate R-parity

$$W_{RPV} = \frac{1}{2} \lambda_{ijk} L_i L_j E_k^c + \lambda'_{ijk} L_i Q_j D_k^c + \kappa_i L_i H_u + \frac{1}{2} \lambda''_{ijk} U_i^c D_j^c D_k^c$$

trilinear coupling
bilinear coupling
trilinear coupling



Protect proton decay by e.g. not violating B and L simultaneously.

unstable LSP \rightarrow no obvious SUSY dark matter candidate, detector signature w/o E_T^{miss}

| Model | Signature | Searches | Yes | 20.3 | Mass scale [TeV] | Assumptions | Reference |
|-------|---|------------------------|---------------------|------|------------------------|--|---------------------|
| RPV | LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/e\tau/\mu\tau$ | $e\mu, e\tau, \mu\tau$ | - | 3.2 | $\tilde{\nu}_\tau$ | $\lambda'_{511}=0.11, \lambda_{132/133/233}=0.07$ | 1607.08079 |
| | Bilinear RPV CMSSM | $2 e, \mu$ (SS) | 0-3 b | Yes | \tilde{q}, \tilde{g} | $m(\tilde{q})=m(\tilde{g}), c\tau_{LSP} < 1$ mm | 1404.2500 |
| | $\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow e\bar{\nu}, e\nu, \mu\bar{\nu}, \mu\nu$ | $4 e, \mu$ | - | Yes | $\tilde{\chi}_1^{\pm}$ | $m(\tilde{\chi}_1^0) > 400$ GeV, $\lambda_{12k} \neq 0$ ($k=1,2$) | ATLAS-CONF-2016-075 |
| | $\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\bar{\nu}_e, e\nu_\tau$ | $3 e, \mu + \tau$ | - | Yes | $\tilde{\chi}_1^{\pm}$ | $m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^{\pm}), \lambda_{133} \neq 0$ | 1405.5086 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq$ | 0 | 4-5 large- R jets | - | \tilde{g} | $m(\tilde{\chi}_1^0) = 1075$ GeV | SUSY-2016-22 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq$ | $1 e, \mu$ | 8-10 jets/0-4 b | - | \tilde{g} | $m(\tilde{\chi}_1^0) = 1$ TeV, $\lambda_{112} \neq 0$ | 1704.08493 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$ | $1 e, \mu$ | 8-10 jets/0-4 b | - | \tilde{g} | $m(\tilde{t}_1) = 1$ TeV, $\lambda_{323} \neq 0$ | 1704.08493 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$ | 0 | 2 jets + 2 b | - | \tilde{t}_1 | | 1710.07171 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\ell$ | $2 e, \mu$ | 2 b | - | \tilde{t}_1 | $BR(\tilde{t}_1 \rightarrow b\ell/\mu) > 20\%$ | 1710.05544 |
| Other | Scalar charm, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$ | 0 | 2 c | Yes | \tilde{c} | $m(\tilde{\chi}_1^0) < 200$ GeV | 1501.01325 |

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

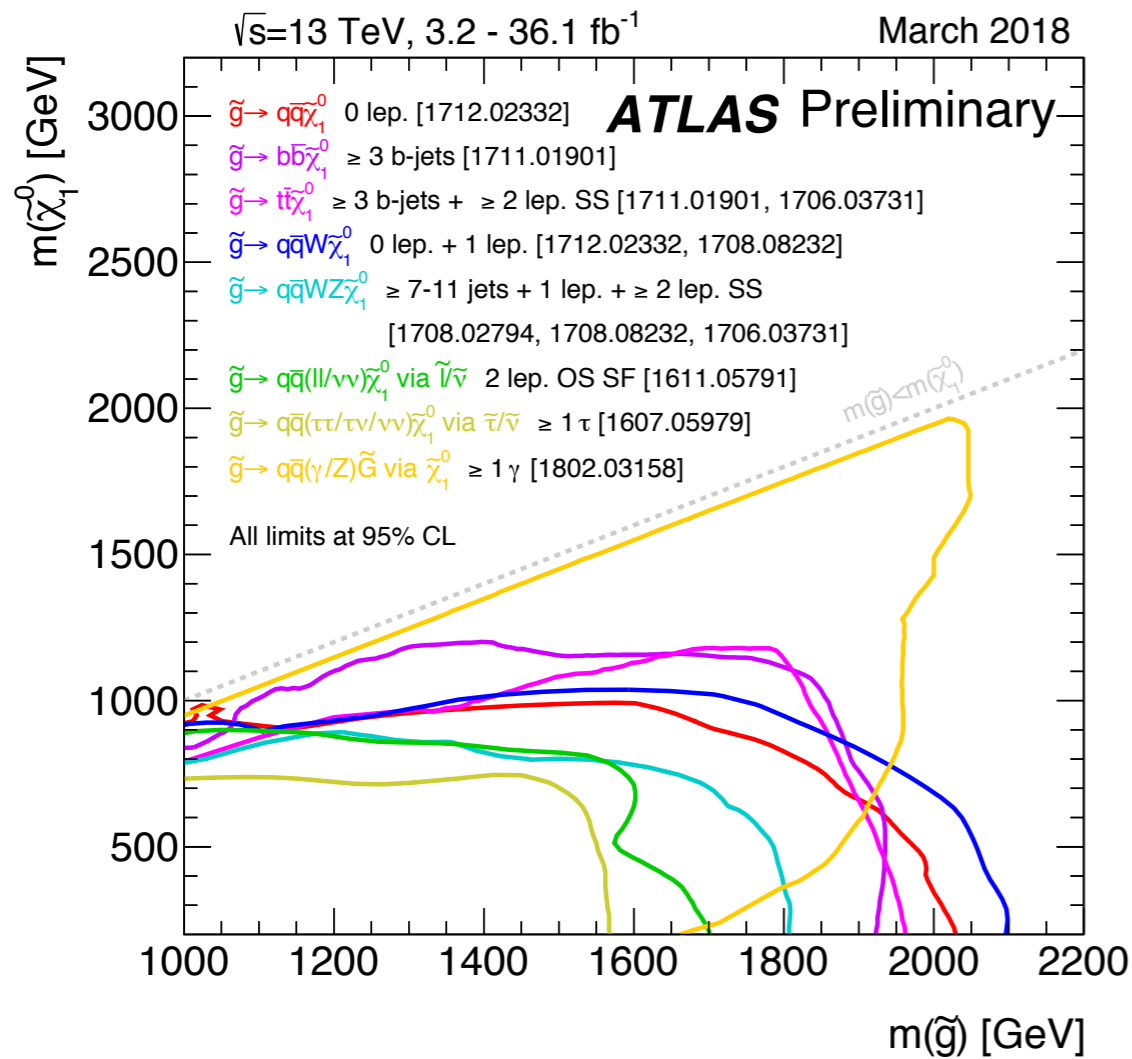


Frontiers

Complex final states and searches for rare & challenging signatures

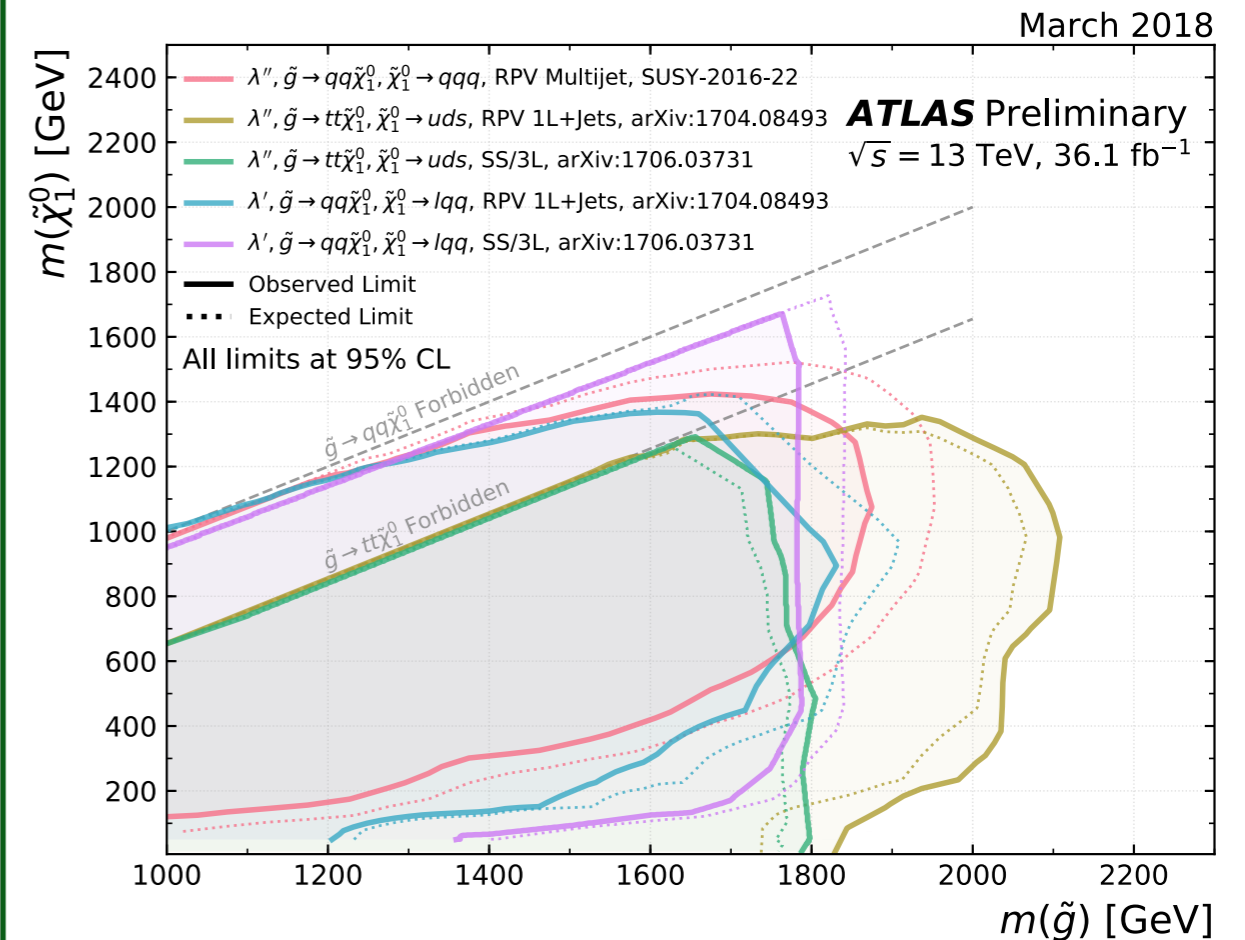
From R -parity conserving to R -parity violating

Exclusion summary for gluino searches in **RPC SUSY** (with E_T^{miss}) scenarios



target prompt decays to jets + X + E_T^{miss}

Exclusion summary for gluino searches in **RPV SUSY** (without E_T^{miss}) scenarios

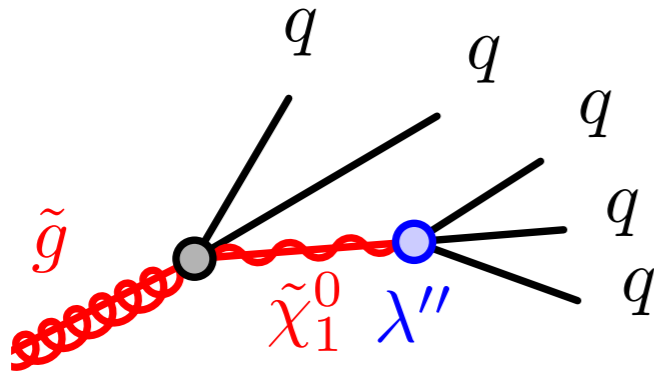


target prompt decays to jets or leptons

From R -parity conserving to R -parity violating

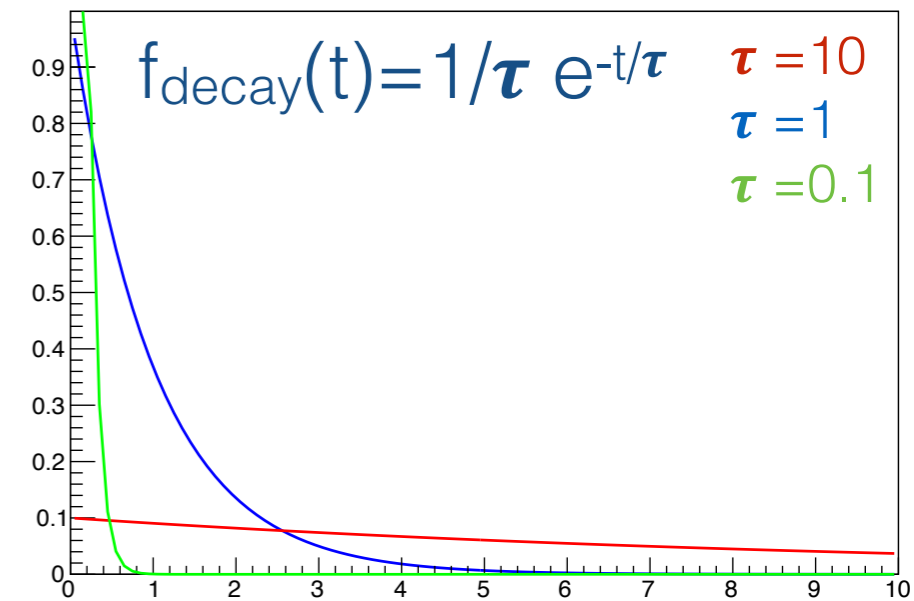
mean decay length for a bino-like lightest neutralino is approximately [Nucl. Phys. B 365 (1991) 597]:

$$L(\text{cm}) = \frac{0.9\beta\gamma}{\lambda''^2} \left(\frac{m(\tilde{q})}{100 \text{ GeV}} \right)^4 \left(\frac{1 \text{ GeV}}{m(\tilde{\chi}_1^0)} \right)^5 \quad \langle L \rangle = \beta\tau = \beta\gamma\tau^0$$



χ_1^0 RPV decay can be non-prompt.

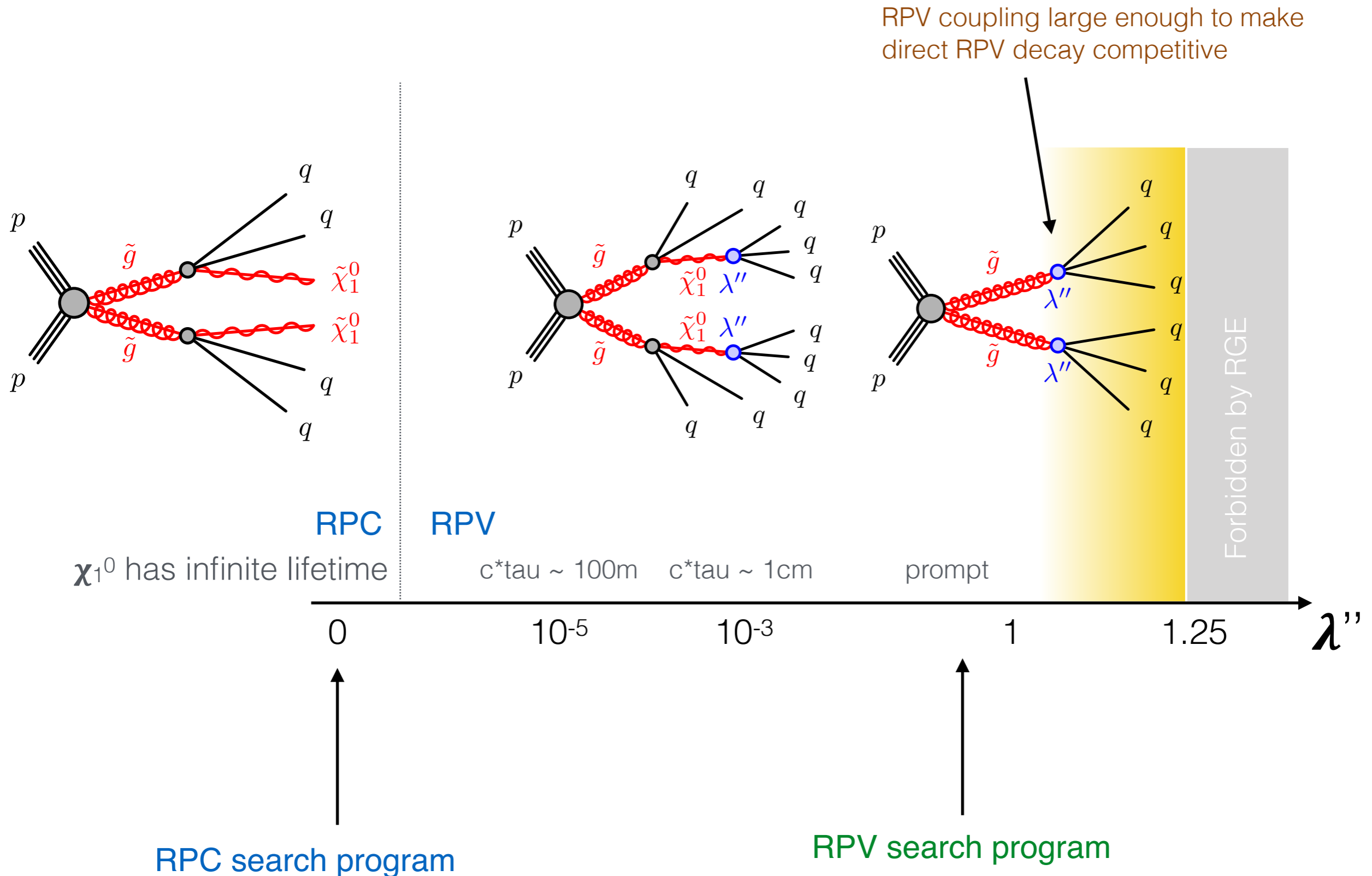
No genuine E_T^{miss}



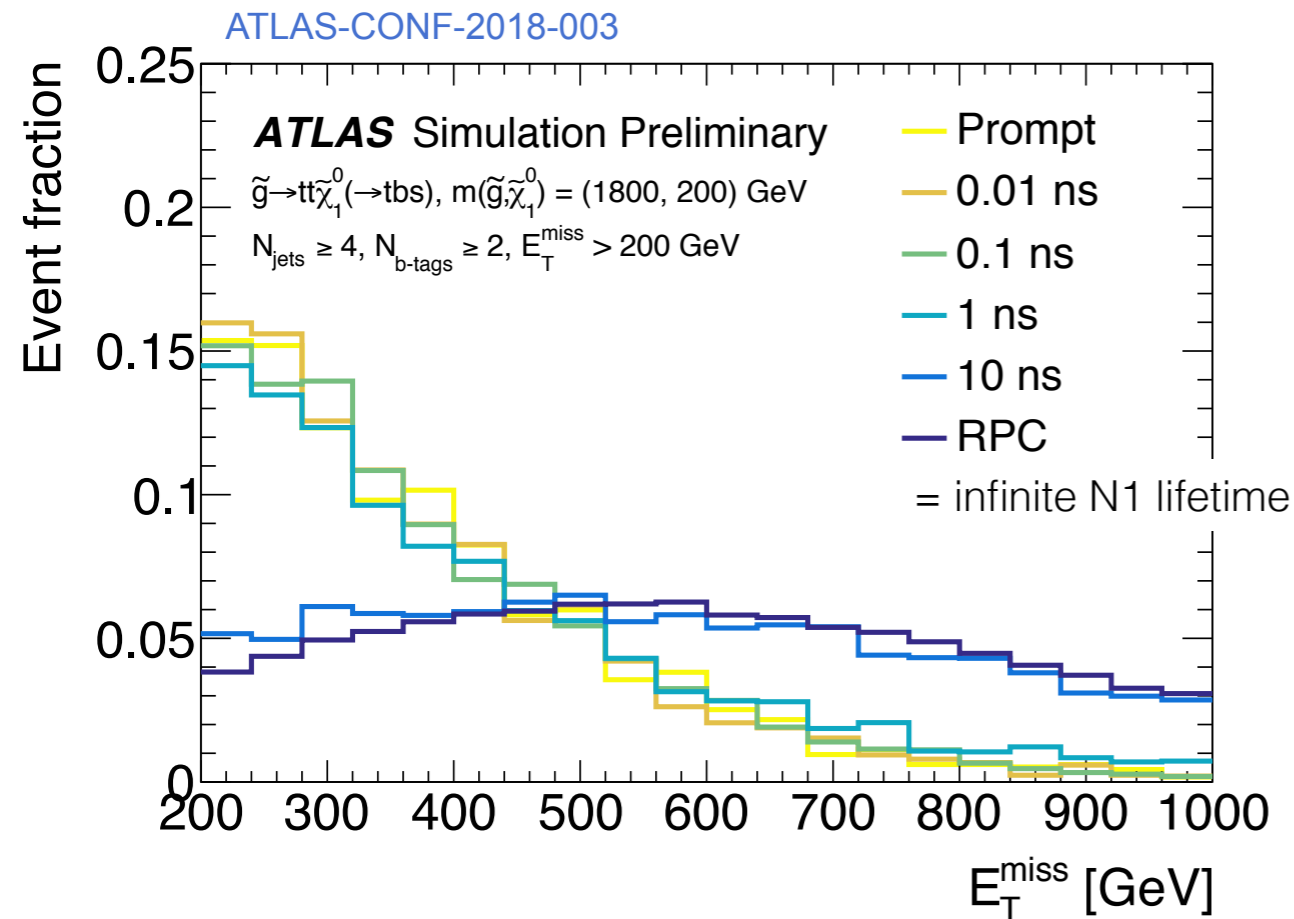
Long-lived particles search program

- typically requires large E_T^{miss} for online event selection (trigger)
- sensitivity above minimal displacement (analysis dependent)

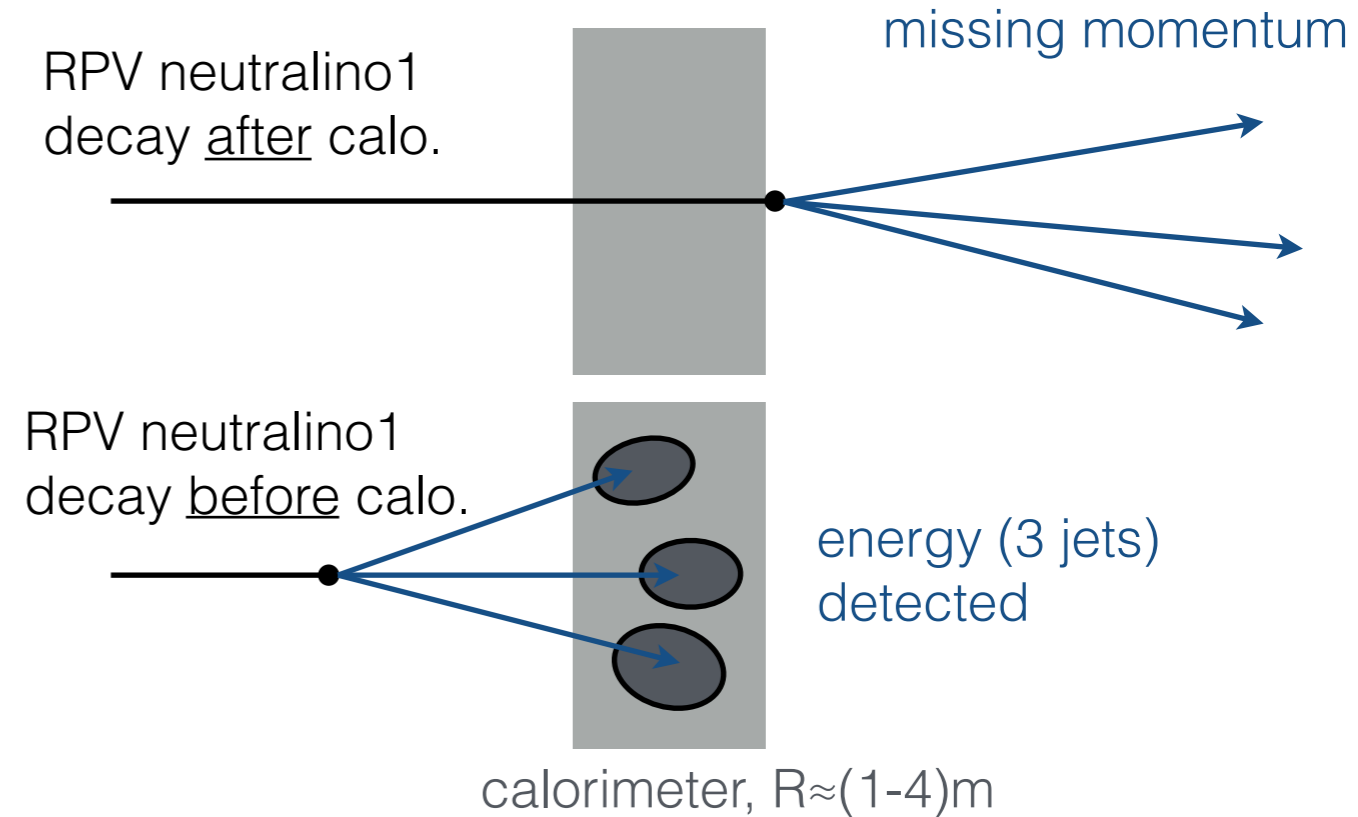
From R -parity conserving to R -parity violating



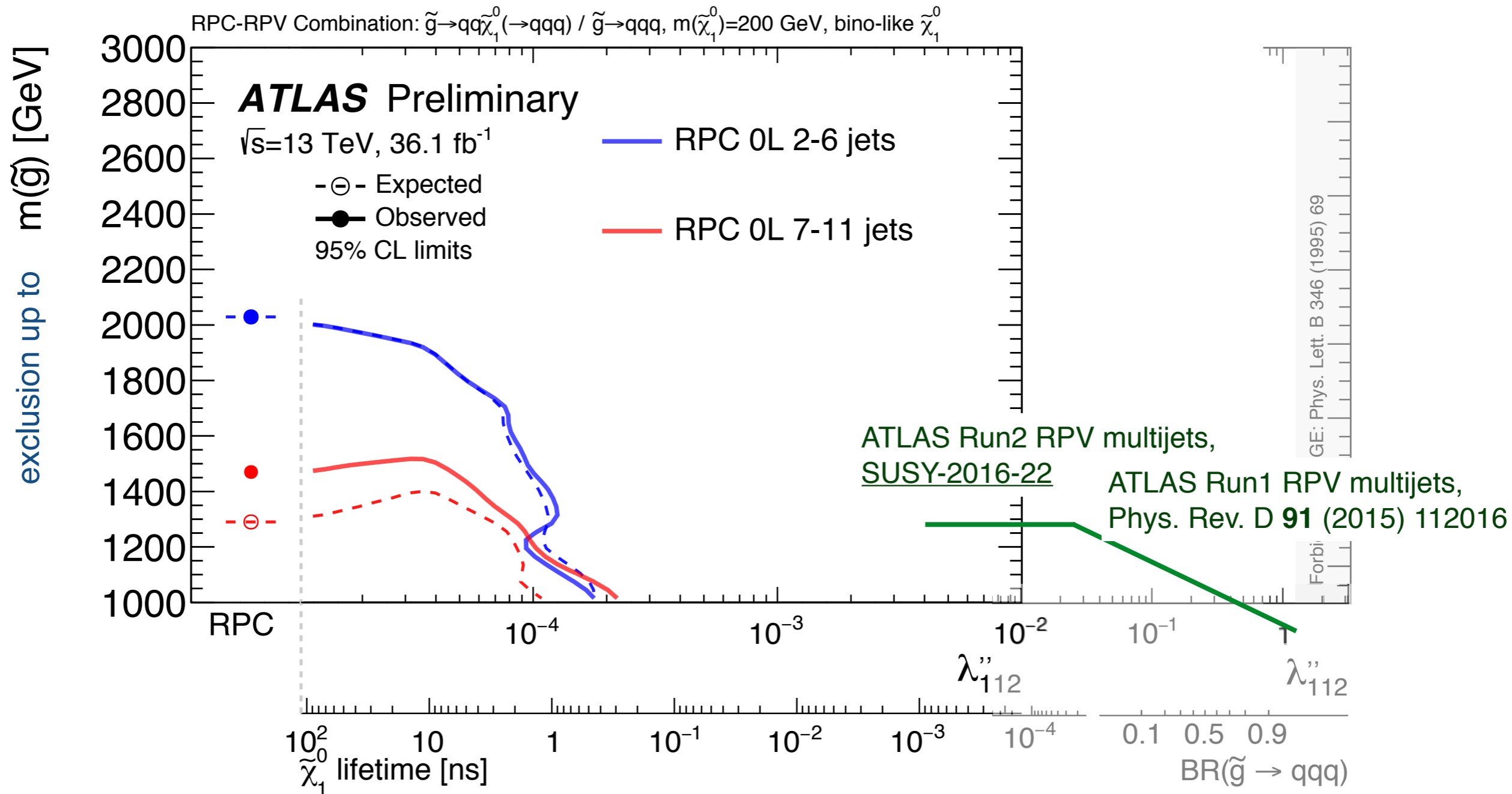
Characteristic observables



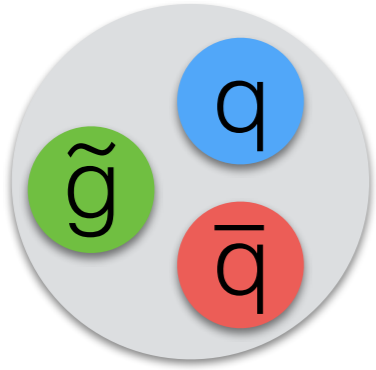
Reconstruct jets once neutralino decays are before calorimeter.



From R -parity conserving to R -parity

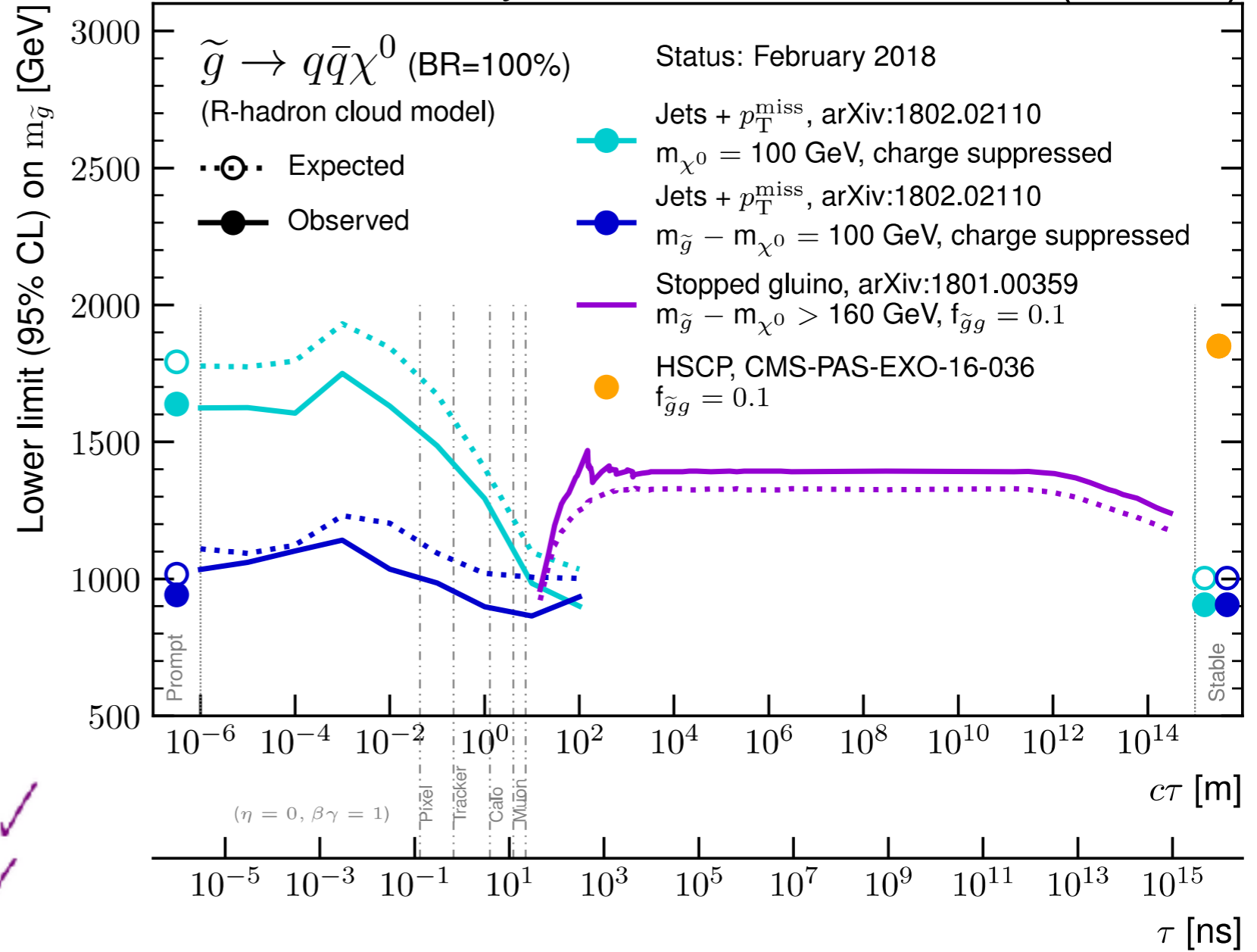


R-hadron scenario

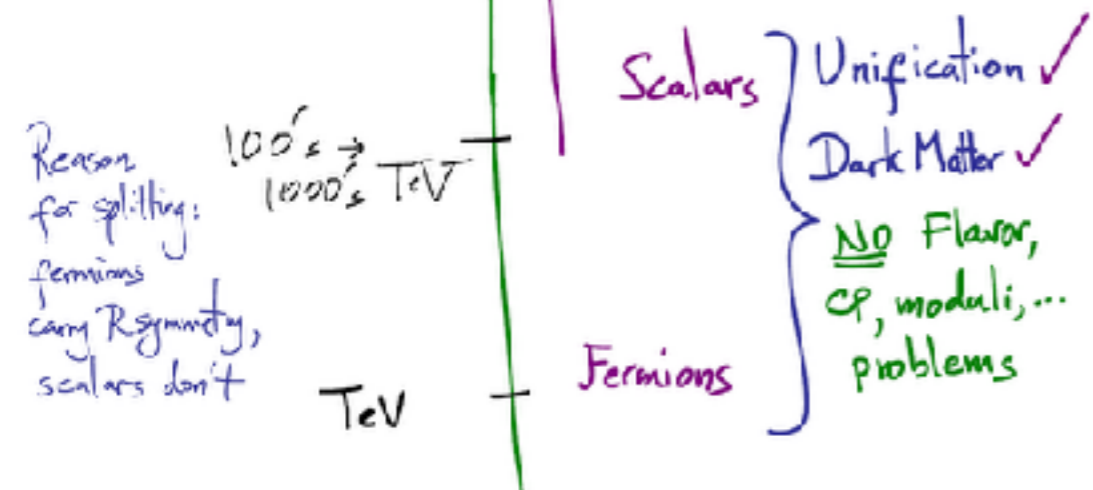


CMS Preliminary

13–39 fb^{-1} (13 TeV)

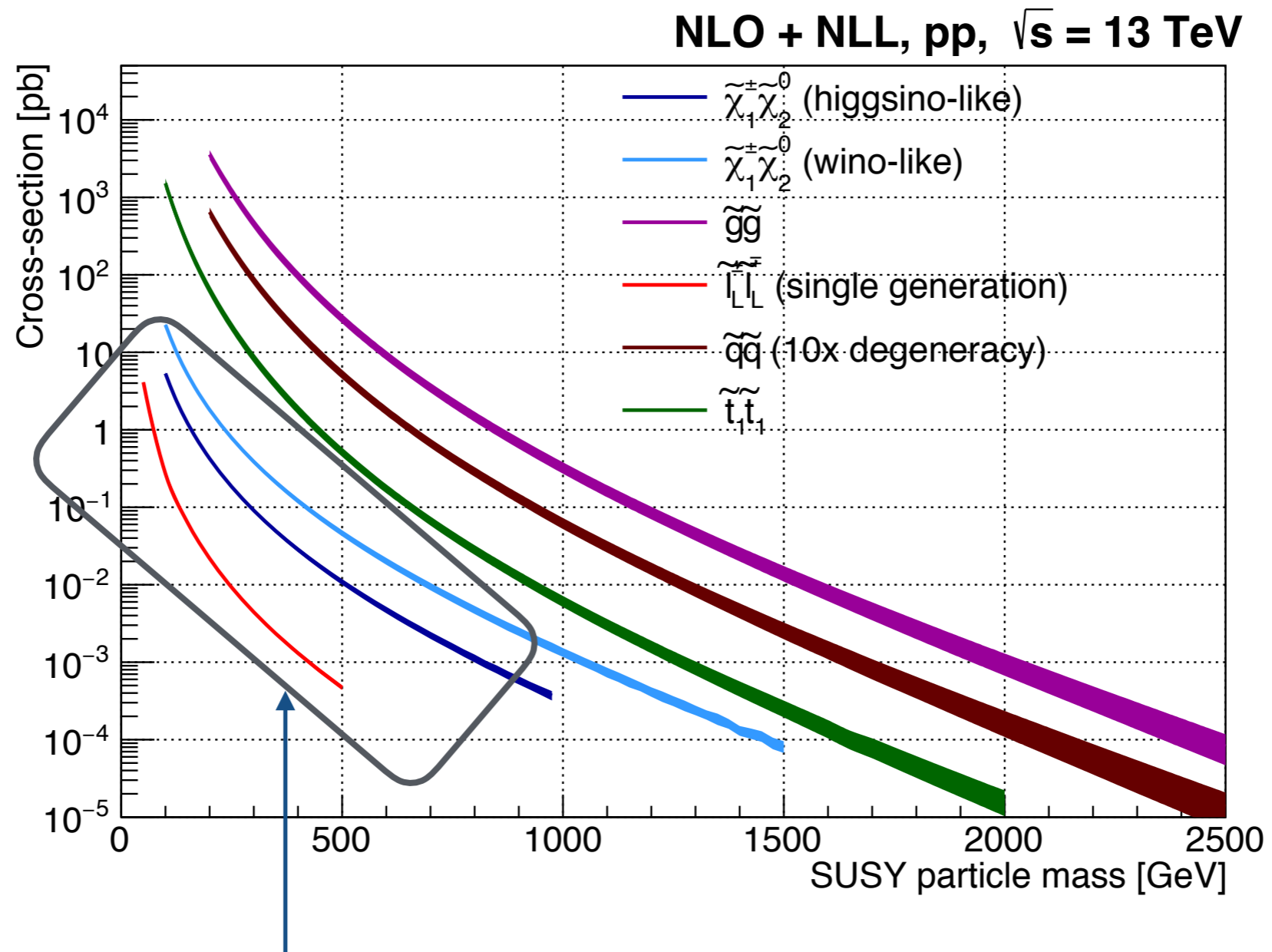


SPLIT SUSY



From Arkani-Hamed's talk

Frontier: Low rates



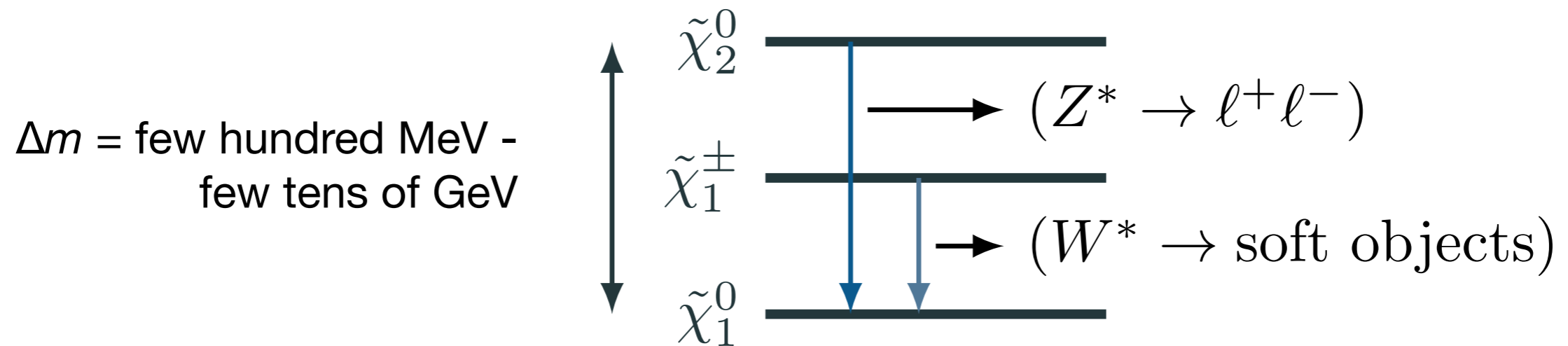
Production cross section of SUSY particles
w/o color charge, **electroweak production**

Compressed electroweak spectrum

Dark matter: co-annihilation is compression

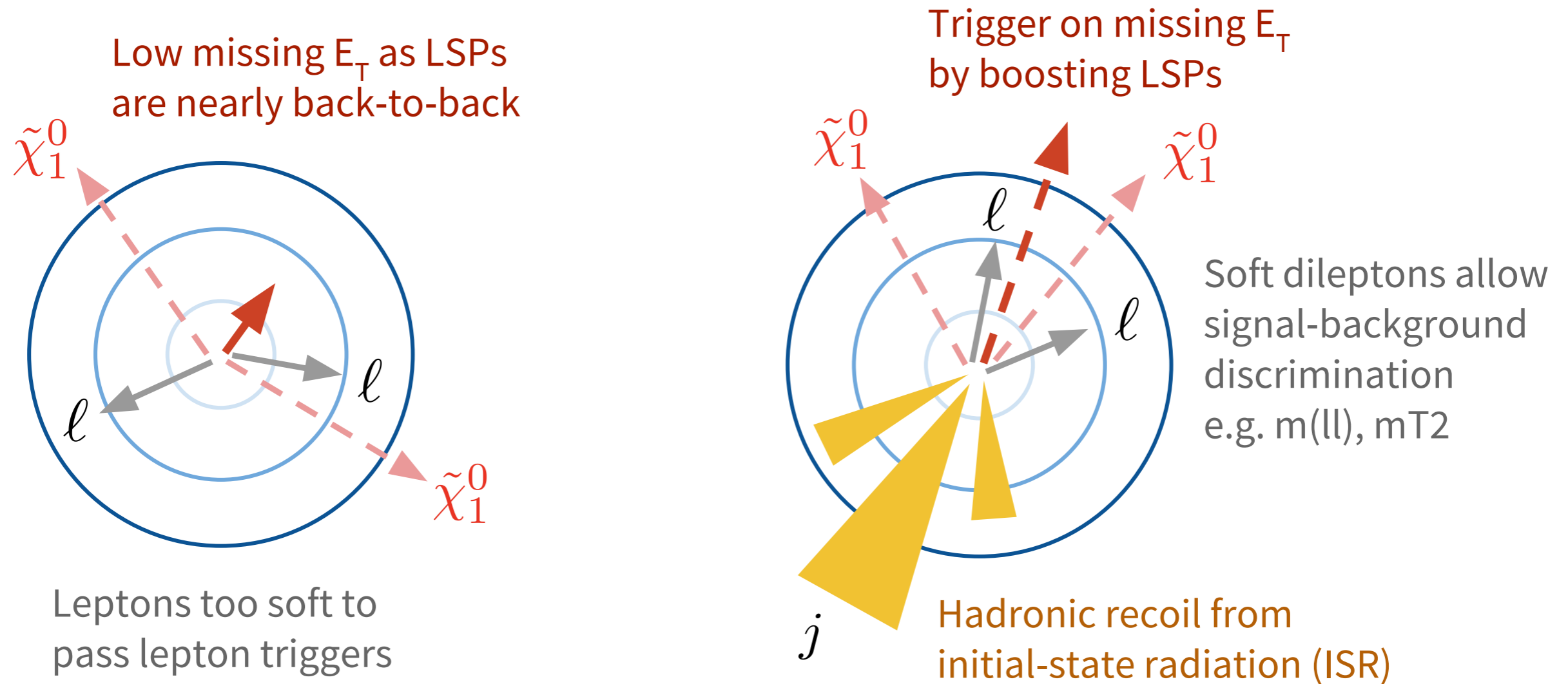
Naturalness requires Higgsinos near weak-scale

Higgsinos realised as multiplet of neutralinos & charginos



Challenge to reconstruct intra-Higgsino soft decay products

Strategy: initial-state-radiation (ISR)



Guidice et al [1004.4902], Gori et al [1307.5952], Han et al [1401.1235], Baer et al [1409.7058], Barr et al [1501.02511]. . .

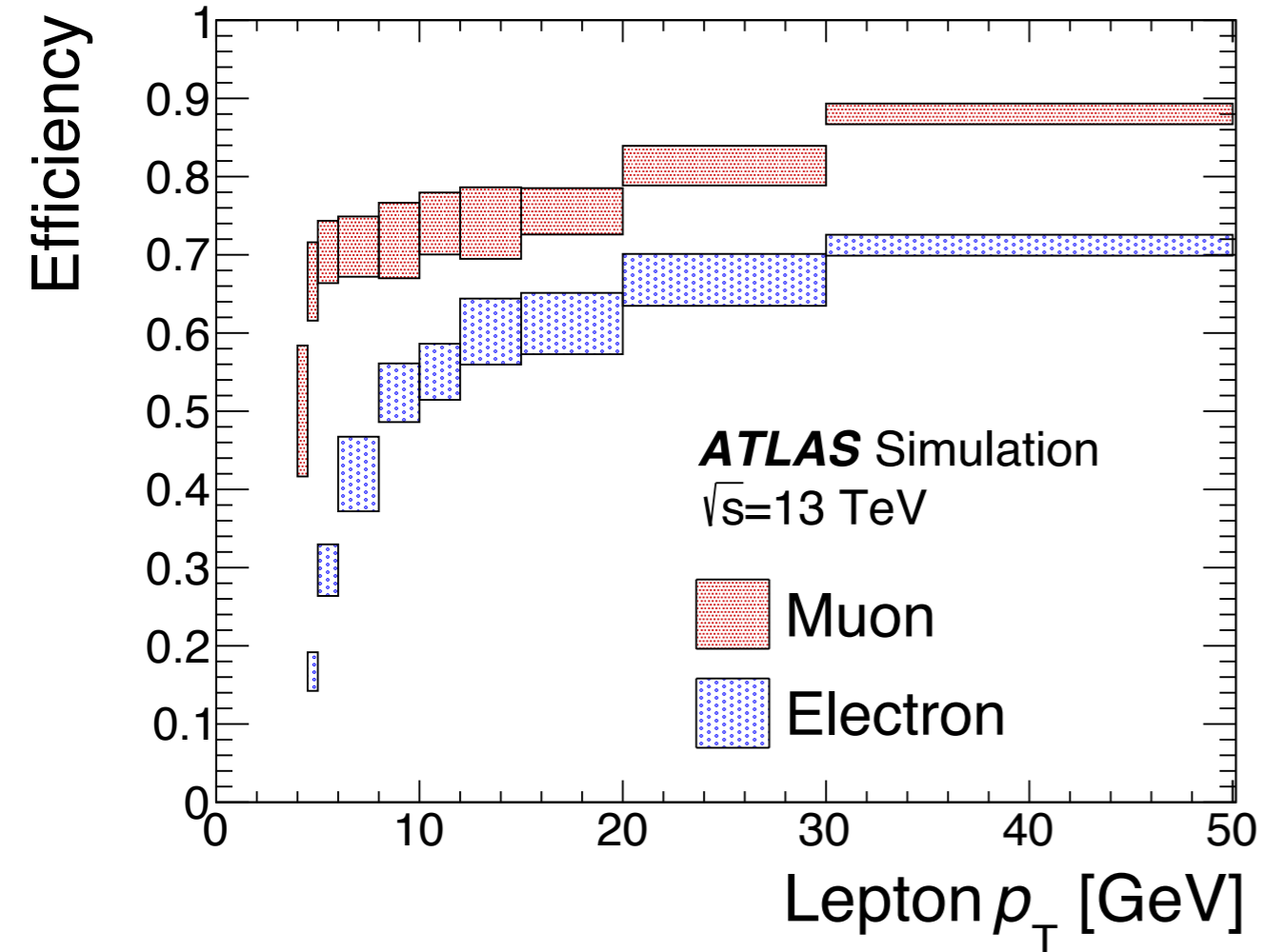
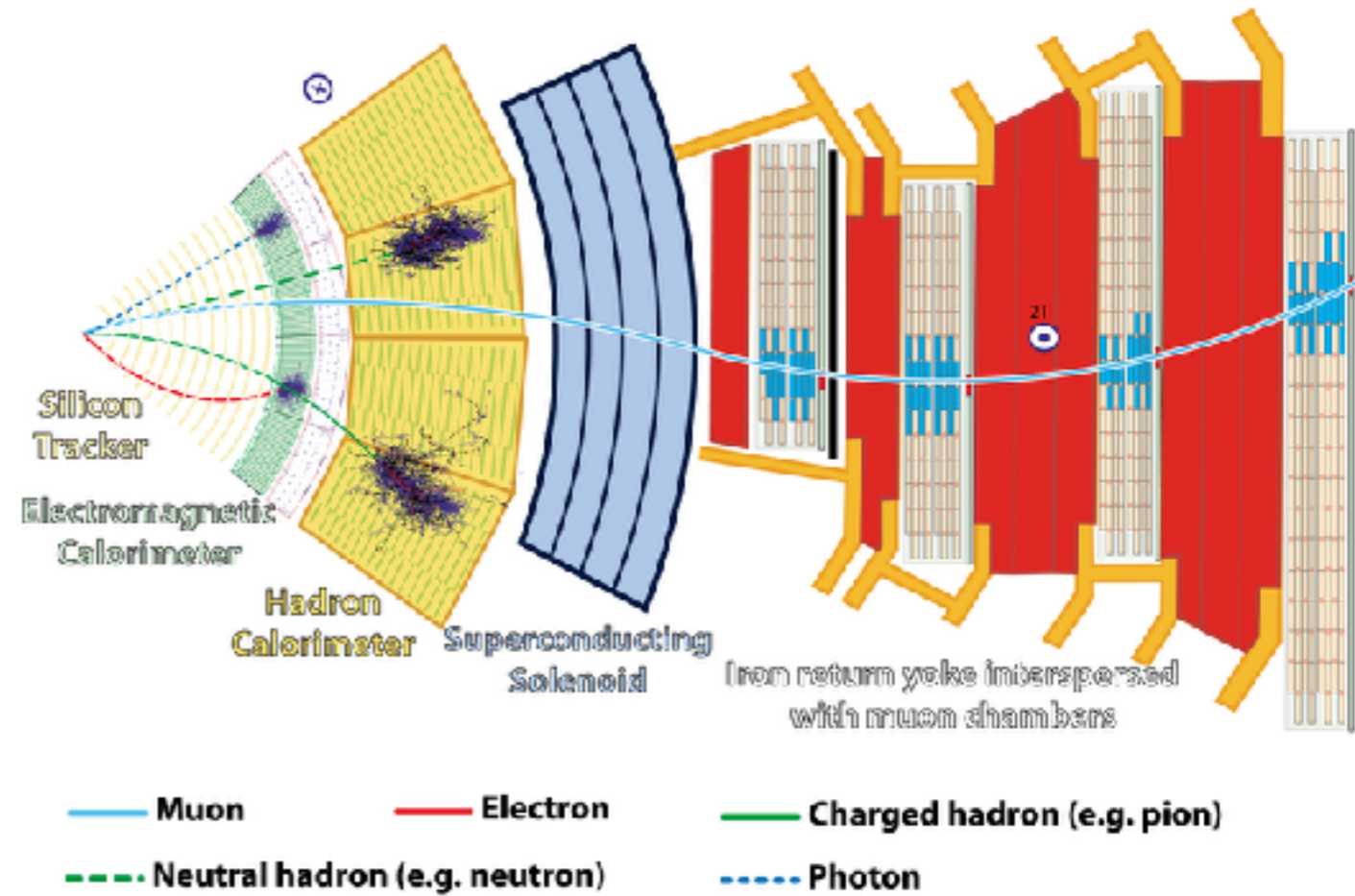
Adopted by ATLAS [1712.08119] and CMS [1801.01846] 13 TeV, 36.1 fb⁻¹

Sketches by Jesse Liu

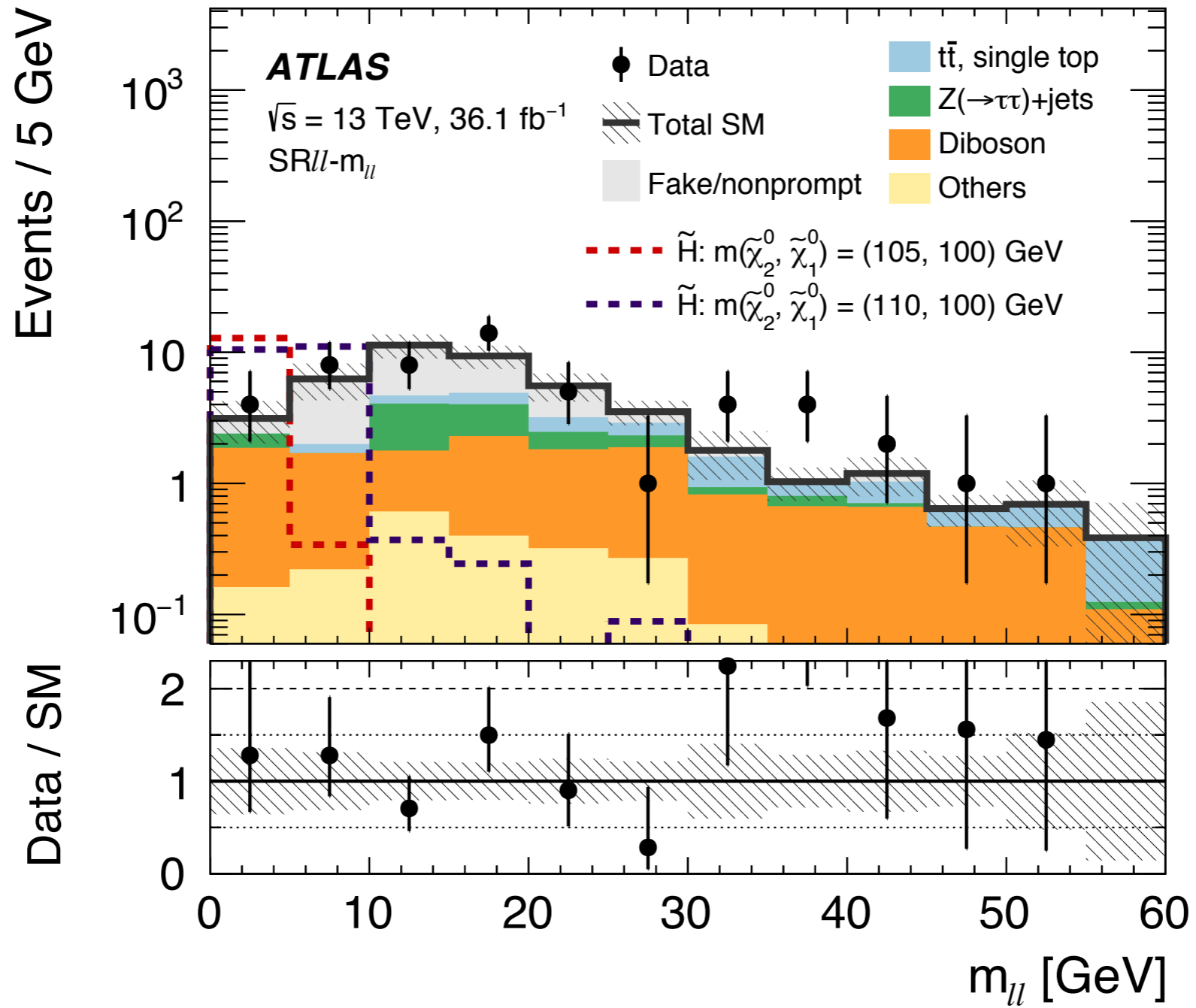
Detector challenges

Muon energy loss in calorimeter ~ 3 GeV

Electron cluster reconstruction in calorimeter, ≈ 5 GeV

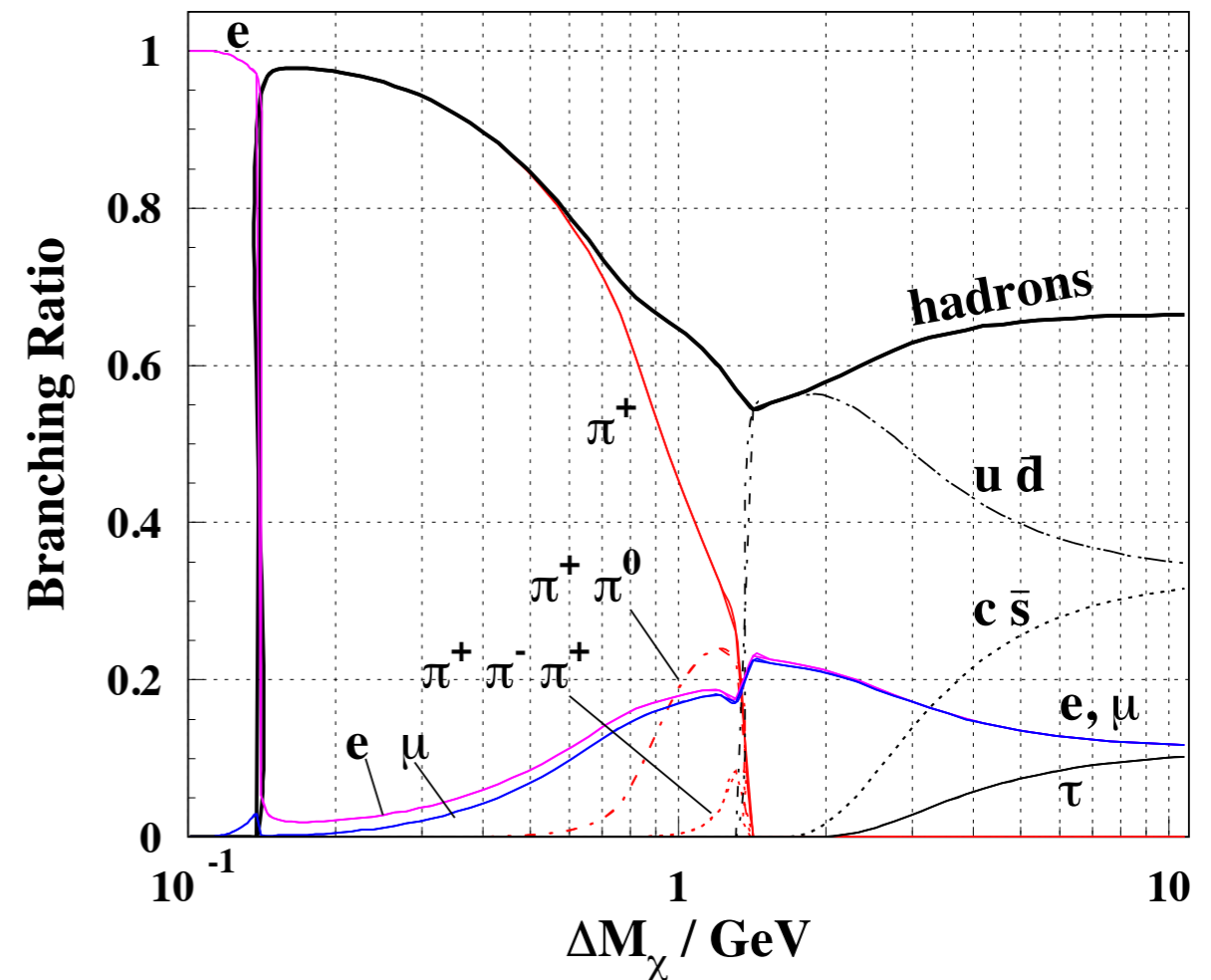
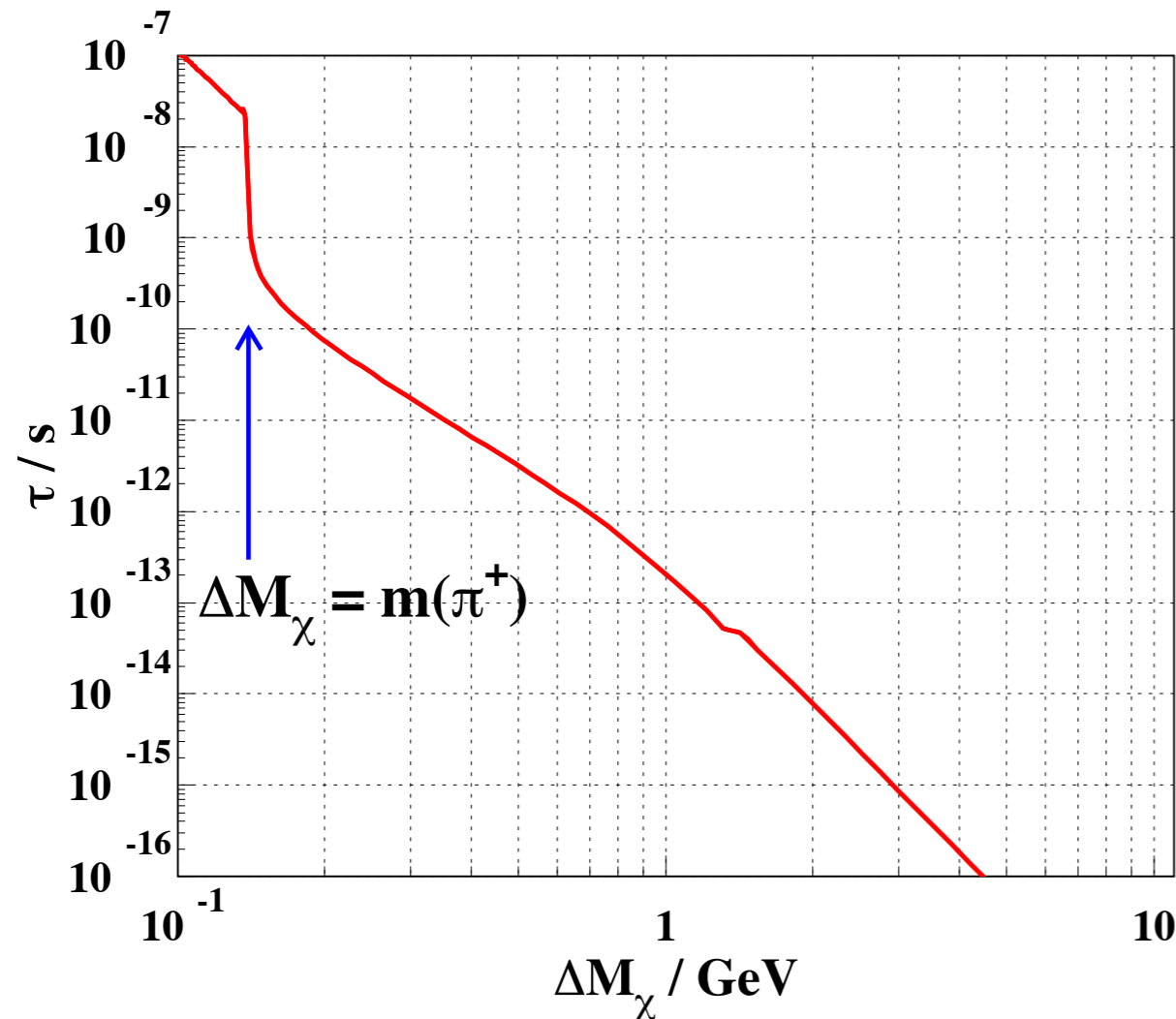
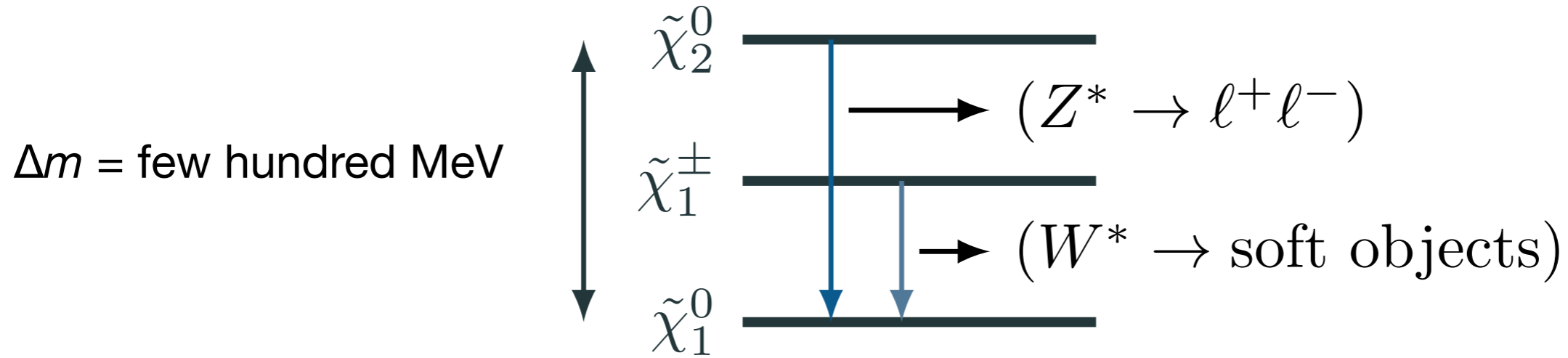


Inspecting data



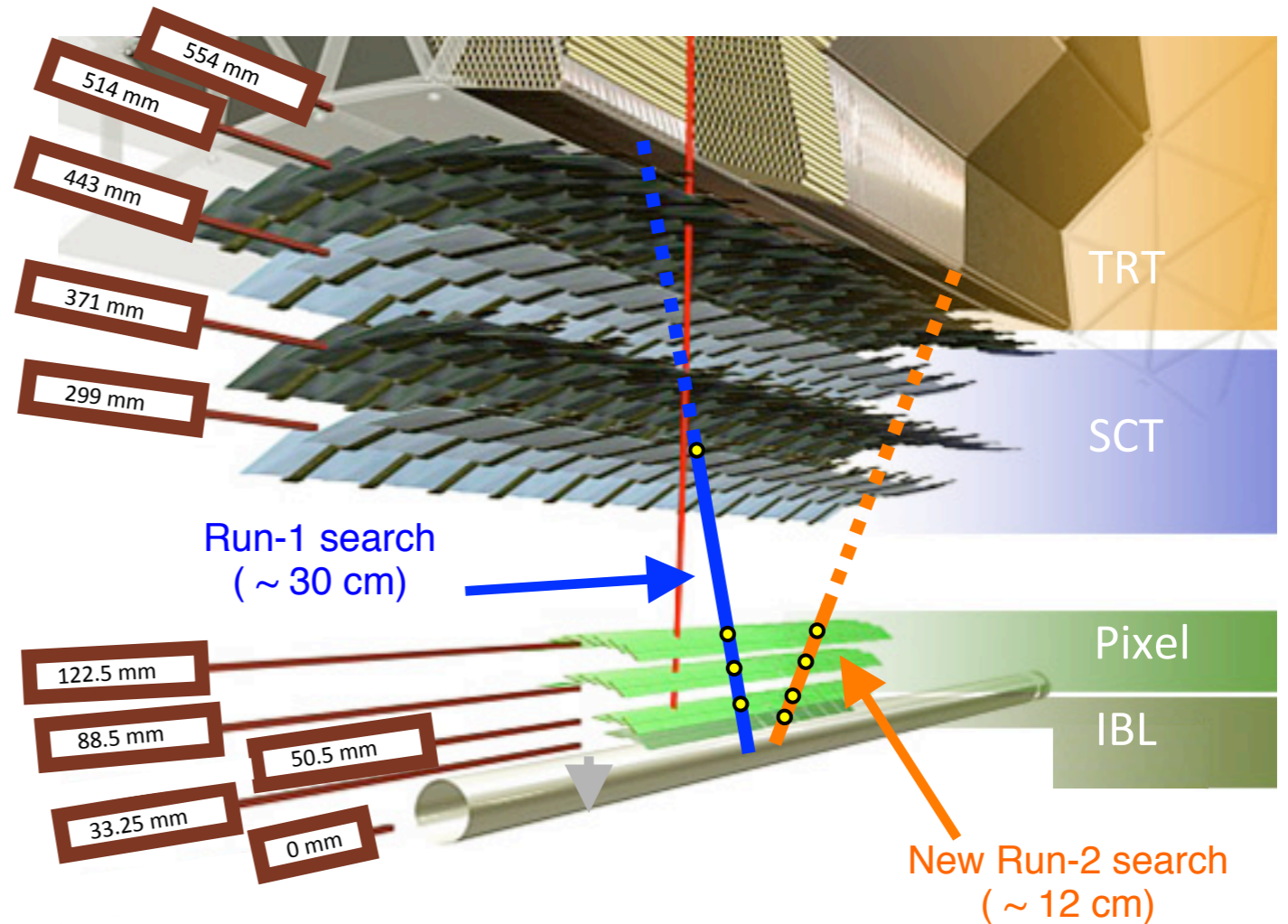
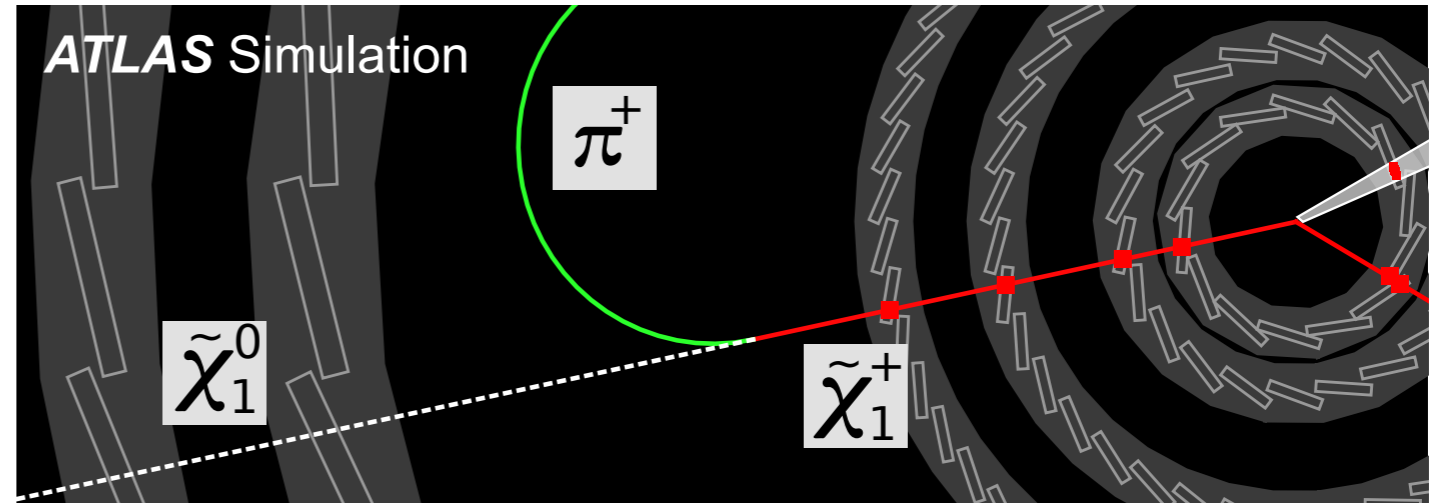
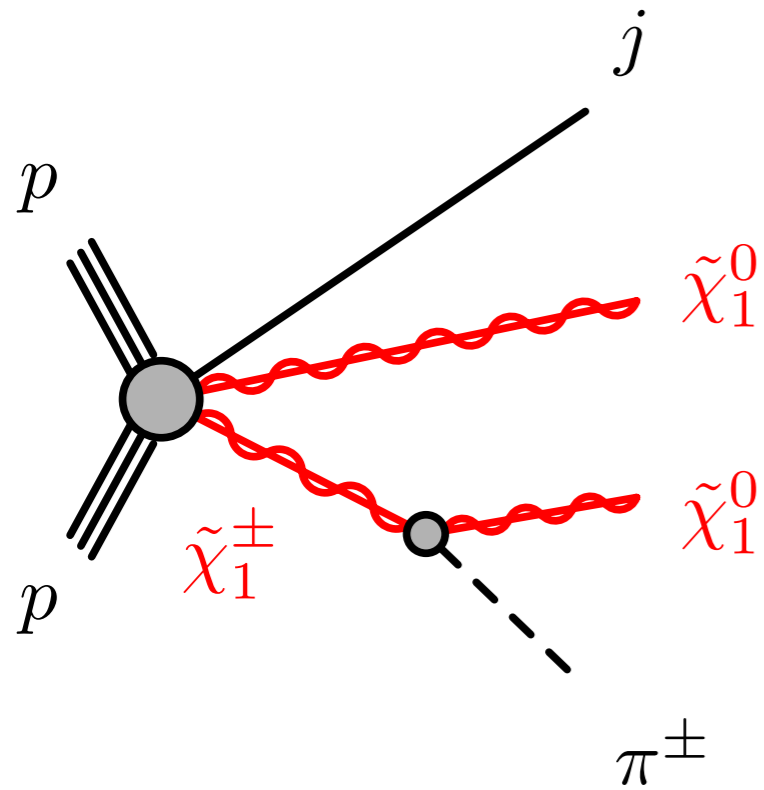
Signal has kinematic endpoint $m_{ll} < \Delta m(\chi^0_2, \chi^0_1)$

Ultra compressed spectrum



Disappearing track

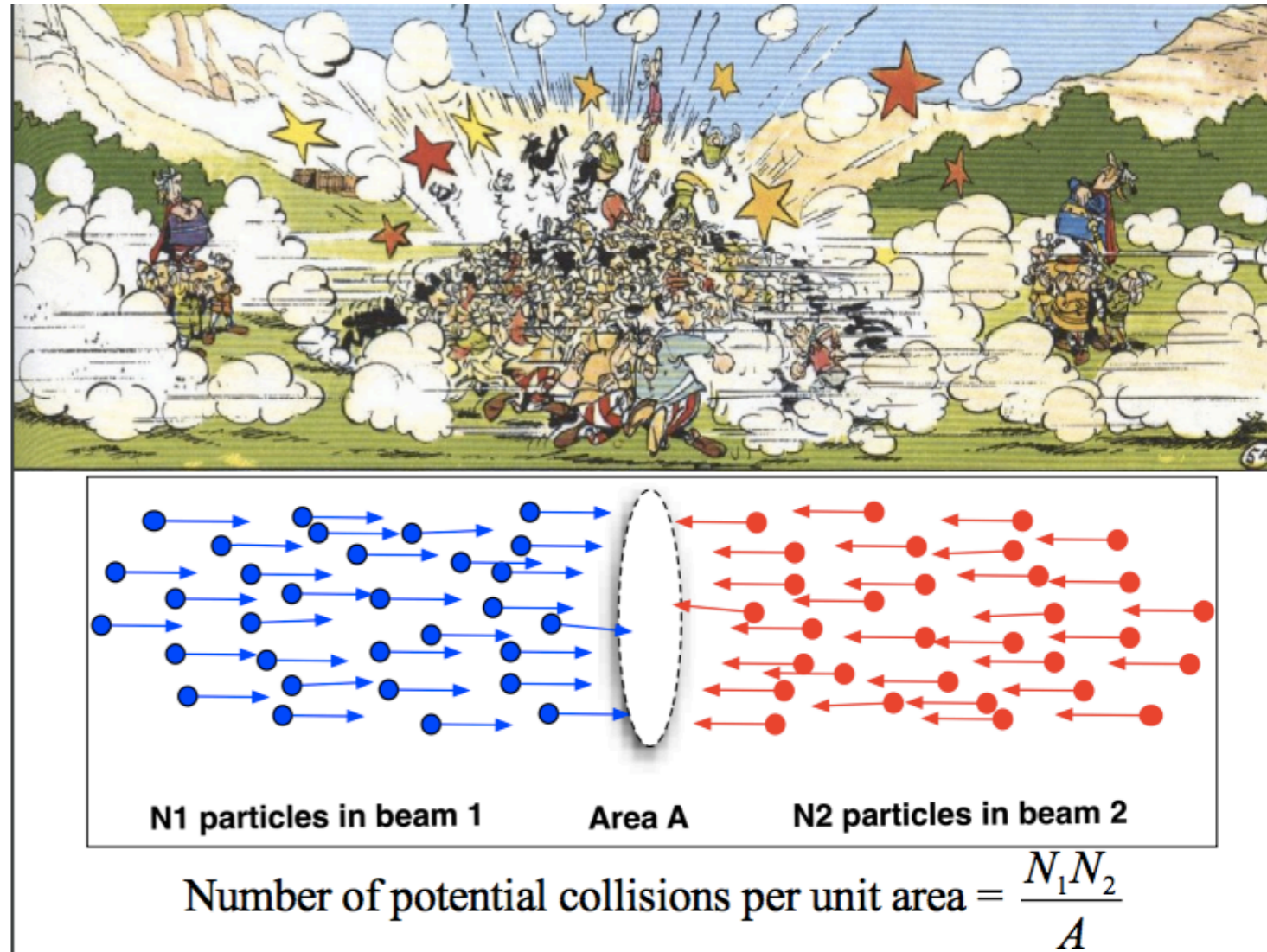
Hadronic recoil (ISR), trigger!



Experimental challenge: *pileup*

Many pp interactions (i.e. collisions) in the same bunch crossing = pileup.

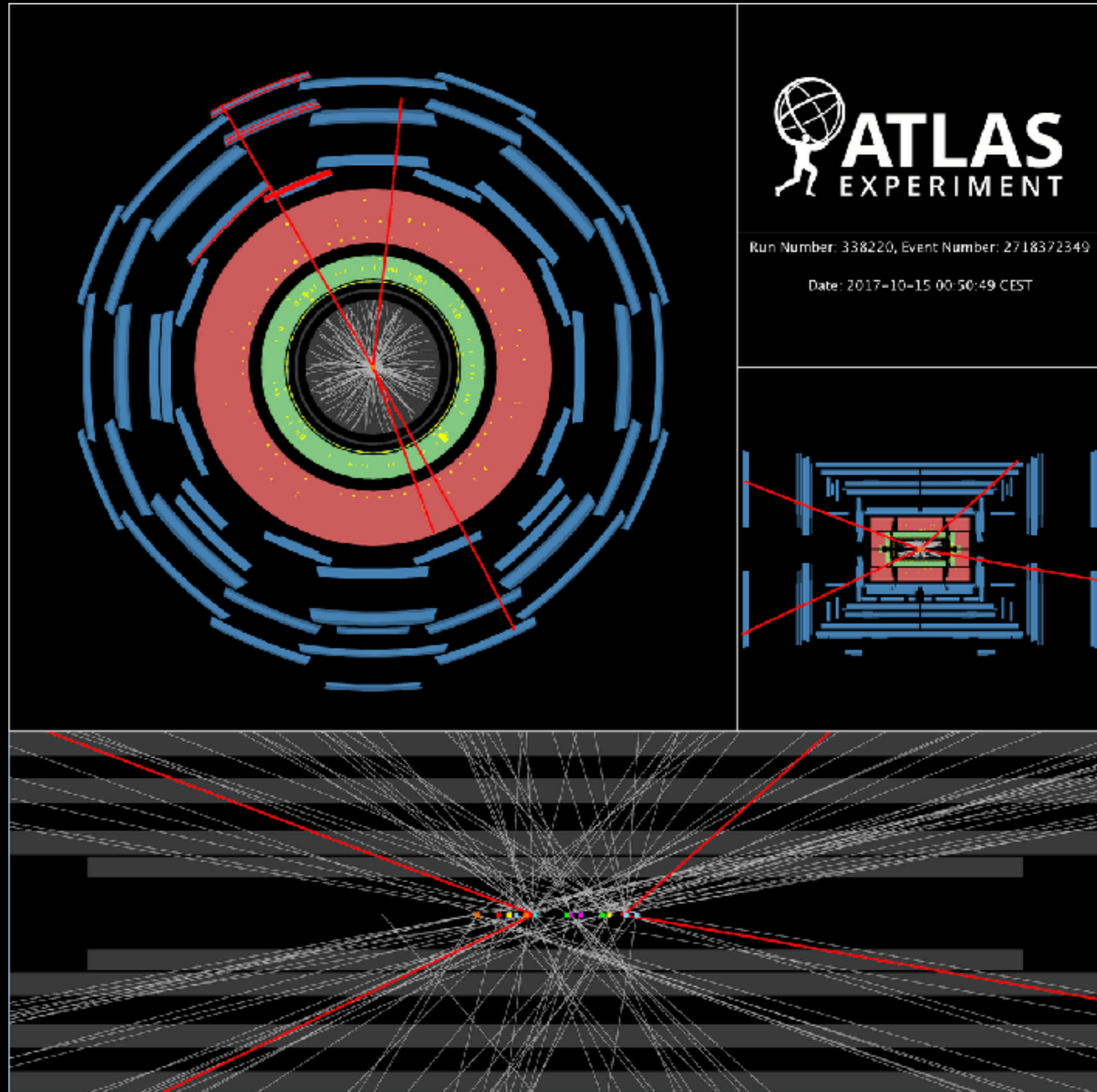
Average of 38 collisions per crossing in 2017.



Many charged soft particles that leave hits in inner-detector.
Challenge for disappearing track (random hits bkg), and many other areas!

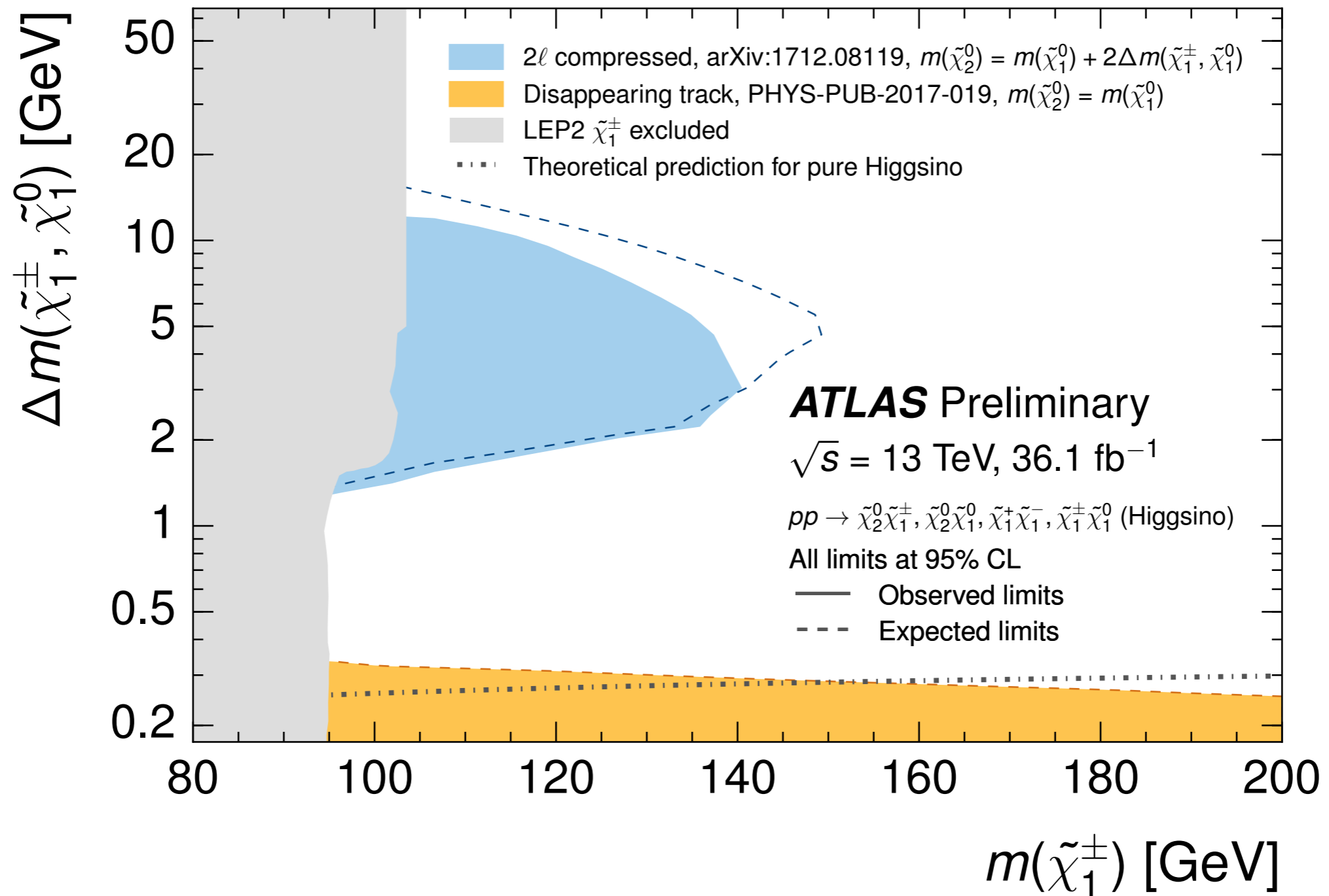
High pileup event display

ATLAS 2017 collision data event with **two Z-boson** candidates each decaying to two muons and originating from well separated pp interactions in the same LHC bunch crossing. The production vertices of the two Z boson candidates are separated by 67 mm.



Hadron collider extends nearly 20 years old LEP limits

March 2018



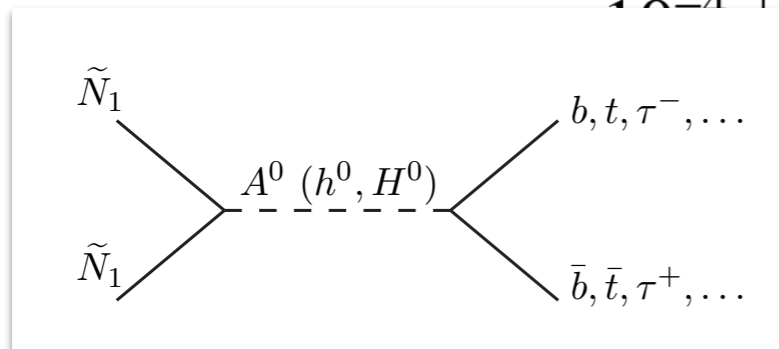
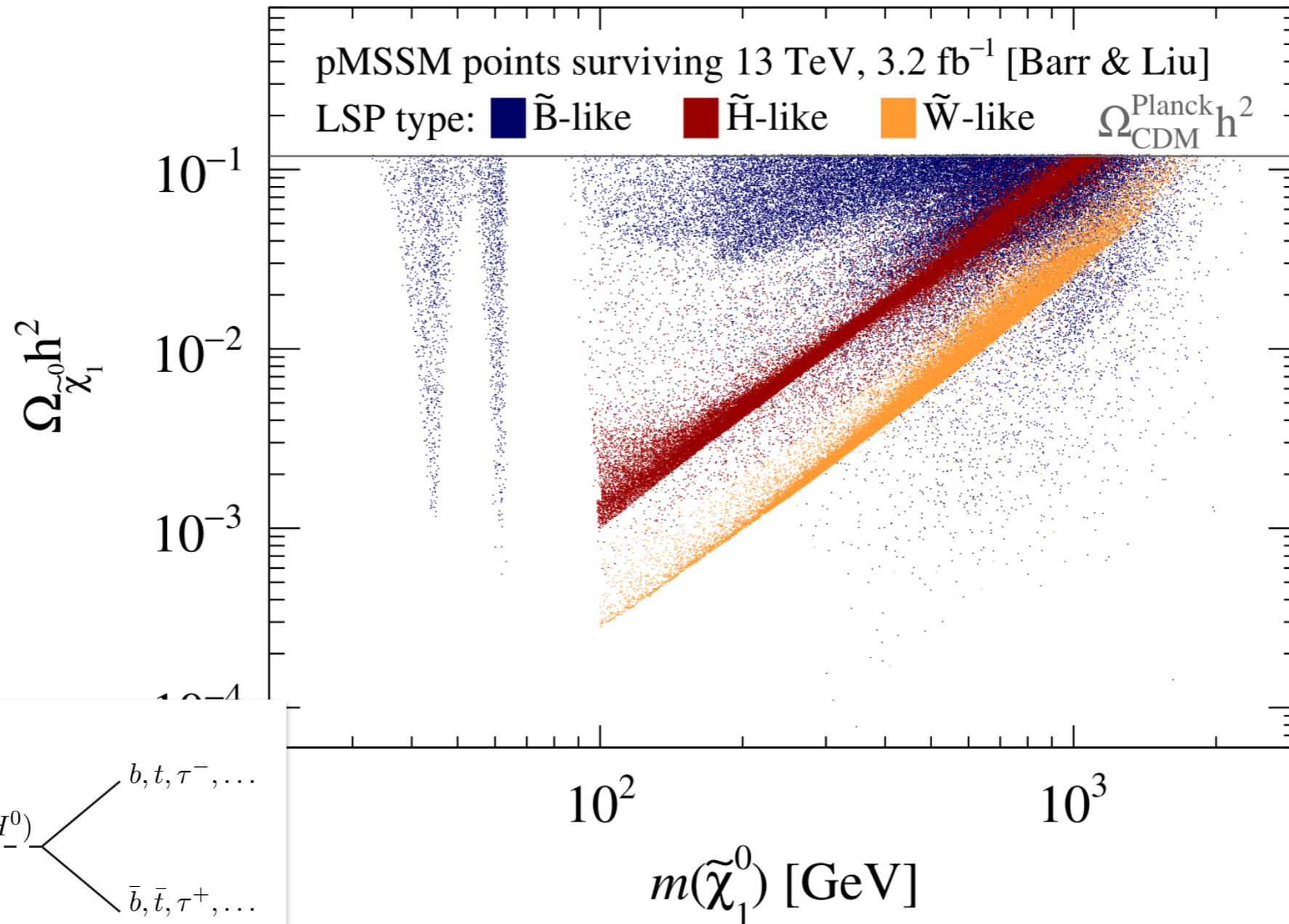


Outlook

Next steps in the hunt for supersymmetry.

SUSY comes with Dark Matter WIMP

Neutralino dark matter: composition controls relic abundance



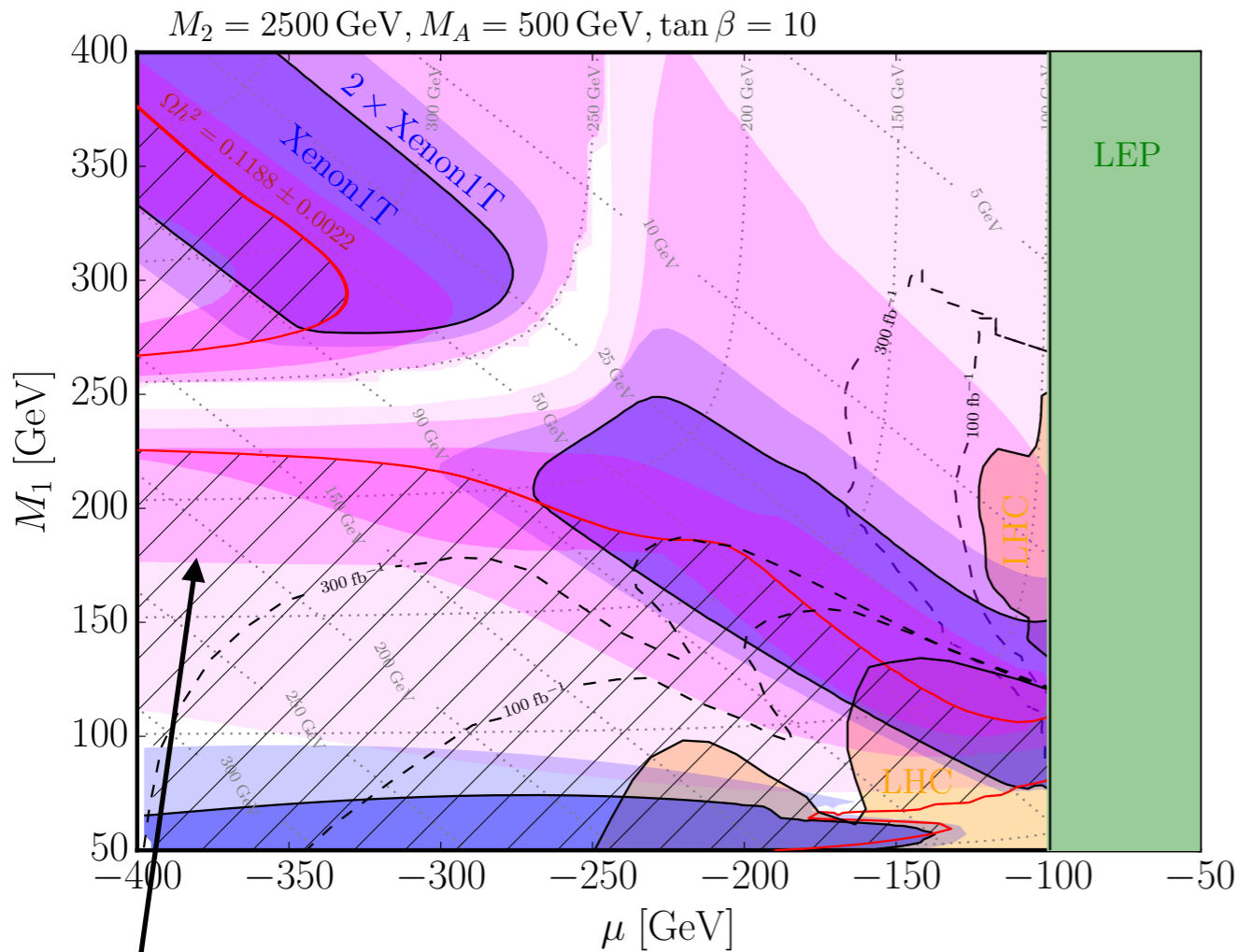
Bino dark matter annihilation suppressed, other processes ensure $\Omega_{\tilde{\chi}_1^0} \leq \Omega_{\text{CDM}}^{\text{Planck}}$

Higgsino & wino dark matter bands: mass controls abundance $\Omega_{\tilde{\chi}_1^0} h^2 \sim \langle \sigma v \rangle^{-1} \sim m_{\tilde{\chi}_1^0}^2$

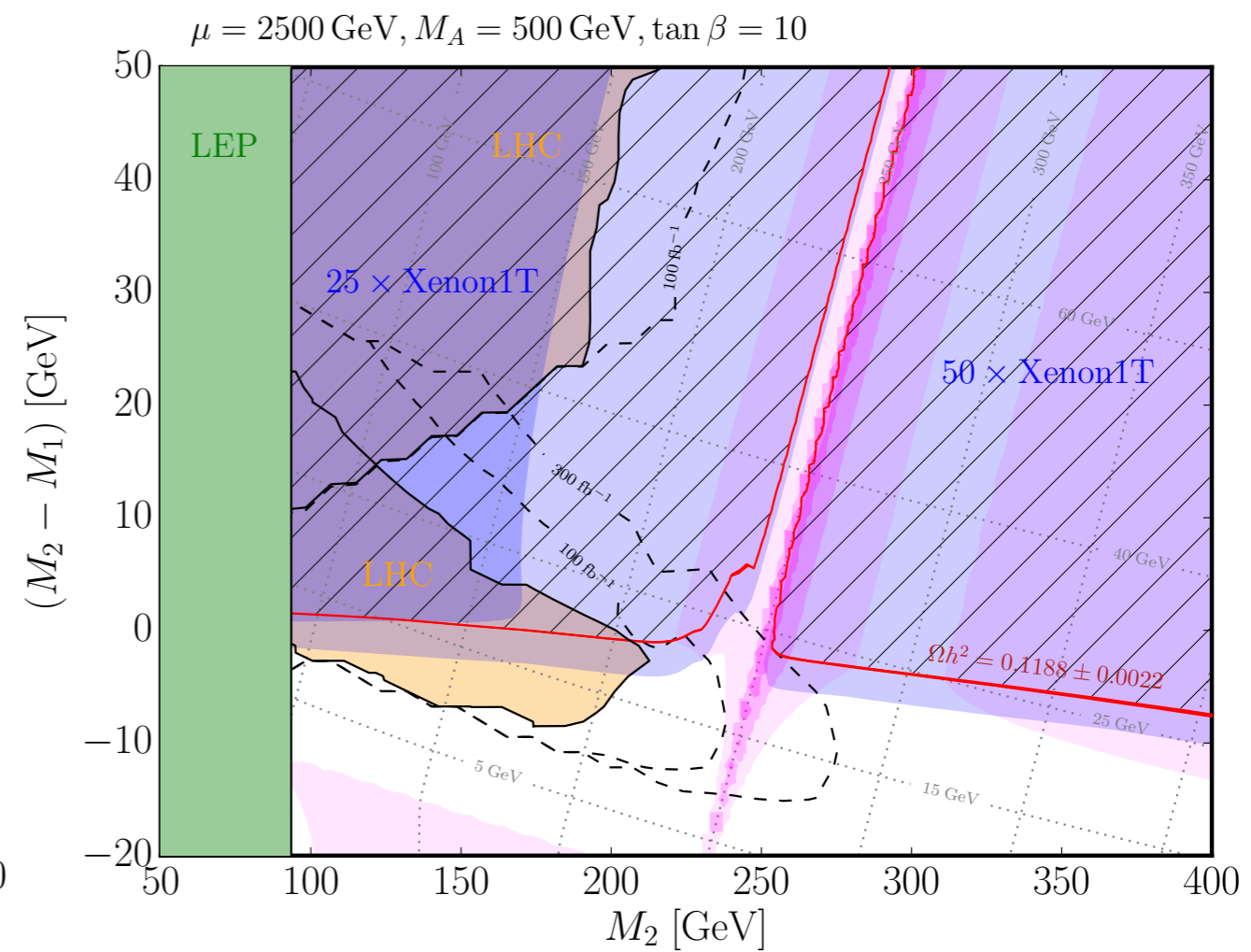
Barr & Liu [1608.05379]

Well Tempered Neutralino

well-tempered bino-Higgsino neutralino



well-tempered bino-wino neutralino



Profumo, Stefaniak, Haskins [1706.08537]

Projected limits from **DM indirect** detection experiments scaled by current limit is scaled by factors of 10, 100 and 1000, with high to low opacity.

For large μ values direct detection quickly becomes ineffective.

Complementarity of dark matter programs

Estimates of future LHC sensitivity at 95% CL

| | ~35 ifb 13 TeV | Run2 (140 ifb) 13 TeV | Run3 (~300 ifb) 13/14 TeV | HL-LHC (3000 ifb) 13/14 TeV |
|-----------------------------------|---|------------------------------|--|---|
| gluino | 2 TeV (prel.) | ~2.3 TeV (my est.) | 2.4 TeV (upgrade), 2.4/2.6 TeV (my est.) | 2.9 TeV (upgrade), 2.9/3.1 TeV (my est.) |
| squark (x8) decoupled | 1.5 TeV (prel.) | 1.75 TeV (my est.) | 1.9/2.0 TeV (my est.) | 2.3/2.4 TeV (my est.) |
| stop | 1-1.1 TeV (prel.) | ~1.3 TeV (my est.) | 1.4x TeV (my est.) | 1.85 TeV (my est.) |
| wino C1N2 to WZ bino | ~600 GeV (prel.) | 670 GeV (my est.) | 780 GeV (my est.), 750 GeV (CMS est.), 840 GeV (upgrade) | 1150 GeV (my est.), 1.2 TeV (CMS est.), 1.1 TeV (upgrade) |
| wino C1C1 to WW bino | 225 GeV (based on x-section ratio for 180 GeV Run1) | ~320 GeV (my est.) | 380-400 GeV (my est.) | ~630 GeV (my est.) |
| wino LSP, pixel-trklet | 420 GeV (prel.) | 580 GeV (my est.) | 680 GeV (my est.) | 1030 GeV (my est.) |
| higgsino LSP, DM=3-20 GeV | 150 GeV | 200-250 GeV (my est.) | 250-300 GeV (my est.) | 450-500 GeV (my est.) |
| higgsino LSP, DM=0.3-3 GeV | ? | | | |
| slepton (sel, smu) | ~500 GeV (prel.) | 670 GeV (my est.) | | |
| stau | ? | | | 700 GeV (upgrade) |

Summary

Harvesting large LHC dataset
wide Supersymmetry search program; so far no clear sign of new physics.

SUSY frontiers at LHC

Complex scenarios: new clever ideas,
sophisticated analysis techniques.

(uncovered signatures .. long-lived, R -parity
violating, etc.)

Electroweak SUSY production,
complementing dark matter direct & indirect detection

(analysed ~1% of the expected LHC p-p dataset)



A very exciting and puzzling time for particle physics.