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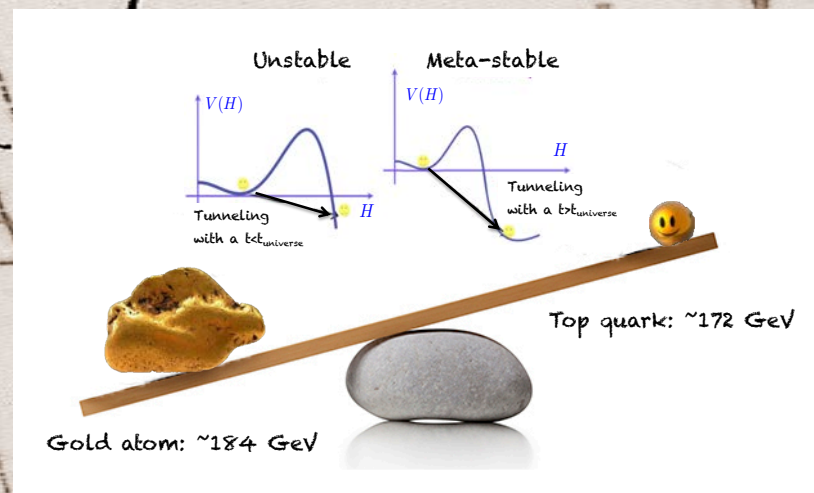


VNIVERSITAT  
D VALÈNCIA

The top quark is *beautiful, pretty charming*  
and still very *strange*

(the *ups* and *downs* of the top mass)

Juan A. Fuster Verdú – IFIC, València  
DESY-Zeuthen Colloquium  
Zeuthen, July 10<sup>th</sup> 2019



# Thanks and acknowledgements

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Special thanks to Katerina Lipka as DESY host

and for the support of

Sven-Olaf Moch (Hamburg), Peter Uwer (Berlin), Norbert Wermes (Bonn)



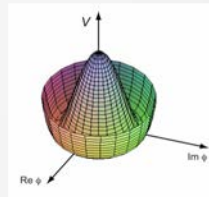
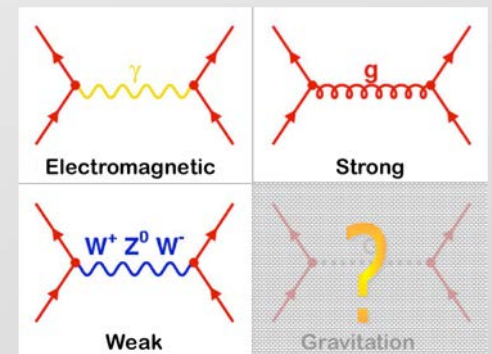
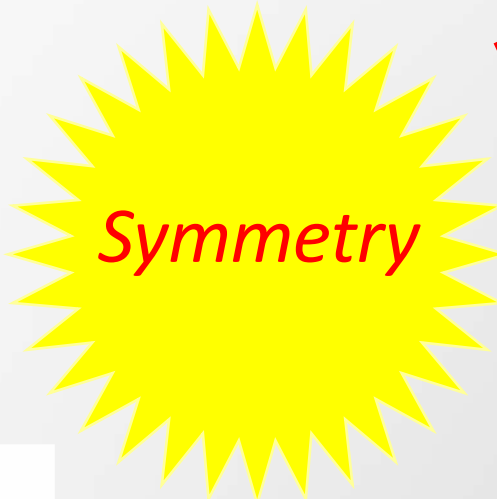
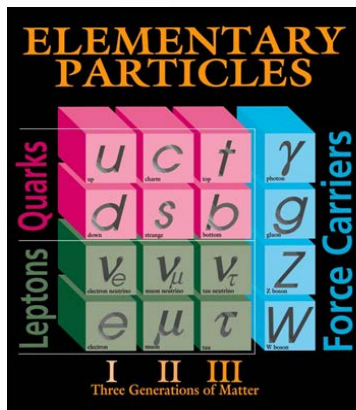
“Never underestimate the pleasure people get when they listen to something they already know”, (E. Fermi).

# The Standard Model

## Standard Model

$$\mathcal{L}_{\text{SM}} = \mathcal{L}_{\text{Gauge}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{Yukawa}}$$

(~1980)



Higgs: EWSB

Particles

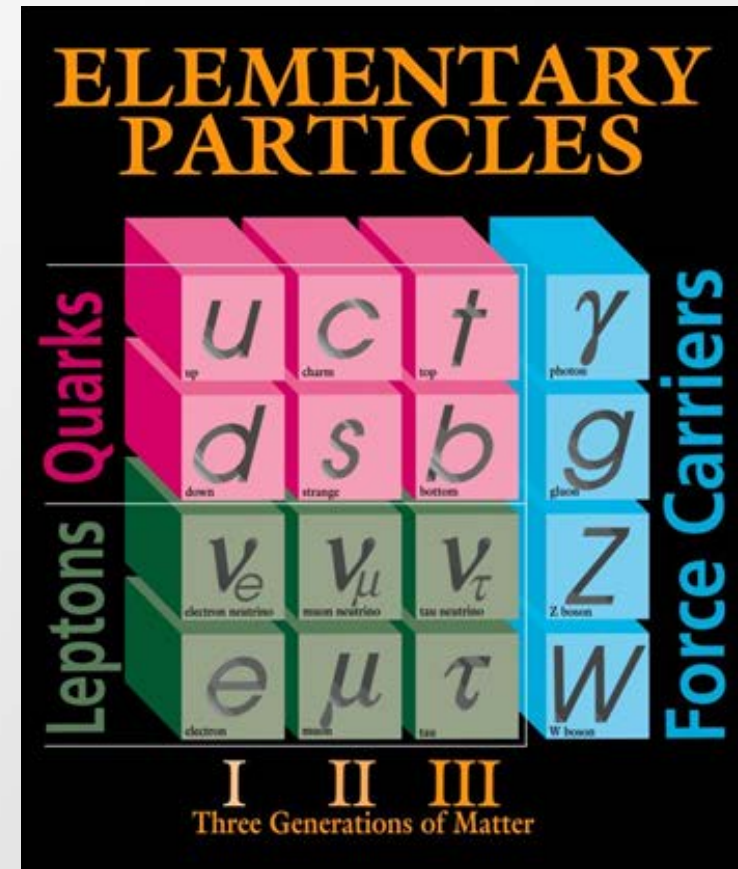
Interactions/Forces

# Why the Standard Model is as it is ?

- First and last column are necessary but what about the others !
- **Averroes** (Córdoba XII a.C.) says: “Nothing in nature is superfluous”
- We want to understand the reasons behind this structure,

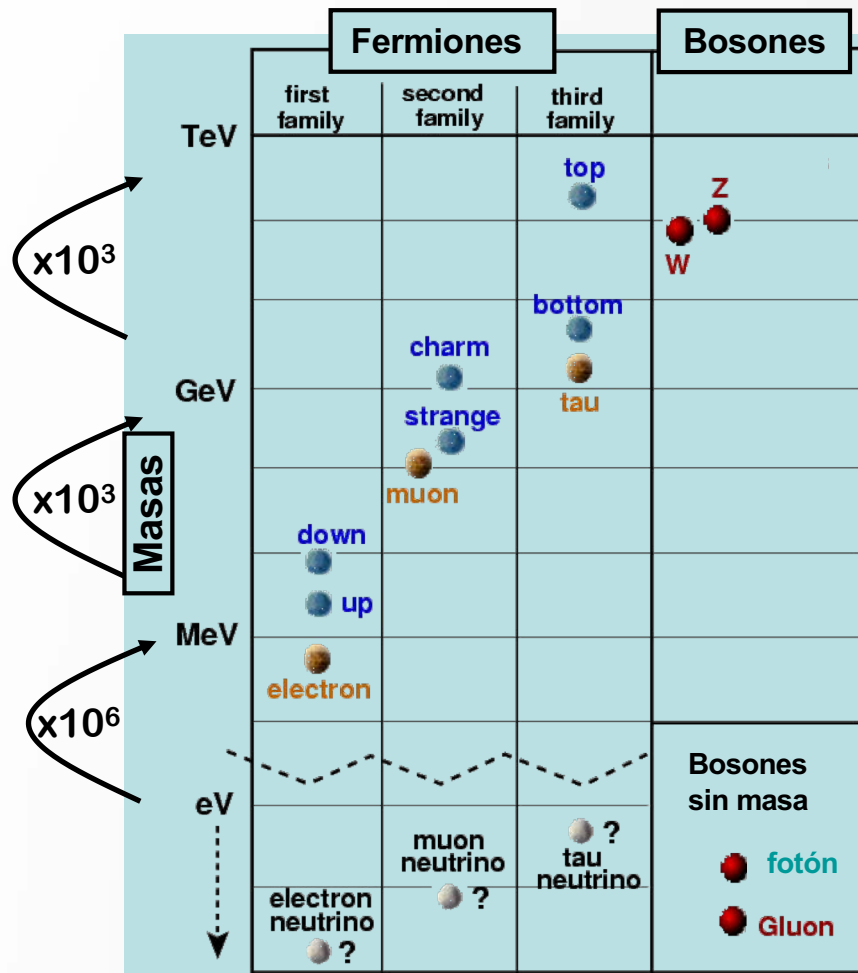
**Why ?**

**Why the families are duplicated but with different mas ?**



# The masses of the elementary particles

## Large spectrum of masses



$$m_\gamma = 0 \quad i.e. < 6 \cdot 10^{-26} \quad \text{GeV}$$

$$m_{W^\pm} = 80.425 \pm 0.038 \quad \text{GeV}$$

$$m_{Z^0} = 91.1876 \pm 0.0021 \quad \text{GeV}$$

$$m_g = 0$$



Averroes: why ?

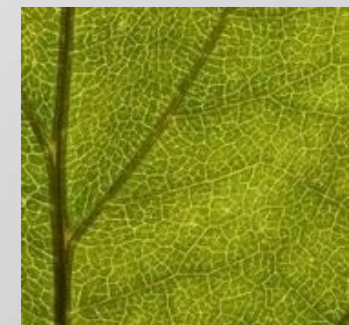
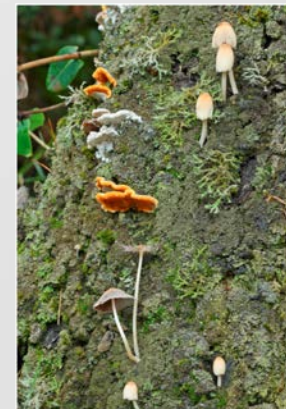
# The masses of the elementary particles

- Masses cover **at least ~12 orders of magnitude**
- What does it mean ?
- Similar differences in scale imply very complex objects but we call all them **“elementary”**
- **Averroes:** are they really elementary ?

- Earth =  $12,74 \times 10^{+6}$  m



- Vegetal cells  $\sim 1 \times 10^{-[4,5]}$  m



# The masses of the elementary particles

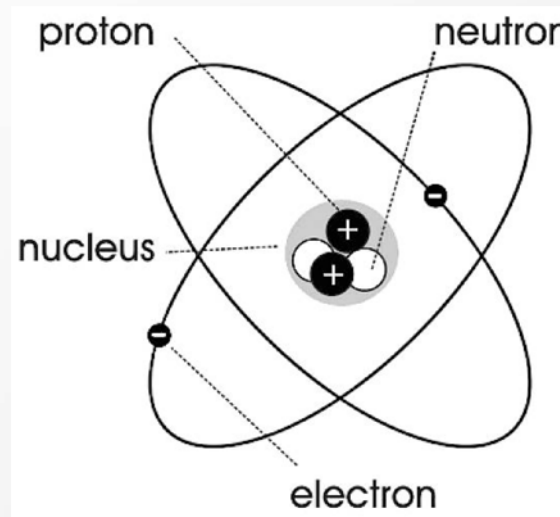
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- Is it by **chance** or is it a **necessity** ? In other words can we change them without any consequence in the Universe ?
- There are cases in which such changes imply dramatic consequences
- Let's modify "slightly" the masses of the 1<sup>st</sup> family  
**u d e**



# The masses of the elementary particles

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- For instance a **neutrón udd** could be lighter than a **proton uud**. But then **Hydrogen** would not be stable.
- $M_{\text{electron}} \neq 0$ . If the mass of the electron would be zero the nucleus could NOT capture electrons. No atoms.
- **Averroes: why?**

# The different faces of the mass



- **1687 Newton:** inertial mass, resistance to movement, laws of gravitation



$$\vec{F} = m\vec{a}$$

# The different faces of the mass



- **1687 Newton:** inertial mass, resistance to movement, laws of gravitation



$$\vec{F} = m\vec{a}$$

- **1905 Einstein:** equivalence between mas-energy, proton mass is condensed energy of quarks-gluons



$$E = mc^2$$

# The different faces of the mass



- **1687 Newton:** inertial mass, resistance to movement, laws of gravitation



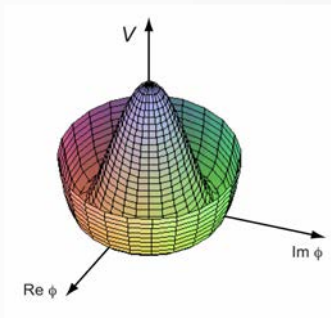
$$\vec{F} = m\vec{a}$$

- **1905 Einstein:** equivalence between mass-energy, proton mass is condensed energy of quarks-gluons



$$E = mc^2$$

- **1964 Brout-Englert-Higgs:** connects mass and the vacuum. The vacuum is not the “absence of things” but the “level of minimum energy”



$$|\phi| = a = \left(\frac{-\mu^2}{2\lambda}\right)^{1/2}$$

$$\phi(x) = a + \frac{\phi_1(x) + i\phi_2(x)}{\sqrt{2}}$$

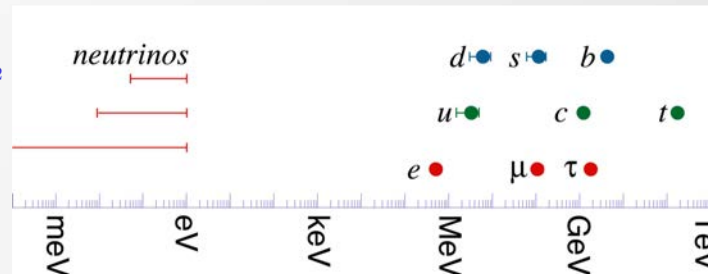
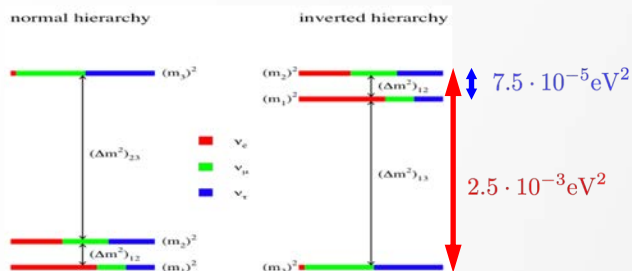
**Averroes: why ?**

# The masses of the elementary particles

The Standard Model boundaries:  
**massive  $\nu$**



The Standard Model boundaries:  
**Connection: top-quark mass, Higgs mass and  $\alpha_s$**



$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U_{PMNS}(\theta_{12}, \theta_{23}, \theta_{13}, \delta, \dots) \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$$\theta_{12} \sim 34^\circ$$

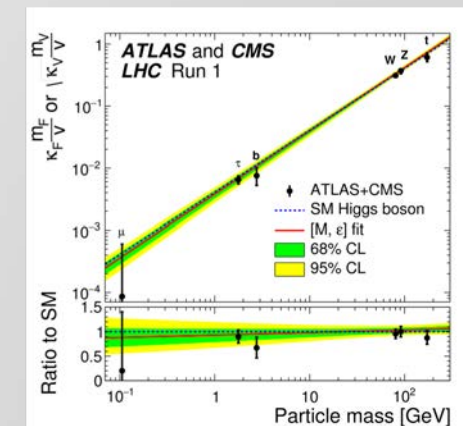
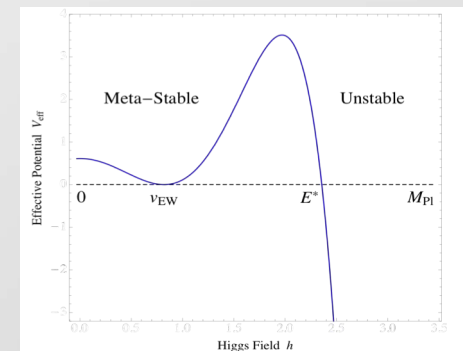
$$\theta_{23} \sim 42^\circ \text{ o } 48^\circ$$

$$\theta_{13} \sim 8.5^\circ$$

$$\delta \sim ?$$

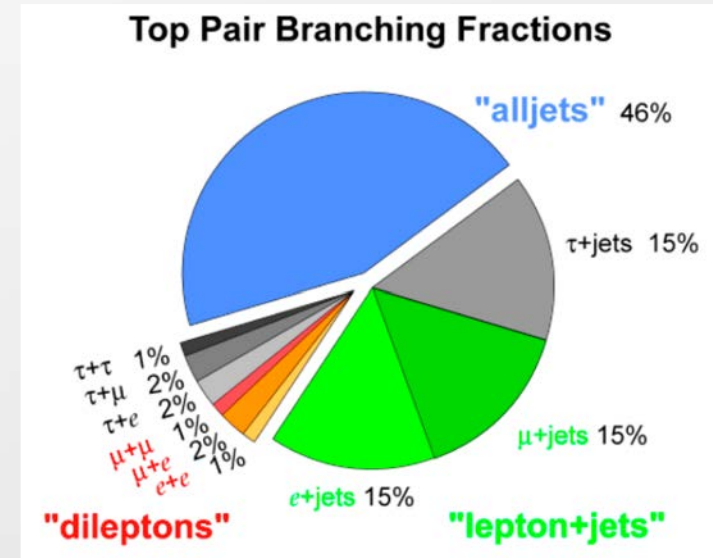
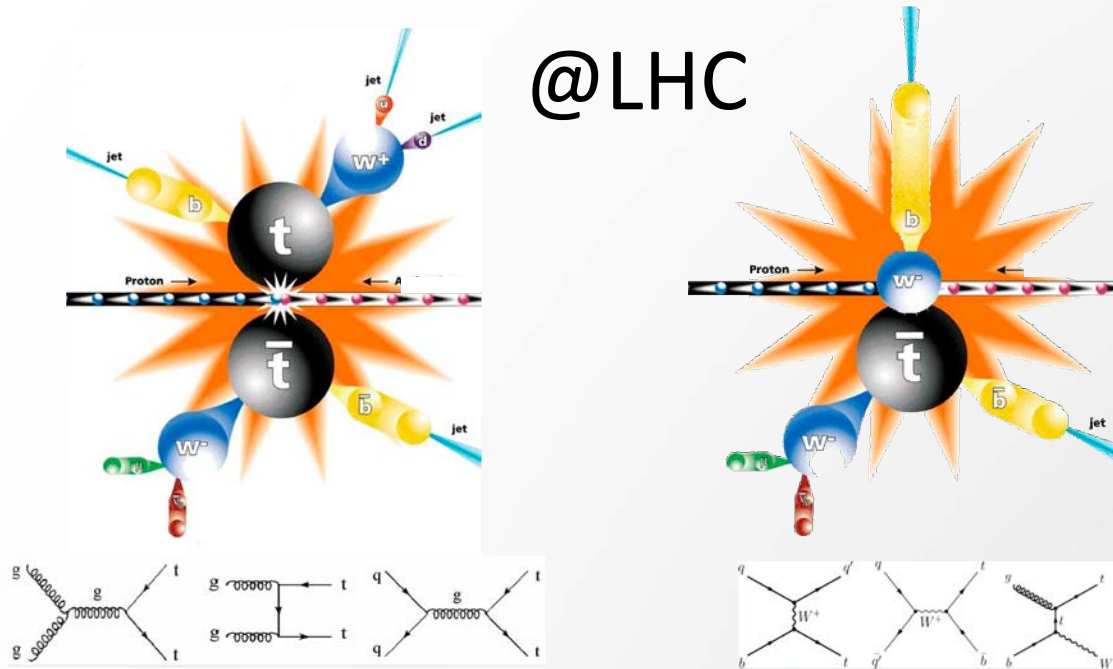
The **borders** of the mass spectrum region represent a golden place to test the validity of Standard Model therefore to look for its consistency and/or new physics.

**Averroes: Neutrinos and top-quark**



# The top quark is *beautiful and charming*

@LHC



- The heaviest known elementary particle
- Yukawa coupling to Higgs boson  $y_t = \mathcal{O}(1)$ : privilege position to test EWSB
- Special role in many BSM: a window to new physics that couples preferentially to top quarks
- Decays before hadronizing: the only “naked” quark

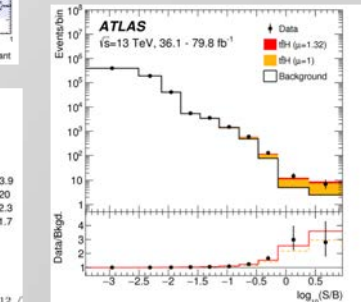
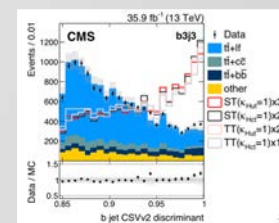
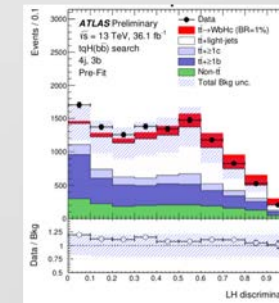
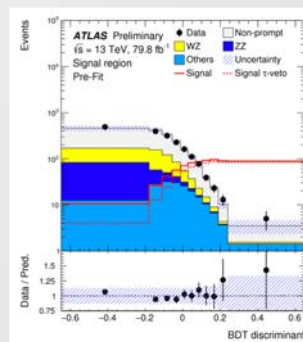
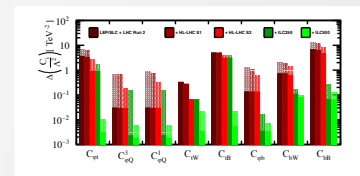
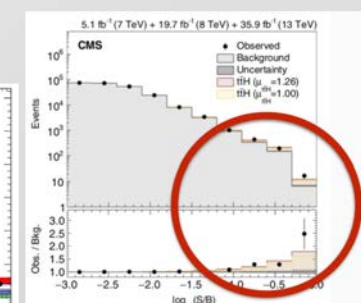
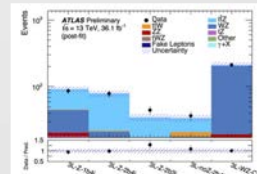
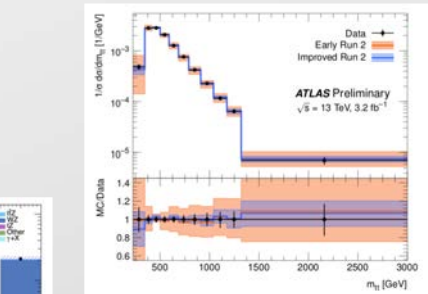
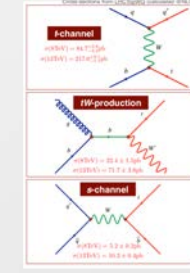
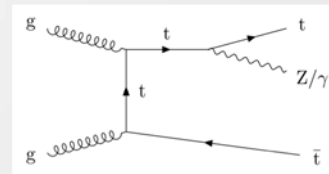
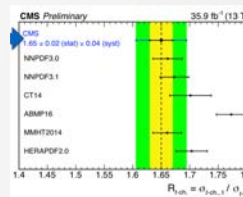
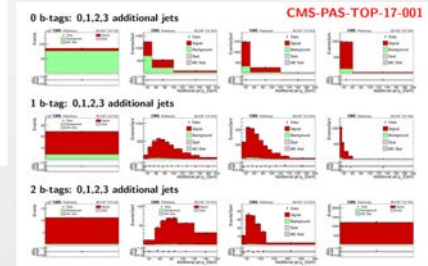
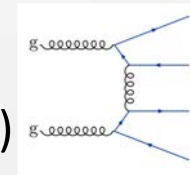
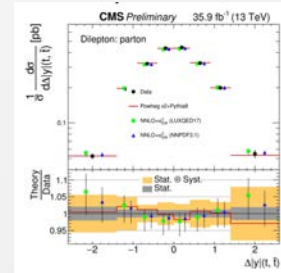
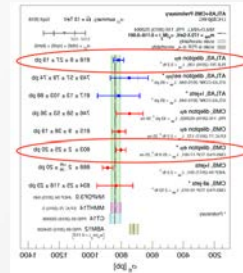
$$\tau_{\text{had}} \approx h/\Lambda_{\text{QCD}} = 2 \cdot 10^{-24} \text{ s}$$

$$\tau_{\text{top}} \approx h/\Gamma_{\text{top}} = 1/(G_F m_t^3 |V_{tb}|^2 / 8\pi\sqrt{2}) = 5 \cdot 10^{-25} \text{ s}$$

$$\tau_{\text{bottom}} \approx 10^{-12} \text{ s}$$

# The top quark is *beautiful and charming*: rich physics programme

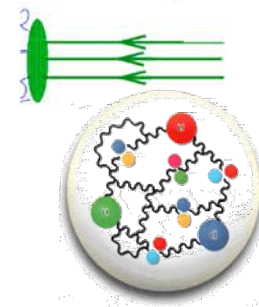
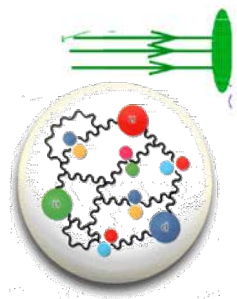
- Top properties:
  - Cross/differential sections ( $tt+0$ jets,+1jets,+2jets,+3jets)
  - Spin correlations
  - Charge asymmetries
  - Top Yukawa
  - Colour flow
- Single top production
- New Physics
  - Anomalous couplings
  - Flavour Changing neutral Currents ( $tqZ,tqH$ )
  - $tt\gamma$
- Top mass
  - “Direct” measurements
  - “Alternative” measurements



Disclaimer: This talk only covers part of the “top mass” topic

# The top quark is *beautiful and charming* but LHC is challenging

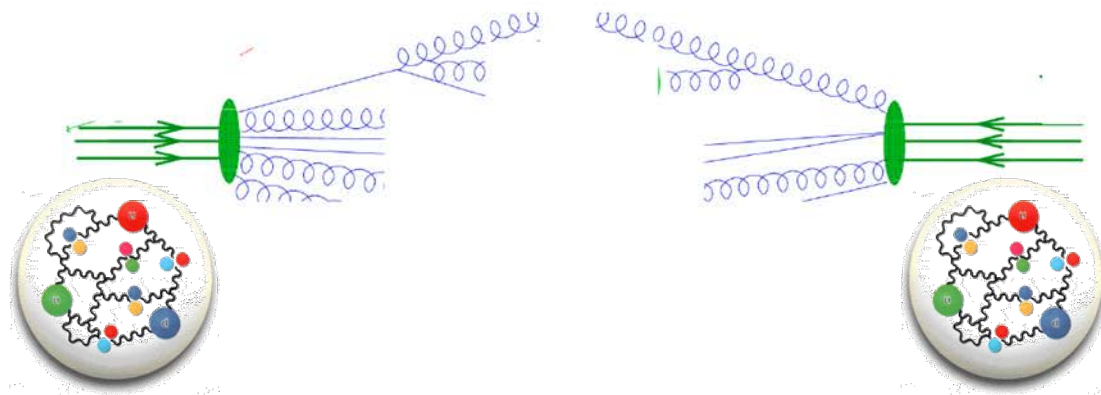
- Parton Distribution Functions (PDFs)





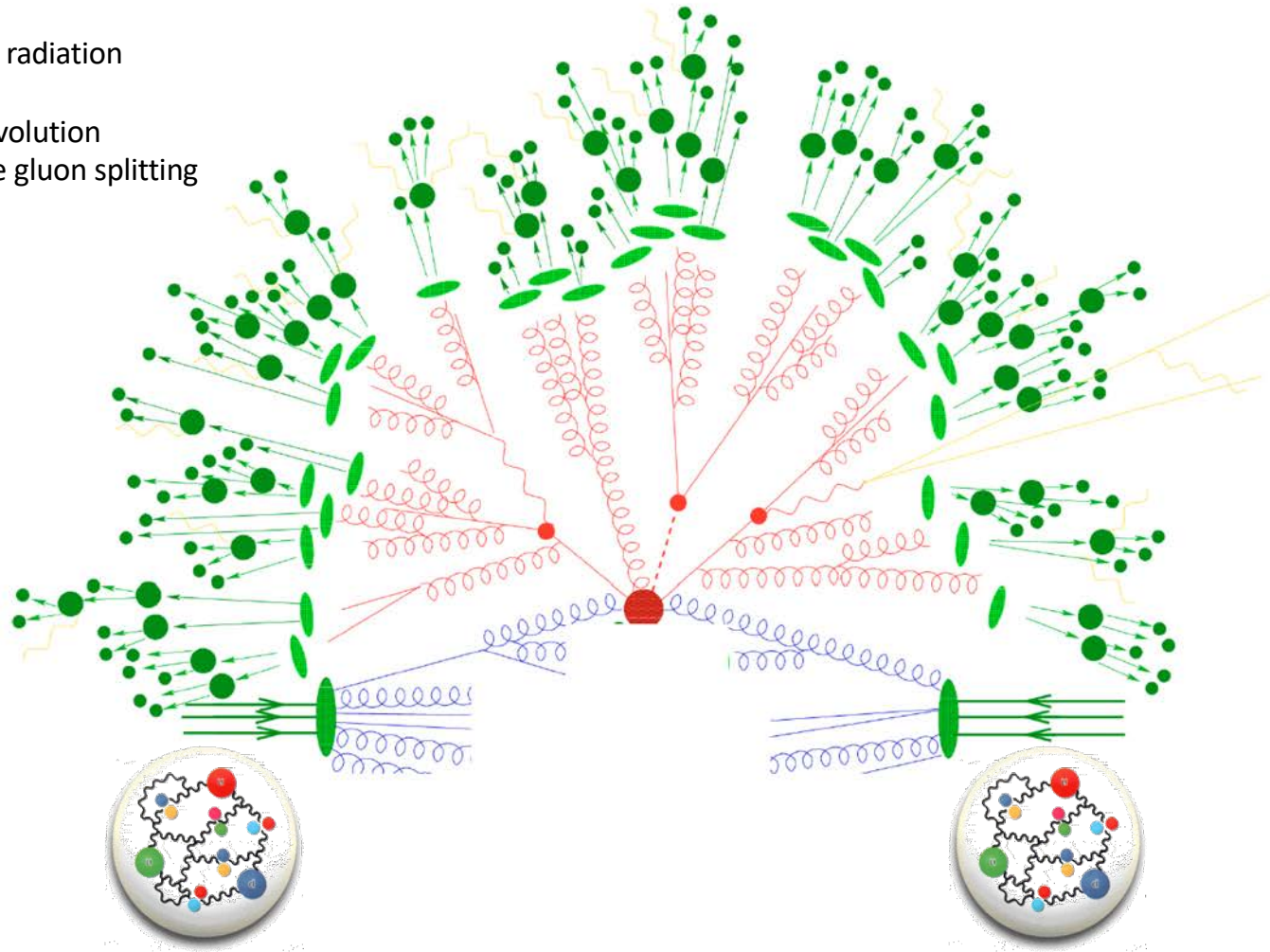
# The top quark is *beautiful and charming* but LHC is challenging

- Parton Distribution Functions (PDFs)
- Hard process
- Initial/Final state radiation



# The top quark is *beautiful and charming* but LHC is challenging

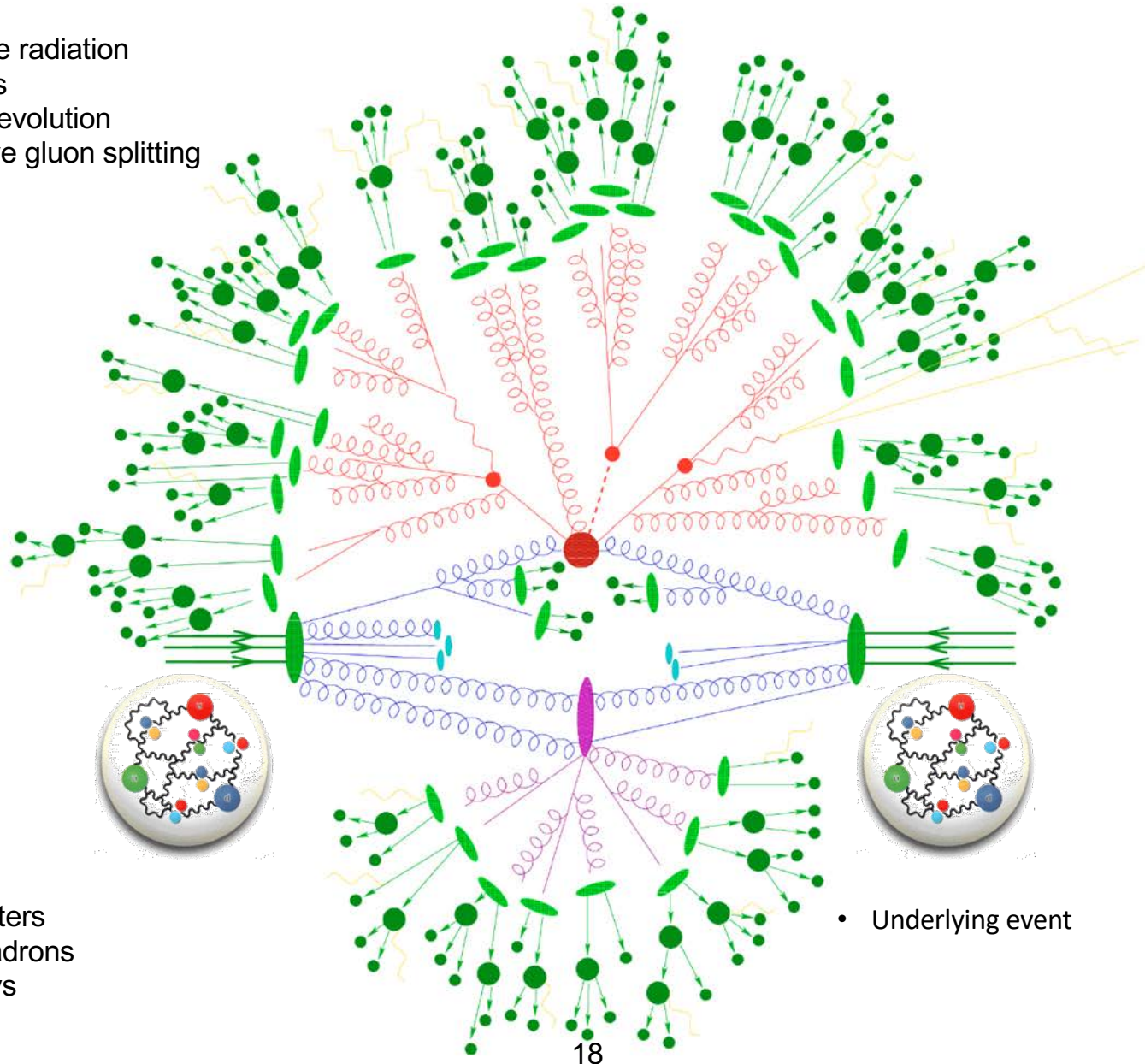
- Parton Distribution Functions (PDFs)
- Hard process
- Initial/Final state radiation
- Partonic decays
- Parton shower evolution
- Non-perturbative gluon splitting
- Colour singlets



- Colourless clusters
- Clusters into hadrons
- Hadronic decays

# The top quark is *beautiful and charming* but LHC is challenging

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• Underlying event

# The top quark is *beautiful and charming* but LHC is challenging

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# The top quark is still *strange*: the consistency of the Standard Model

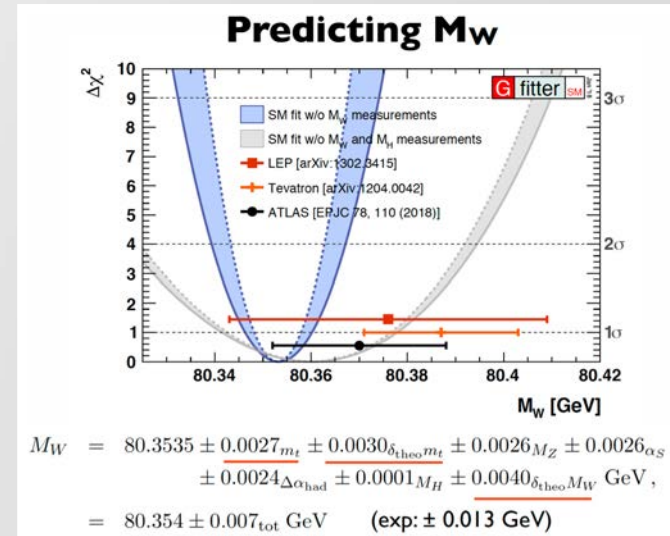
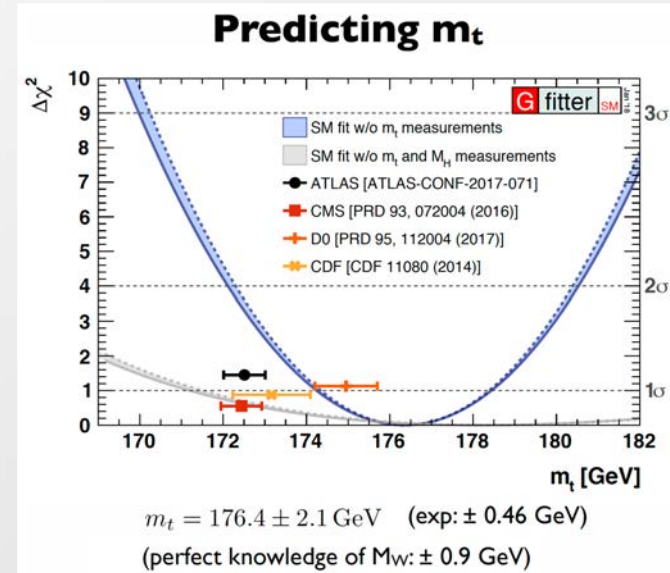
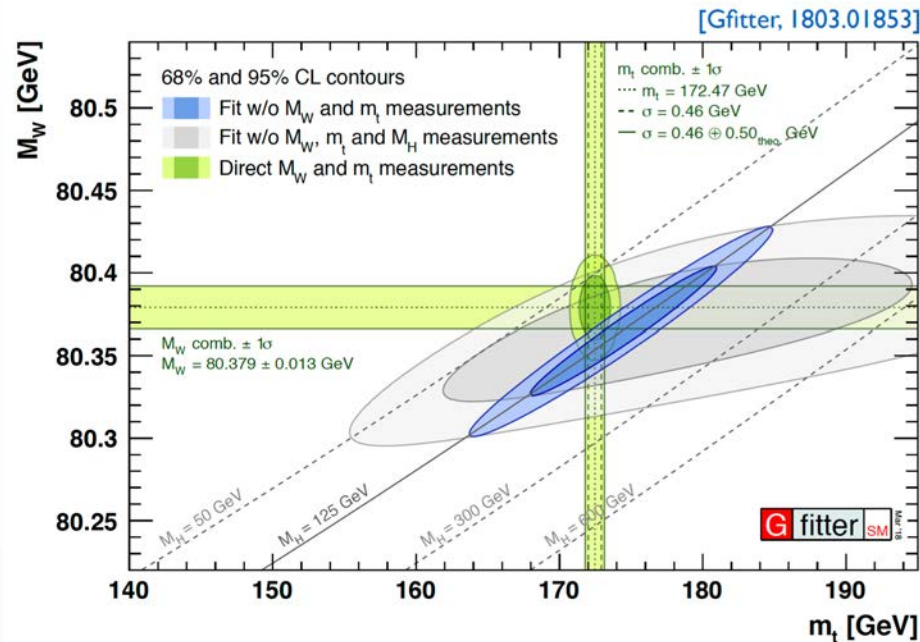
EW consistency between:  $M_W \rightleftharpoons M_H \rightleftharpoons M_t$

Roman Kogler Madrid 2019,

arXiv: 1509.00672;1708.06355;1803.01853

Gfitter group

[http://project-gfitter.web.cern.ch/project-gfitter/Standard\\_Model/](http://project-gfitter.web.cern.ch/project-gfitter/Standard_Model/)



$\delta M_W$  (indirect) = 7 MeV

$$M_W = M_W^{LO} + \Delta r_{top} + \Delta r_H$$

Large contributions to  $\delta M_W$  from top and unknown higher-order EW corrections.

$\delta M_W$  (direct) = 13 MeV

# The top quark is still *strange*: the consistency of the Standard Model

EW consistency between:  $M_W \text{ } \& \text{ } M_H \text{ } \& \text{ } M_t$

Roman Kogler Madrid 2019,

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## Experimental Input

### Fit is overconstrained

- ▶ All free parameters measured ( $\alpha_s(M_Z)$  unconstrained in fit)
  - Most input from  $e^+e^-$  colliders
    - $M_Z : 2 \cdot 10^{-5}$
  - Crucial input from Tevatron and LHC:
    - $m_t : 4 \cdot 10^{-3}$
    - $M_H : 2 \cdot 10^{-3}$
    - $M_W : 2 \cdot 10^{-4}$
  - Remarkable precision,  $O(0.1\%)$
- ▶ Require precision calculations (NNLO corrections available)

→	$M_H$ [GeV]	$125.1 \pm 0.2$	LHC
→	$M_W$ [GeV]	$80.379 \pm 0.013$	Tev.+LHC
	$\Gamma_W$ [GeV]	$2.085 \pm 0.042$	
	$M_Z$ [GeV]	$91.1875 \pm 0.0021$	LEP
	$\Gamma_Z$ [GeV]	$2.4952 \pm 0.0023$	
	$\sigma_{\text{had}}^0$ [nb]	$41.540 \pm 0.037$	SLD
	$R_\ell^0$	$20.767 \pm 0.025$	
	$A_{\text{FB}}^{0,\ell}$	$0.0171 \pm 0.0010$	Tev. (+LHC?)
	$A_\ell$ (*)	$0.1499 \pm 0.0018$	
	$\sin^2\theta_{\text{eff}}^\ell(Q_{\text{FB}})$	$0.2324 \pm 0.0012$	SLD
	$\sin^2\theta_{\text{eff}}^\ell(\text{TEV})$	$0.23148 \pm 0.00033$	
	$A_c$	$0.670 \pm 0.027$	LEP
	$A_b$	$0.923 \pm 0.020$	
	$A_{\text{FB}}^{0,c}$	$0.0707 \pm 0.0035$	LEP
	$A_{\text{FB}}^{0,b}$	$0.0992 \pm 0.0016$	
	$R_c^0$	$0.1721 \pm 0.0030$	low E
	$R_b^0$	$0.21629 \pm 0.00066$	
→	$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$	$2760 \pm 9$	Tev.+LHC
	$\bar{m}_c$ [GeV]	$1.27^{+0.07}_{-0.11}$	
	$\bar{m}_b$ [GeV]	$4.20^{+0.17}_{-0.07}$	
→	$m_t$ [GeV]( $\nabla$ )	$172.47 \pm 0.68$	

# The top quark is still *strange*: the consistency of the Standard Model

EW consistency between:  $M_W \approx M_H \approx M_t$

Roman Kogler Madrid 2019,

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- ( $\alpha_s(M_Z)$  unconstrained in fit)
- Most input from e+e-
  - $M_Z : 2 \cdot 10^{-5}$
- Crucial :
  - Tevatron
  - $m_t : 2 \cdot 10^{-5}$
  - $M_H : 2 \cdot 10^{-5}$
  - $M_W : 2 \cdot 10^{-5}$
- Remarkable precision,  $O(0.1\%)$
- ▶ Require precision calculations (NNLO corrections available)

e+e- and polarization, many years later, still constraint the Standard Model to

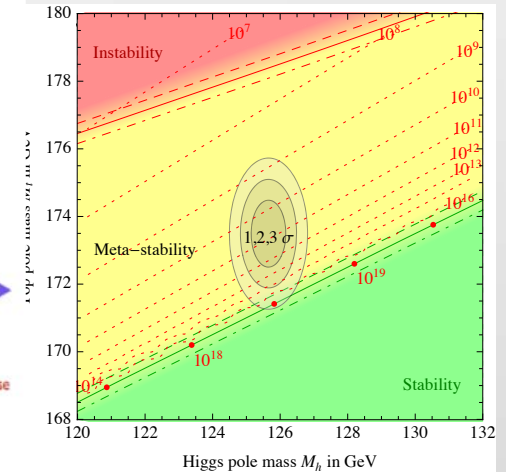
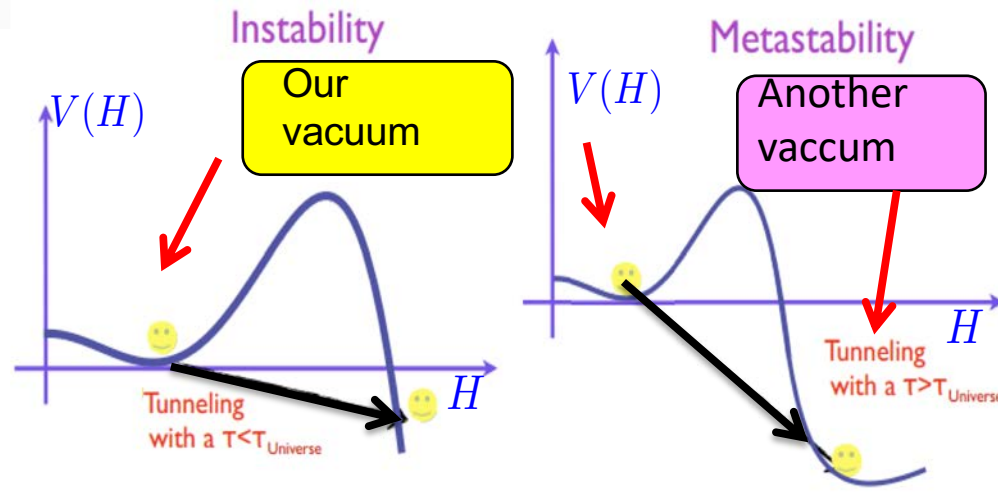
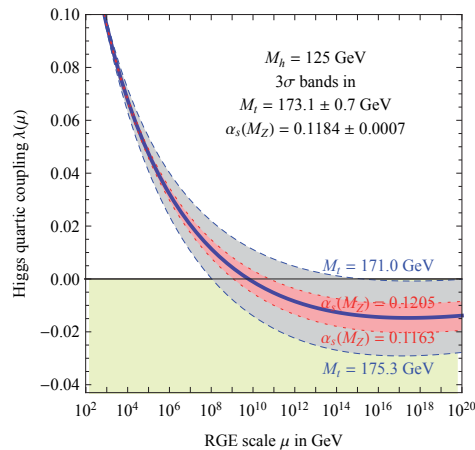
$M_H$ [GeV]		LHC
$M_t$ [GeV]		Tev.+LHC
$\alpha_s(M_Z)$	$0.1181 \pm 0.0023$	
$\sin^2 \theta_{\text{eff}}^{\ell}(Q_{\text{FB}})$	$0.2312 \pm 0.0037$	LEP
$\sin^2 \theta_{\text{eff}}^{\ell}(\text{TEV})$	$0.23148 \pm 0.00033$	SLD
$A_c$	$0.670 \pm 0.027$	Tev. (+LHC?)
$A_b$	$0.923 \pm 0.020$	SLD
$A_{\text{FB}}^{0,c}$	$0.0707 \pm 0.0035$	LEP
$A_{\text{FB}}^{0,b}$	$0.0992 \pm 0.0016$	LEP
$R_c^0$	$0.1721 \pm 0.0030$	
$R_b^0$	$0.21629 \pm 0.00066$	
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$	$2760 \pm 9$	low E
$\bar{m}_c$ [GeV]	$1.27^{+0.07}_{-0.11}$	
$\bar{m}_b$ [GeV]	$4.20^{+0.17}_{-0.07}$	
$m_t$ [GeV]( $\nabla$ )	$172.47 \pm 0.68$	Tev.+LHC

# The top quark is still *strange*: the vacuum stability

Vacuum Stability ( $\lambda(\Lambda) \geq 0$ )

$\lambda(\Lambda)$  the  $\overline{\text{MS}}$  quartic Higgs Coupling

Degrassi et al, *JHEP* **1208 (2012) 098**  
Butazzo et al, *1307.3536 (2013)*

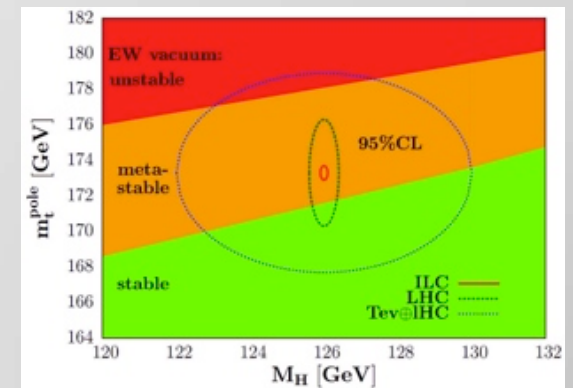


Alekhin et al, *Phys.Lett.* **B716 (2012) 214**

Need to measure  $m_t$  with very high accuracy:

$$\Delta m_t < 100\text{-}150 \text{ MeV}$$

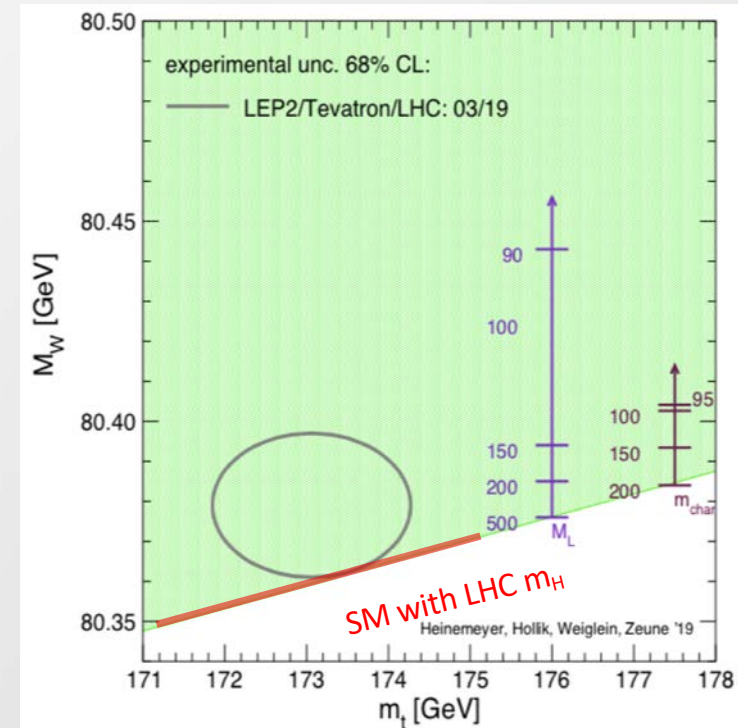
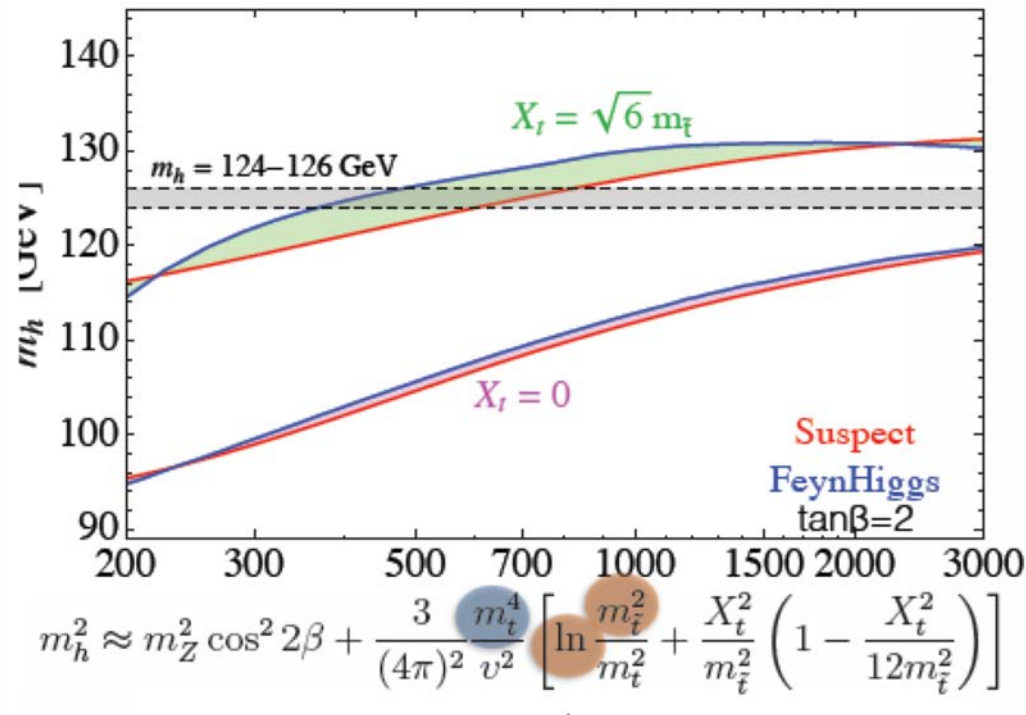
(The existence of New Physics would change the scenario)





# The top quark is still strange: sensitive to New Physics

New Physics, general arguments:



Roberto Franceschini (IFIC seminar, Valencia)

[[www.ifca.unican.es/users/heinemey/uni/plots](http://www.ifca.unican.es/users/heinemey/uni/plots)]

Large mass  $\longleftrightarrow$  Sizeable effects

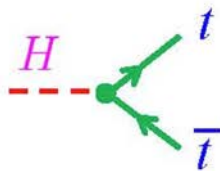
# The ups and downs of the top-quark mass: the mass in the Lagrangian

$$\mathcal{L}_{QCD} = -\frac{1}{4}(\partial^\mu G_a^\nu - \partial^\nu G_a^\mu)(\partial_\mu G_a^\nu - \partial_\nu G_a^\mu) + \sum_f \bar{q}_f^\alpha (i\gamma^\mu \partial_\mu - m_f) q_f^\alpha$$

$$- g_s G_a^\mu \sum_f \bar{q}_f^\alpha \gamma_\mu \left(\frac{\lambda^a}{2}\right)_{\alpha\beta} q_f^\beta - \frac{g_s}{4} f^{abc} (\partial^\mu G_a^\nu - \partial^\nu G_a^\mu) G_\mu^b G_\nu^c - \frac{g_s^2}{4} f^{abc} f^{ade} G_b^\mu G_c^\nu G_\mu^d G_\nu^e$$

$\alpha_s = g_s^2/4\pi$  and quark masses are not predicted by the SM

$m_t$  { Fundamental parameter of the SM interesting “per se”  
 Important for precise tests of the Standard Model, Yukawa coupling  $\sim 1$   
 Test of New Physics scenarios i.e. GUT scenarios, vacuum stability



$$y_t = \frac{\sqrt{2}}{v} m_t = 2^{3/4} G_F^{1/2} m_t = 1 \quad (0.995)$$

## The *ups and downs* of the top-quark mass: mass of quarks

---

- Free quarks are not observed in nature as they are confined into colourless hadrons, so there is no pole in the S-matrix
  - ✓ quark-masses, in particular the top-quark mass, are not “observables” and they are parameters of the underlying theory
    - fit  $O^{\text{exp}}(x)$  with  $O^{\text{th}}(m_t, \alpha_s; x)$  and extract  $m_t$  ←
  - ✓ precise value depends on the definition of the renormalization scheme selected (pole mass,  $\overline{\text{MS}}$ , MSR, etc..)
  - ✓ to fix the renormalization scheme at least a NLO calculation is required
  - ✓ In a way, “quark masses” are kind of “effective coupling constants”

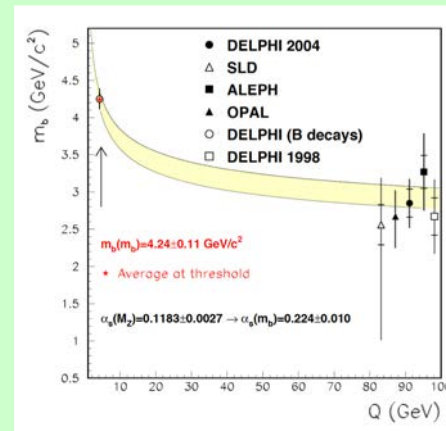
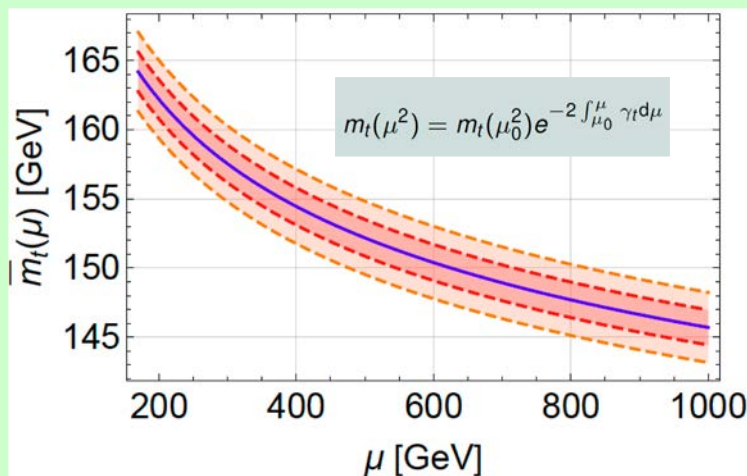
**Corollary:** the mass of quarks cannot be measured as the  $\pi^0$  mass

# The ups and downs of the top-quark mass: mass definitions

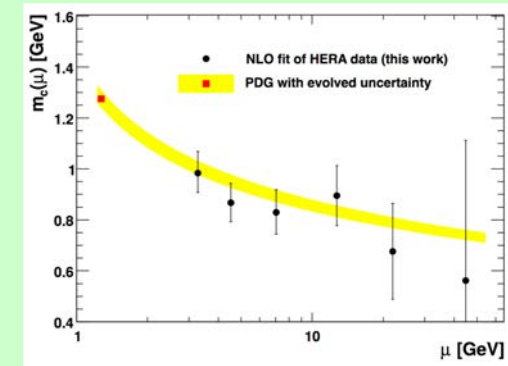
Most common definitions:

$\overline{m}_t(\mu)$ : the SM  $m_t$  renormalised in the  $\overline{\text{MS}}$  /  $\text{MS}$  scheme,

- Well defined order-by-order in pQCD
- It is function of an energy scale parameter  
(but a scale-independent number can be associated to it:  $m_t(m_t)$ )
- Good description of physics far from top-quark production thresholds
- Mass running typically of the form



A. Gizhko et al., *Phys.Lett. B775* (2017) 233-238



Quark masses run !!

# The ups and downs of the top-quark mass: mass definitions

Most common definitions:

$m_t(\mu)$ : the SM  $m_t$  renormalised in the  $\overline{\text{MS}}$  /  $\text{MS}$  scheme,

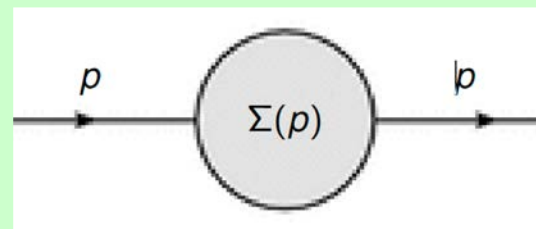
$m_t^{\text{pole}}$ : the mass renormalised in the pole-mass scheme,

- Defined as the pole of the top-quark propagator
- Intrinsic uncertainty in definition (200 MeV renormalon problem)
- Good description of physics near top-quark production thresholds.
- If the top-quark mass would be a stable particle,  $m_t^{\text{pole}}$  would be its physical mass

$$\Pi(p^2) \propto \frac{1}{p^2 - m^2 + \Sigma(p^2) + i0}$$

renormalisation conditions:

$$\Sigma(p^2 = m^2) = 0, \quad \left[ \frac{d\Sigma}{dp^2} \right]_{p^2=m^2} = 0.$$



- $m_t(\mu)$  and  $m_t^{\text{pole}}$  can be related using pQCD

Pole mass  
vs  
running mass

$$m_t = \bar{m}(\mu) \left( 1 + \frac{\alpha_s(\mu)}{\pi} \left[ \frac{4}{3} + \ln \left( \frac{\mu^2}{\bar{m}(\mu)^2} \right) \right] + O(\alpha_s^2) \right)$$

# The ups and downs of the top-quark mass: mass definitions

Most common definitions:

- $m_t(\mu)$ : the SM  $m_t$  renormalised in the  $\overline{\text{MS}}$  /  $\text{MS}$  scheme,
- $m_t^{\text{pole}}$ : the mass renormalised in the pole-mass scheme,
- $m_t^{\text{MSR}}(\mathbf{R})$ : short-distance mass that smoothly interpolates all R scales

- Renormalon free
- Precision in relation to any other short-distance mass:  $\lesssim 20 \text{ MeV} @ O(\alpha_s^4)$

## MSR Mass Definition

Jain, Lepenik, Mateu, Preisser,  
Scimemi, Stewart, AHH  
arXiv:1704.01580

MS Scheme: ( $\mu > \overline{m}(\overline{m})$ )

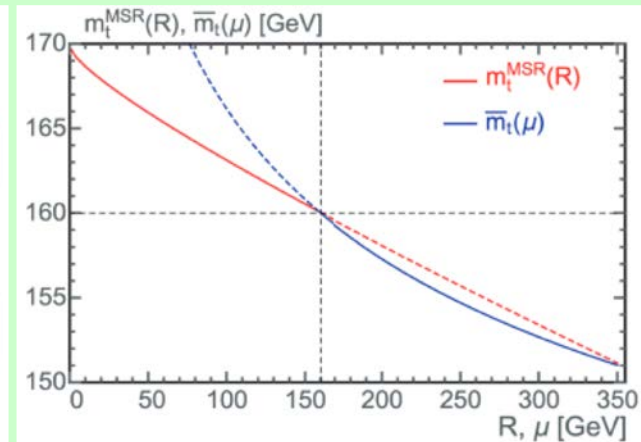
$$\overline{m}(\overline{m}) - m^{\text{pole}} = -\overline{m}(\overline{m}) [0.42441 \alpha_s(\overline{m}) + 0.8345 \alpha_s^2(\overline{m}) + 2.368 \alpha_s^3(\overline{m}) + \dots]$$

MSR Scheme: ( $R < \overline{m}(\overline{m})$ )



$$m_{\text{MSR}}(R) - m^{\text{pole}} = -R [0.42441 \alpha_s(R) + 0.8345 \alpha_s^2(R) + 2.368 \alpha_s^3(R) + \dots]$$

$$m_{\text{MSR}}(m_{\text{MSR}}) = \overline{m}(\overline{m})$$



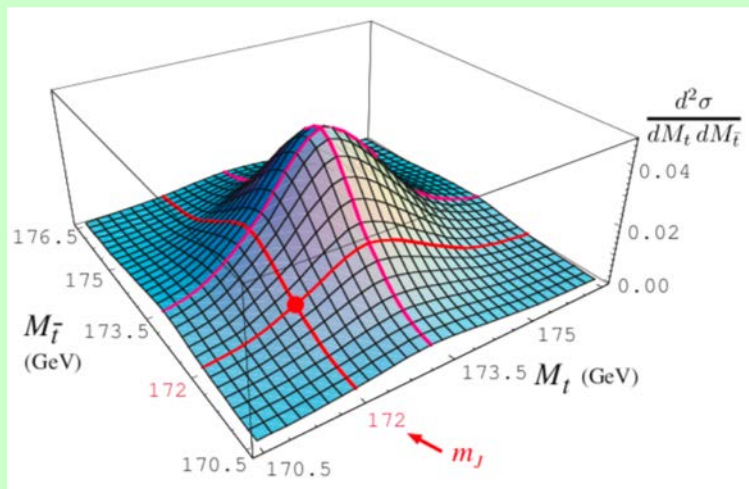
A. Hoang et al., The MSR and the  $O(\Lambda_{\text{QCD}})$  renormalon sum rule, **JHEP 1804 (2018) 003**

# The ups and downs of the top-quark mass: mass definitions

Most common definitions:

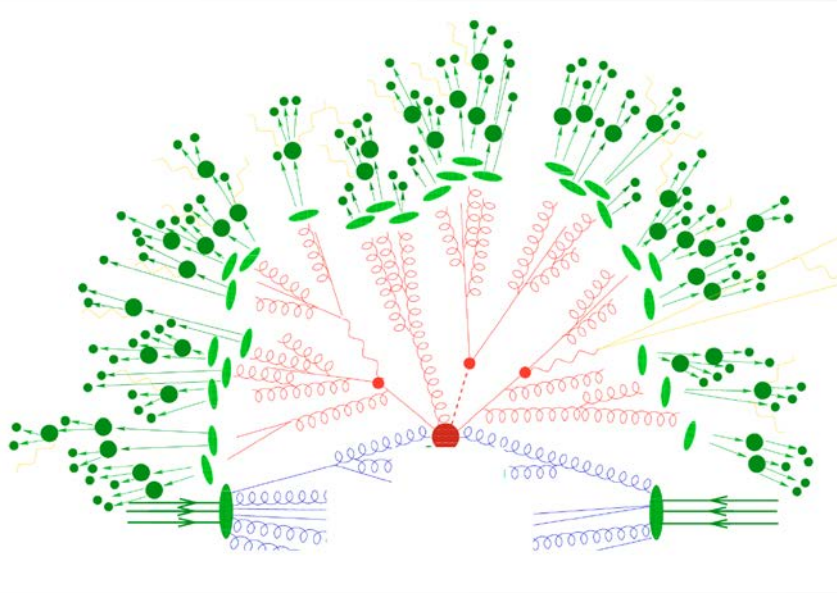
- $m_t(\mu)$ : the SM  $m_t$  renormalized in the  $\overline{\text{MS}}$  /  $\text{MS}$  scheme,
- $m_t^{\text{pole}}$ : the mass renormalized in the pole-mass scheme,
- $m_t^{\text{MSR}}(\mathbf{R})$ : short-distance mass that smoothly interpolates all R scales
- $m_t^{\text{MC}}$ : top mass value as implemented in the Monte Carlo generators

- As an effective parameter used in hard process+parton shower+hadronization
- Not related with any top mass definition of the SM Lagrangian (EW fits, etc..)



Example from Phys. Rev. D77 (2008) 074010  
pQCD calculation generated with  
 $m_j = 172 \text{ GeV}$   
extracted value after shower/hadronization  
 $M_t = 173.5 \text{ GeV}$

# The ups and downs of the top-quark mass: Monte Carlo event generators



Monte Carlo event generators are used to describe these processes:

- Matrix elements (LO/NLO)
- Parton shower (LL)
- Hadronization model

- Experimental observables fully simulated
- QCD-inspired: partly first principles QCD and partly modelled (getting very sophisticated though)
- Modelling parameters are tuned to data getting better accuracy than intrinsic theory
- Top quark in parton shower is treated like a real particle ( $m_t^{\text{MC}} \approx m_t^{\text{pole}} + ?$ )
- Top quark in matrix elements:  $m_t^{\text{MC}} = m_t^{\text{pole}}$

BUT:

- parton showers sum (real & virtual !) perturbative corrections only above the shower cut and not pickup any corrections from below.
- what is the meaning of MC parameters ? (calibration & theory)



# The ups and downs of the top-quark mass: Monte Carlo event generators

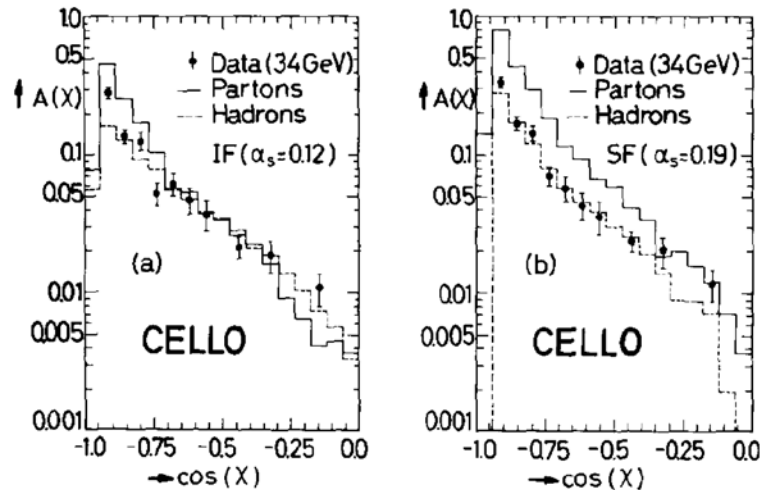


Fig. 1. (a) Corrected asymmetry data compared with the asymmetry of the partons and the generated final state hadrons in the independent fragmentation model (IF). (b) As (a), but now for the string model.

## Already an old story 1984

method	SF	IF0	IF	IF1	SF/IF
asymmetry	0.19	0.12	0.12	0.15	1.58
cluster thrust	0.18	0.13	0.12	0.13	1.50

On the model dependence of the determination of the strong coupling constant in second order QCD from  $e^+e^-$  annihilation into hadrons

CELO Collab., H-J Behrend et al.

**Phys. Lett. 138B (1984) 311-316**

The value of  $\alpha_s$  was shown to have 50% uncertainty due “MC model dependence”. At this time this was Independent Fragmentation vs String fragmentation.

**A huge progress since then in theory and experiment but still modelling in MC is harming !!!**

# The ups and downs of the top-quark mass: Mass definition

- Studies to calibrate/relate  $m_t^{\text{MC}}$ :

- J. Kieseler et al., Phys. Lett, 116 (2016) 162001

$$\Delta_m = m_t - m_t^{\text{MC}} \sim 2 \text{ GeV}$$

This study uses tt cross-sections

- M. Butenschoen et al., Phys.Rev.Lett. 117 (2016) no.23, 232001.

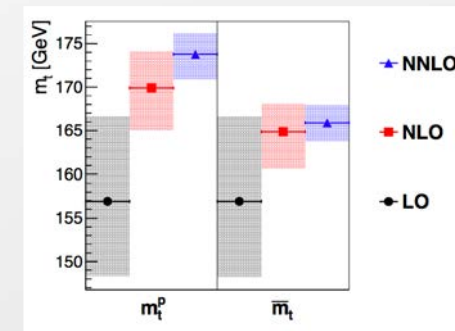
$$m_t^{\text{MC}} = m_t^{\text{MSR}}(1 \text{ GeV}) + (0.18 \pm 0.22) \text{ GeV} \quad m_t^{\text{MC}} = m_t^{\text{pole}} + (0.57 \pm 0.28) \text{ GeV}$$

This study uses 2-jettines distribution in e+e- interactions.

**Expanding the result to pp interactions and different observables remains to be proven**

- Other estimates in literature  $\Delta_m \sim O(0.5-1.0) \text{ GeV}$
- A lot of discussion/controversy on how to interpret these results/differences:

- S. Moch et al., arXiv:1405.4781,
- A. H. Hoang and I. W. Stewart, 500 Nuovo Cimento B123 (2008) 1092–1100,
- A. Buckley et al., arXiv:1101.2599,
- A. H. Hoang, arXiv:1412.3649,
- A. Hoang et al., JHEP 1810 (2018) 200,
- A. Hoang et al, JHEP 1804 (2018) 003,
- M. Dasgupta et al., JHEP 09 (2013) 029,
- A. J. Larkoski et al., JHEP 05 (2014) 146,
- P. Nason, arXiv:1712.02796, arXiv:1901.04737,
- G. Corcella, arXiv:1903.06574,
- S. Ferrairo et al., arXiv:1906.09166,
- S. Ferrairo et al., JHEP 1901 (2019) 203



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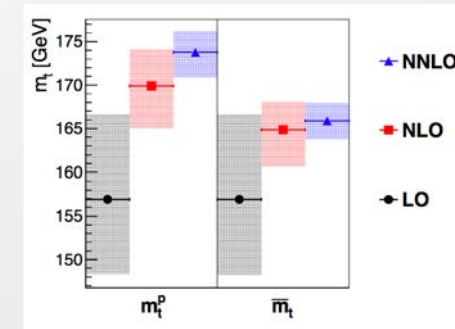
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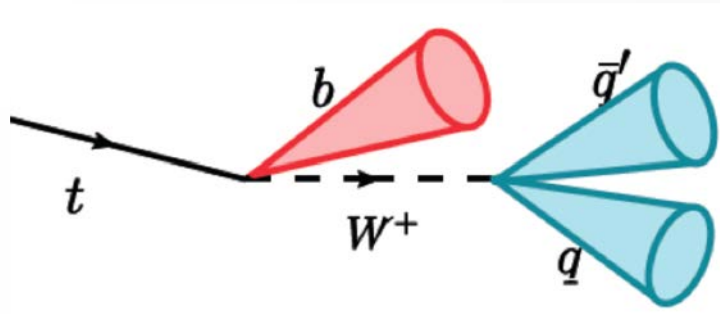
O(0.5-1.0) GeV represent 0.3%-0.6% relative effect to the top mass

Very small but significant given present accuracy

***The prize of being so precise !!!!***

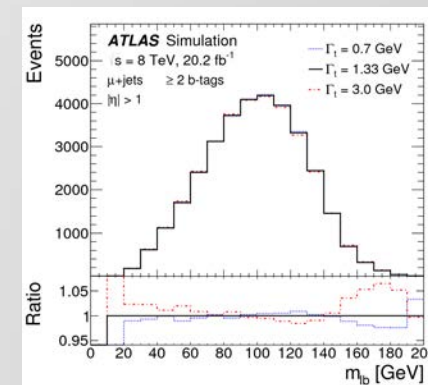
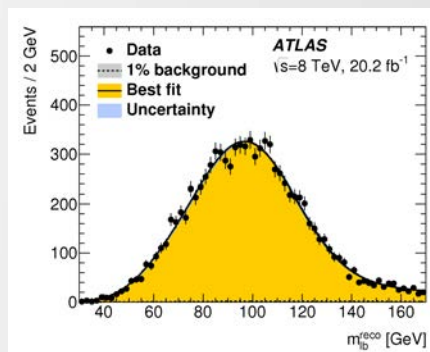
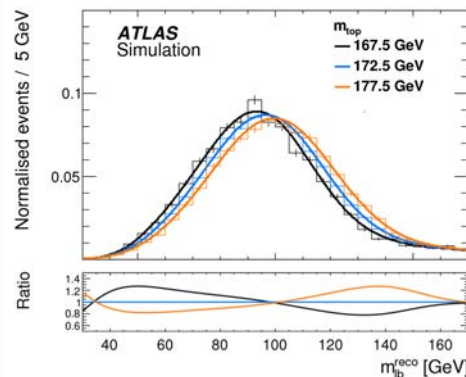
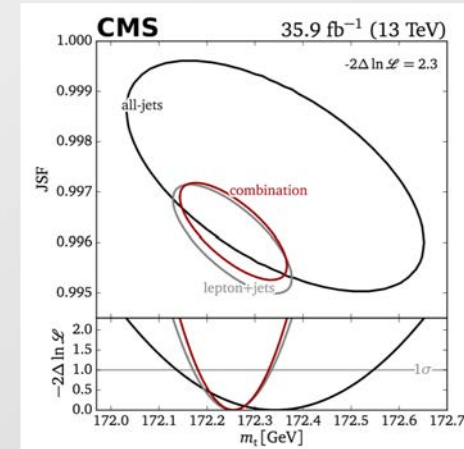
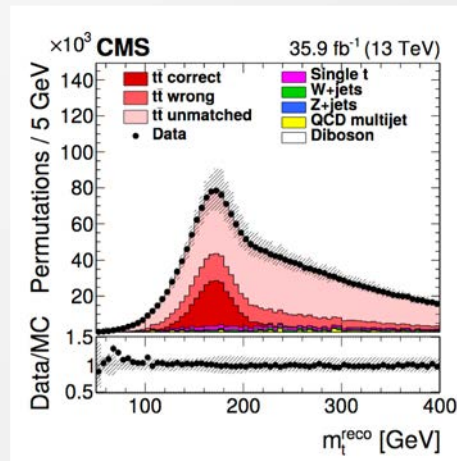
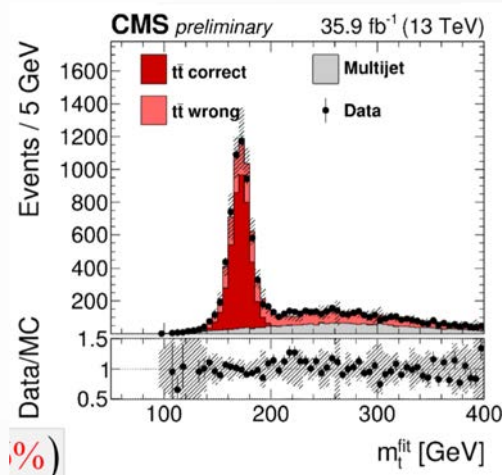
- A. Buckley et al., arXiv:1101.2599,
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# The ups and downs of the top-quark mass: experimental determinations

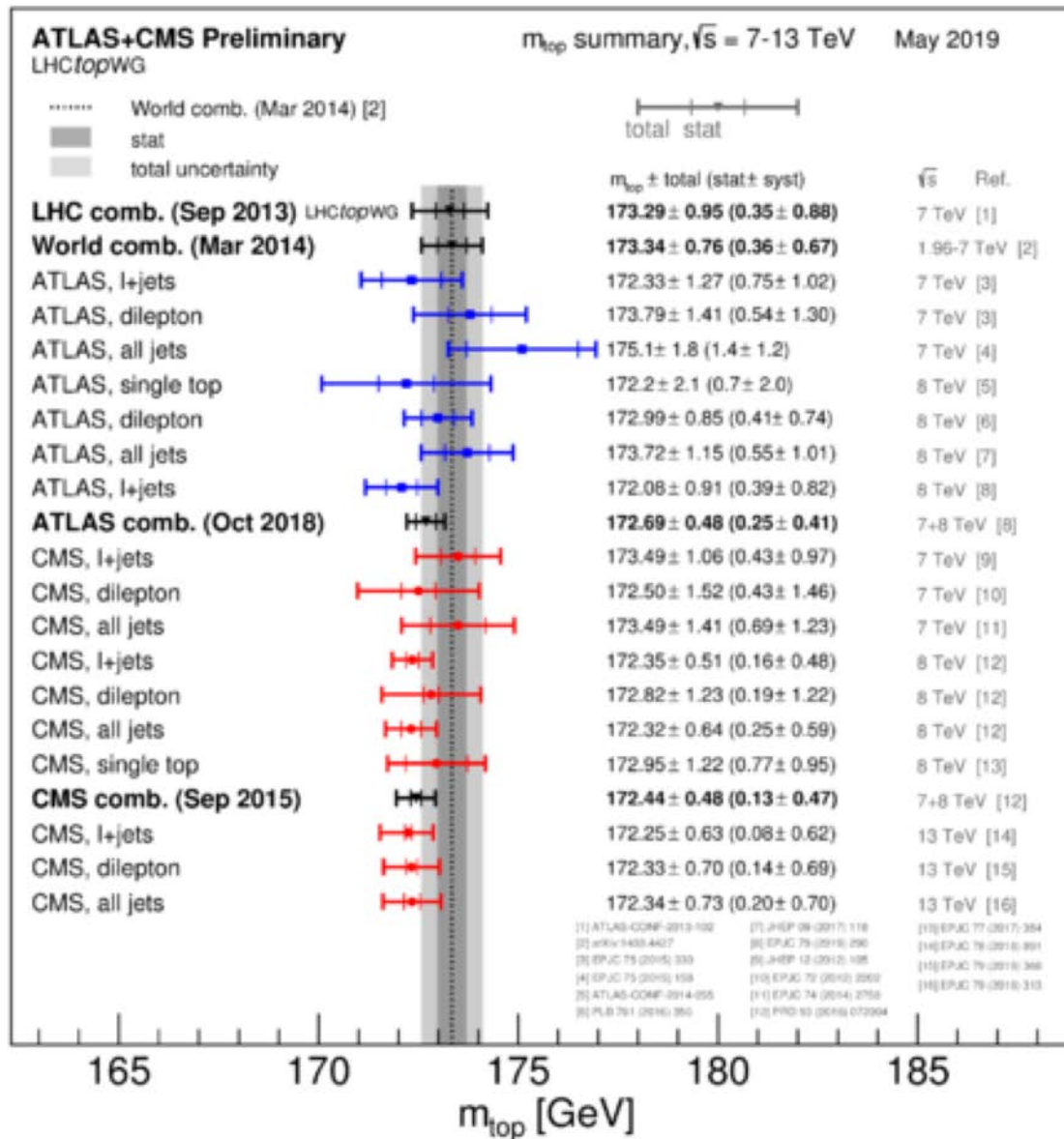


## “Direct” mass measurements

- Reconstruction of top-decay products
- High top mass sensitivity
- Template, matrix-element and ideogram methods
- Determine  $m_t^{MC}$



# The ups and downs of the top-quark mass: experimental determinations



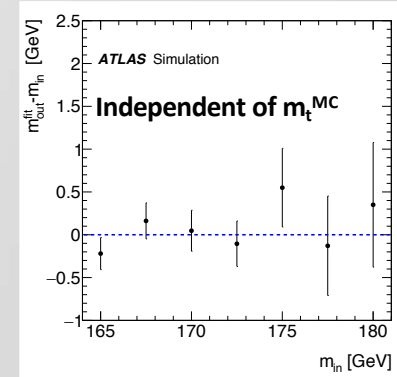
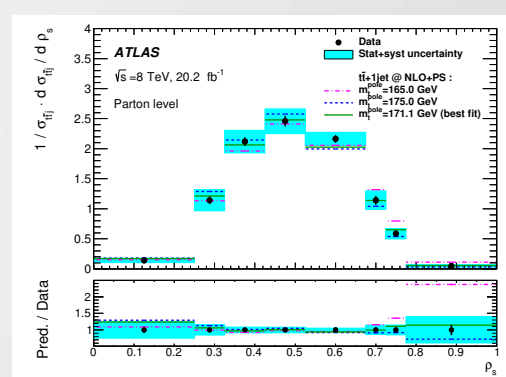
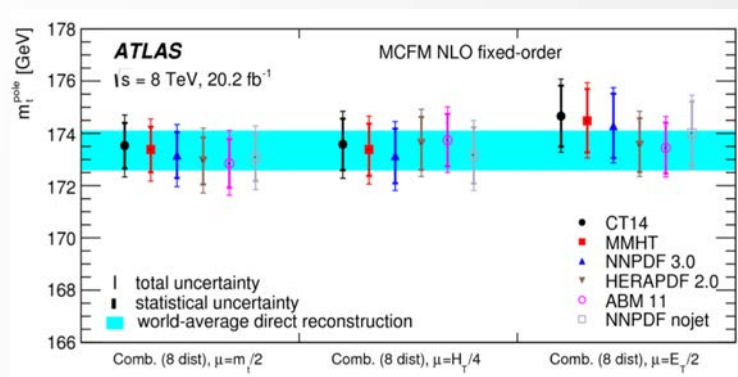
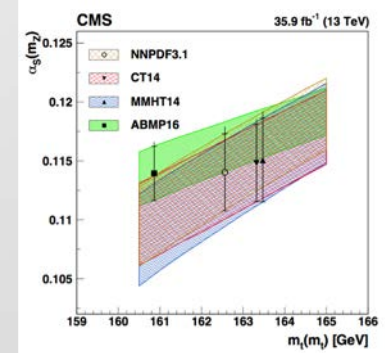
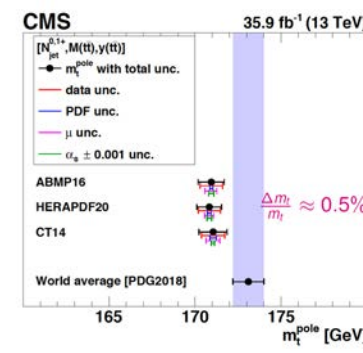
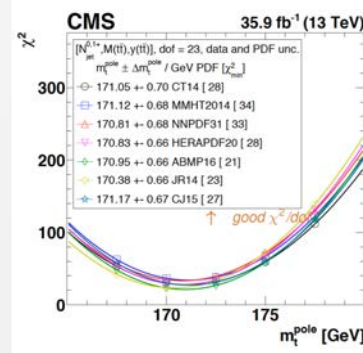
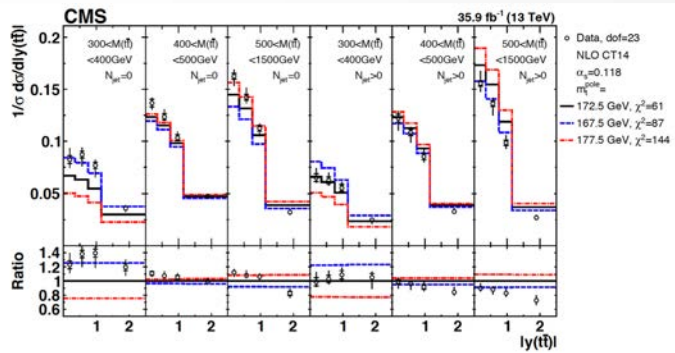
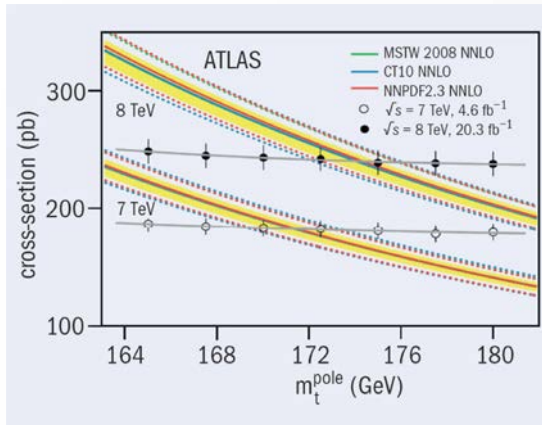
## LHC Summary

## “Direct” measurements

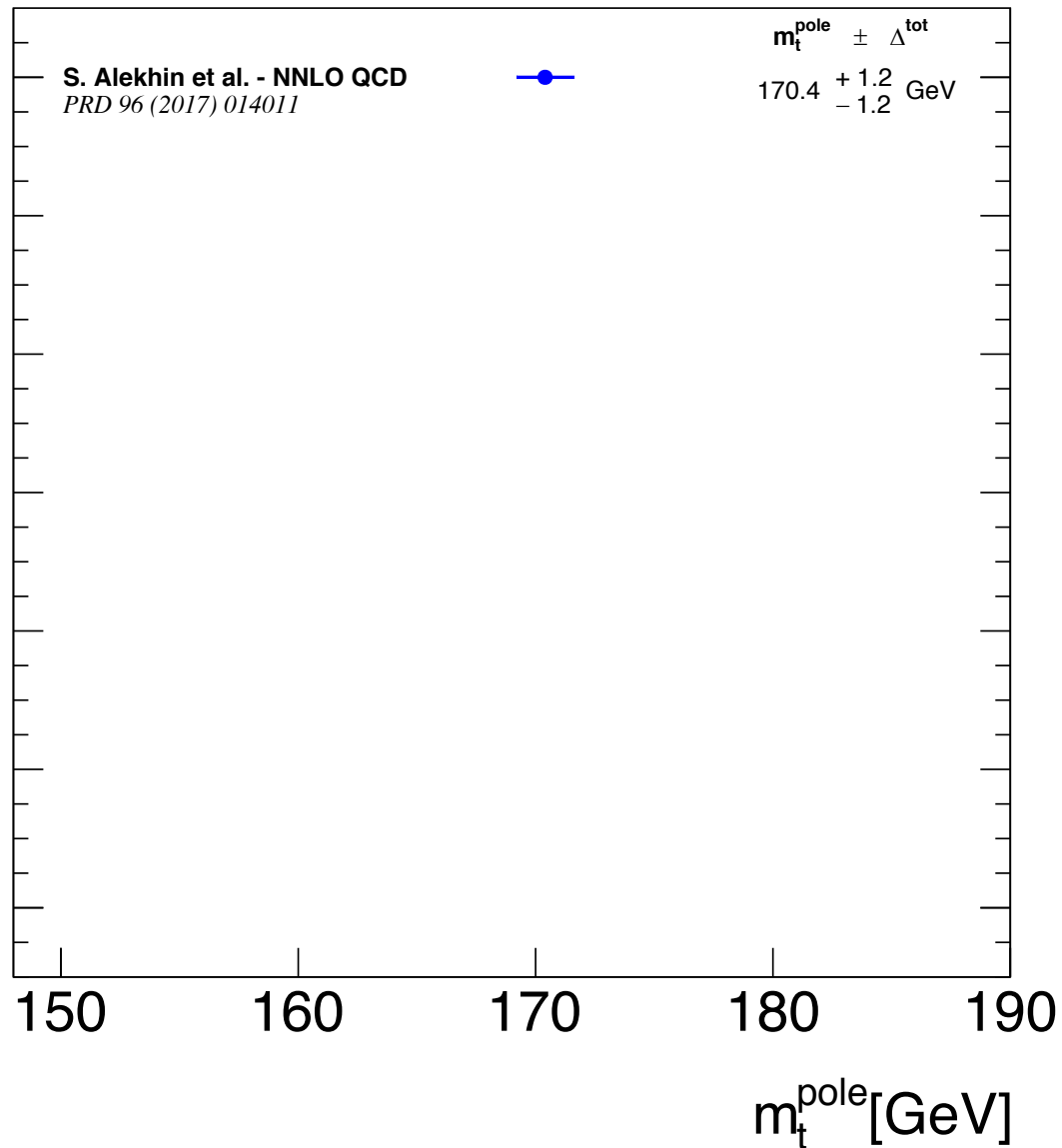
# The ups and downs of the top-quark mass: **experimental determinations**

## “Alternative” mass measurements

- Use total or differential cross-sections
- Compare corrected data to pQCD prediction (NLO/NNLO)
- Determine  $m_t^{\text{pole}}$  and  $m_t(\mu)$
- Usually less sensitivity than “direct” measurements
- But., entering to the sub-GeV accuracy



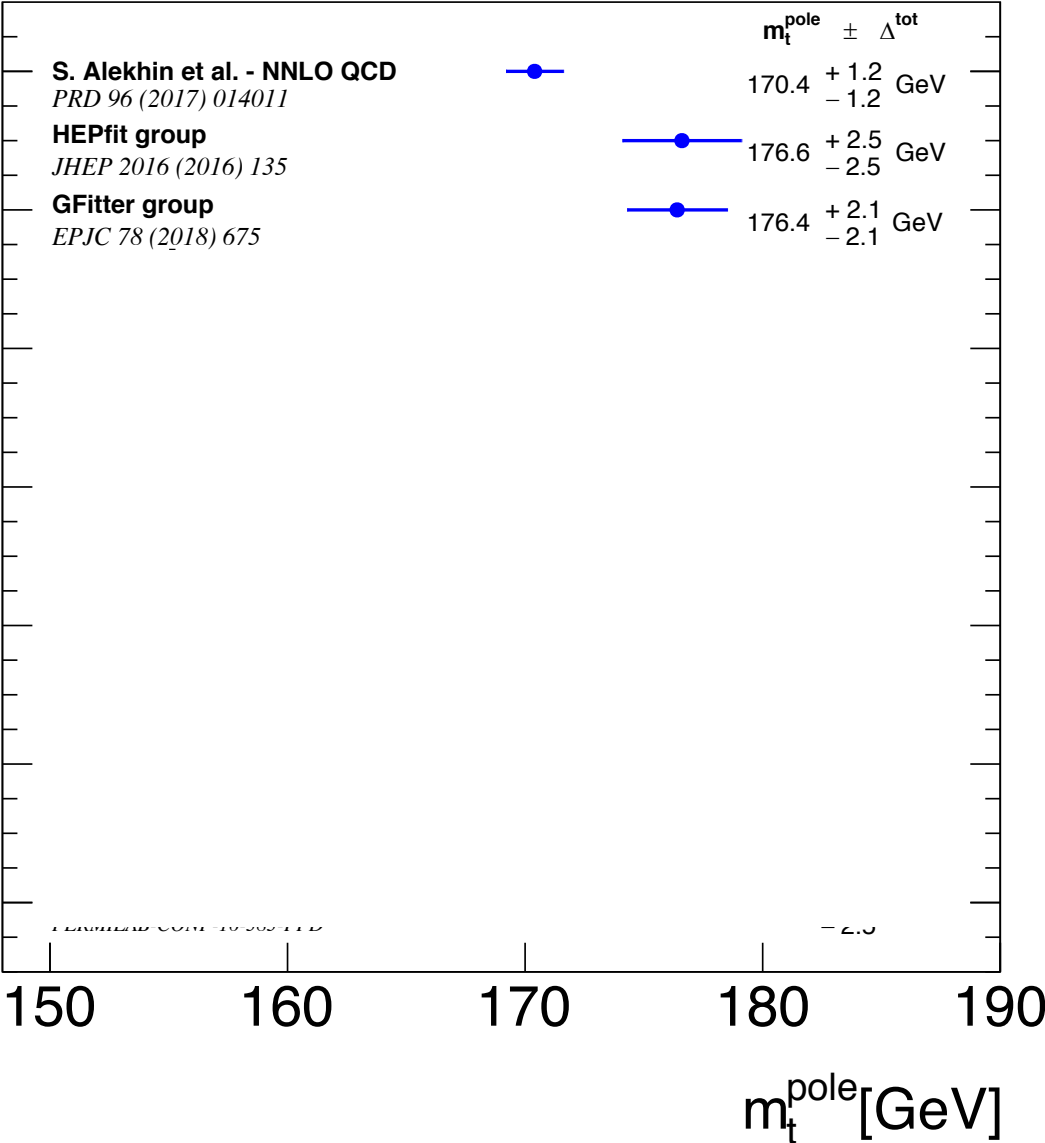
# The ups and downs of the top-quark mass: **experimental determinations**



## Pole mass summary

**Global Fit – NNLO QCD**

# The ups and downs of the top-quark mass: experimental determinations



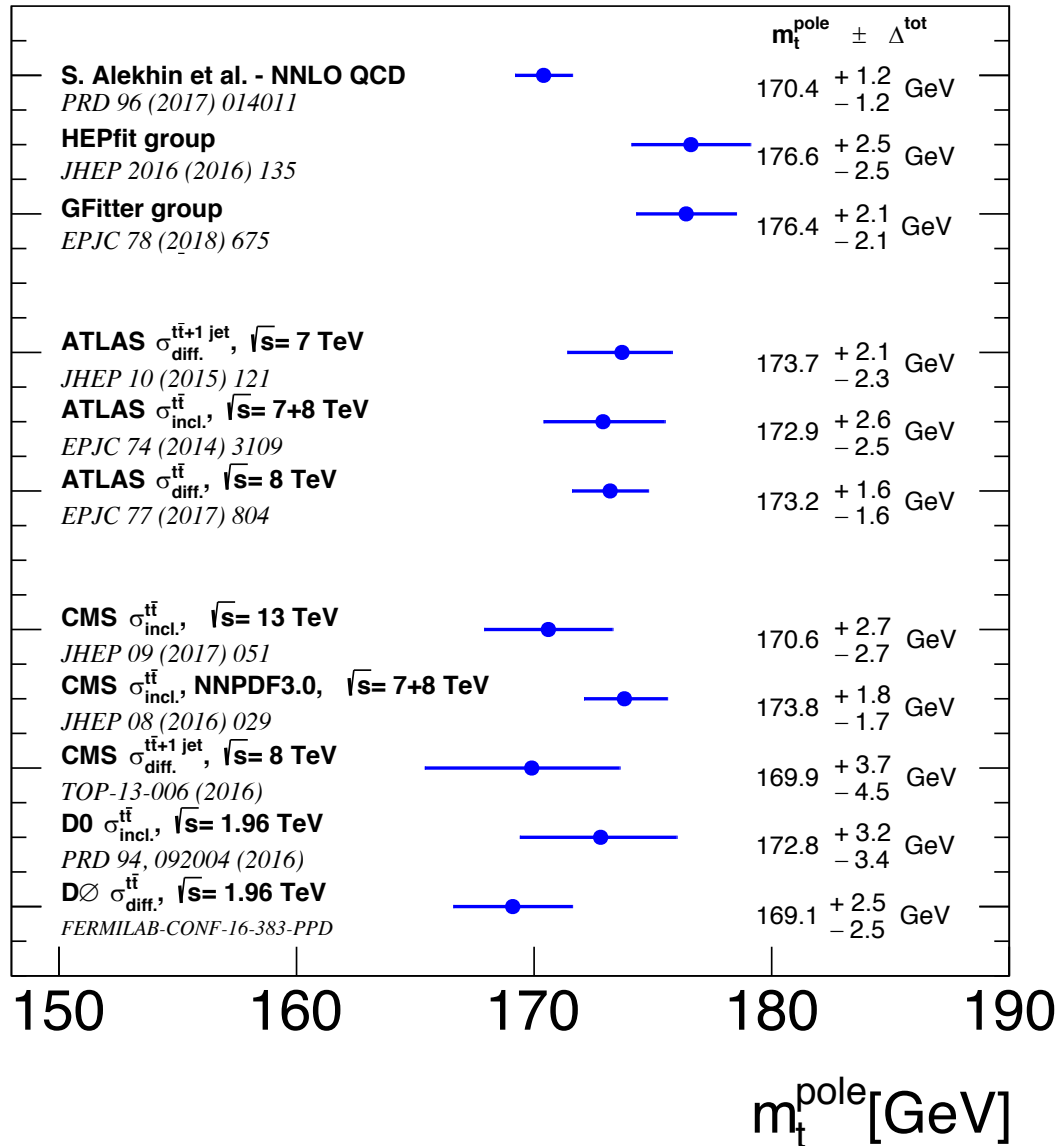
## Pole mass summary

Global Fit – NNLO QCD

EW- Fits: HEPfit. & GFitter



# The ups and downs of the top-quark mass: experimental determinations



## Pole mass summary

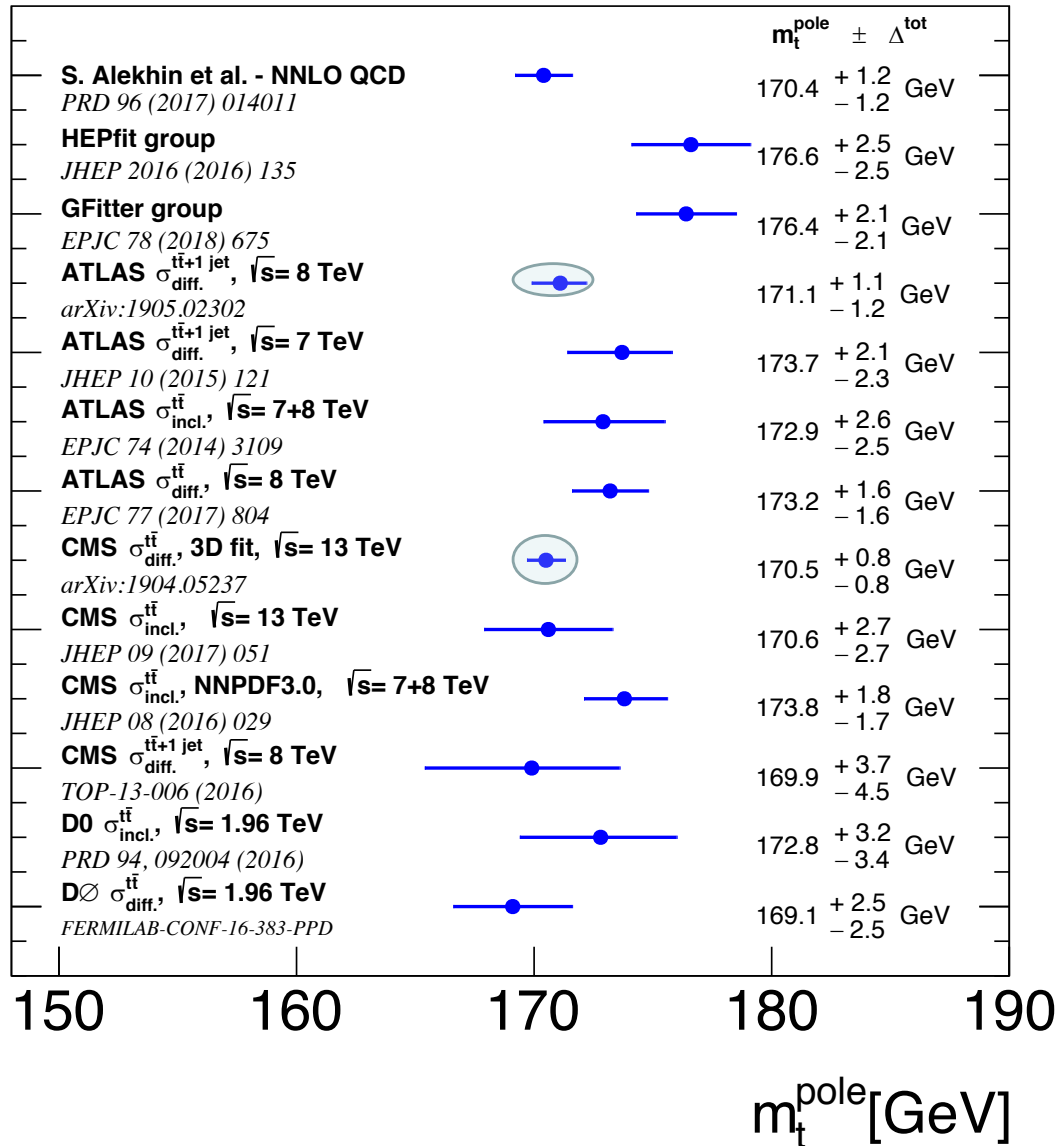
### Global Fit – NNLO QCD

### EW- Fits: HEPfit. & Gfitter

### LHC & Tevatron:

- Cross-section & differential

# The ups and downs of the top-quark mass: experimental determinations



## Pole mass summary

### Global Fit – NNLO QCD

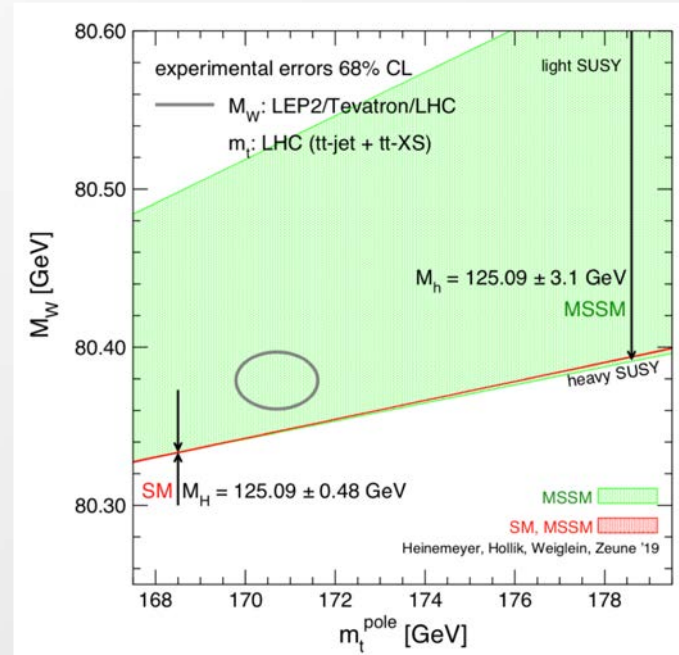
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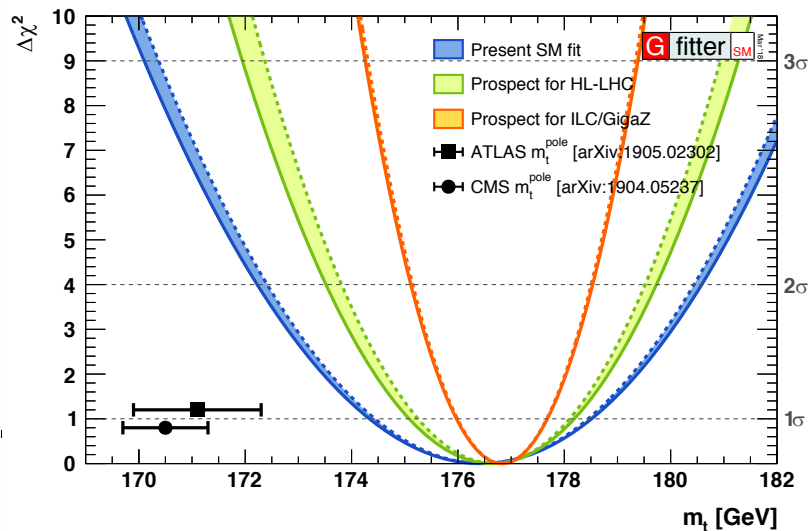
- Cross-section & differential

### LHC latest results sub-GeV !!!

# The ups and downs of the top-quark mass: **Look into the new results**



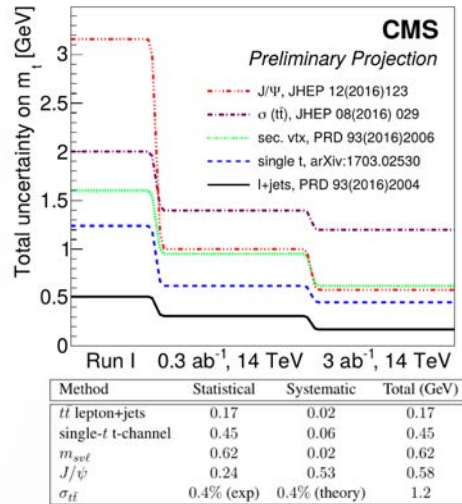
Courtesy of Sven Heinemeyer  
(Using latest ATLAS+CMS results)  
**(naïve combination)**



Courtesy of Roman Kogler - Gfitter  
(Using latest ATLAS+CMS results)  
**(including HL-LHC and ILC projections)**

**Nothing to claim but getting interesting**

# The ups and downs of the top-quark mass: looking into the future

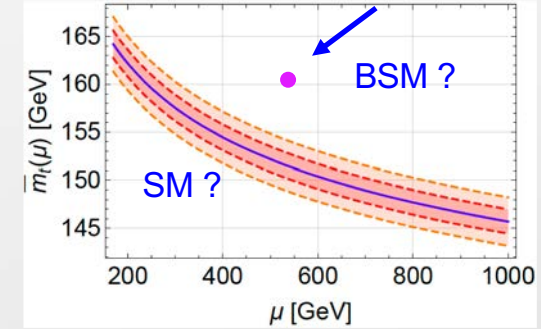


Standard Model Physics at the HL-LHC and HE-LHC, [arXiv:1902.04070](https://arxiv.org/abs/1902.04070)

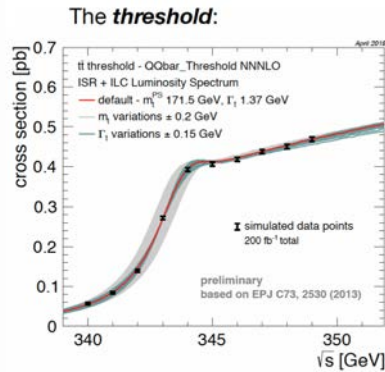
Z.Z. Xing et al. et al, *Phys.Rev. D72 (2008) 113016*

## HL-LHC (higher statistics):

- Possibility to use rare decays ( $J/\Psi$ )
- Restrict phase space regions
- Better control of systematics/modelling
- Expected accuracy: 200-300 MeV



- Need to develop further present theory calculations/predictions
- New observables or/and use of different mass definitions. To be explored



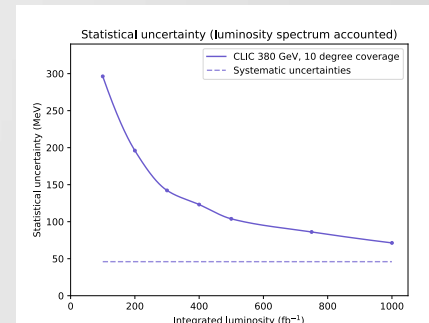
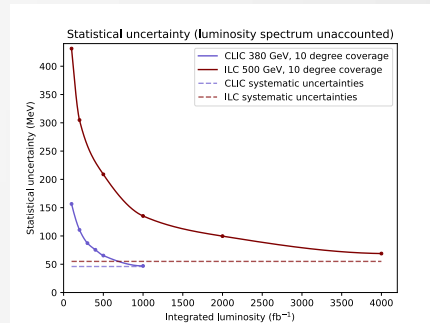
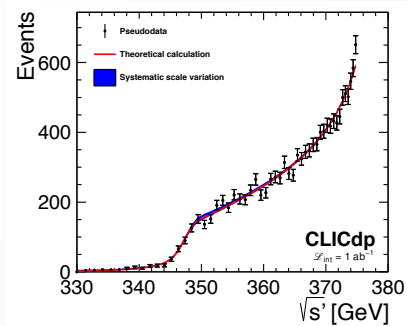
## Collider e+e- (at top threshold):

H. Abramowicz et al., CLICdp Collab., [arXiv:1807.02441](https://arxiv.org/abs/1807.02441)

- Well-defined mass scheme
- Access to top-width and Yukawa coupling
- Expected accuracy:  $m_t \sim 40-75$  MeV;  $\Gamma_t \sim 100$  MeV;  $y_t \sim 15\%$

## Collider e+e- (at continuum above top threshold):

- Well-defined mass scheme
- After 1-2 years data taking better accuracy than LHC/HL-LHC complete programme
- Expected accuracy:  $m_t \sim 100-150$  MeV



$$\sigma(e^+e^- \rightarrow t\bar{t}) = f(s, m_t) \quad \rightarrow \quad \sigma(e^+e^- \rightarrow t\bar{t}\gamma) = f(s', m_t)$$

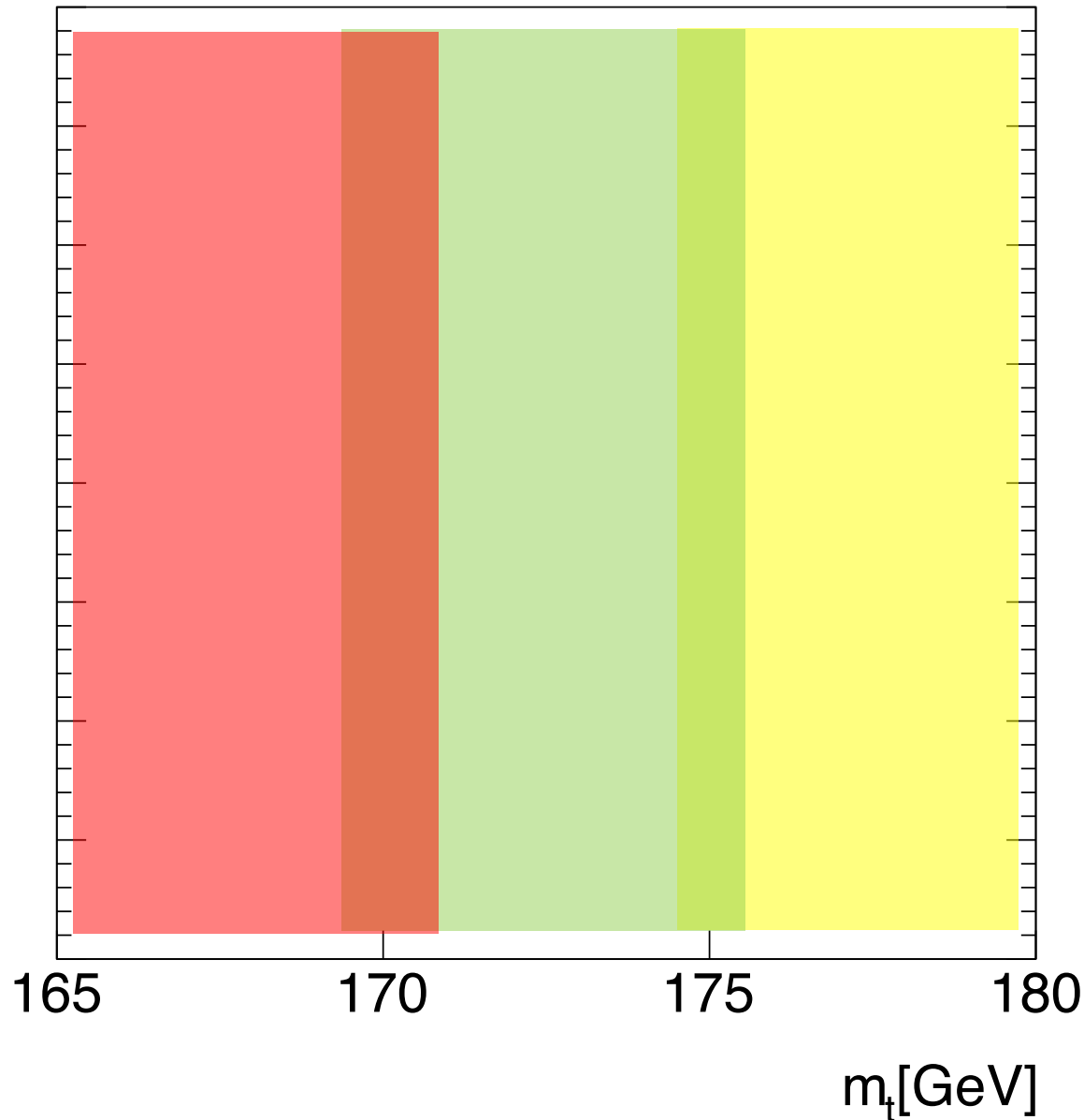
$$s = (p_{e^-} + p_{e^+})^2 \quad \rightarrow \quad s' = (p_{e^-} + p_{e^+})^2$$

$$s' = s \left(1 - \frac{2E_\gamma}{\sqrt{s}}\right)$$

J. Fuster

# The ups and downs of the top-quark mass: **discussion (qualitative)**

Keeping present values/uncertainties of  $m_h$  and  $m_W$



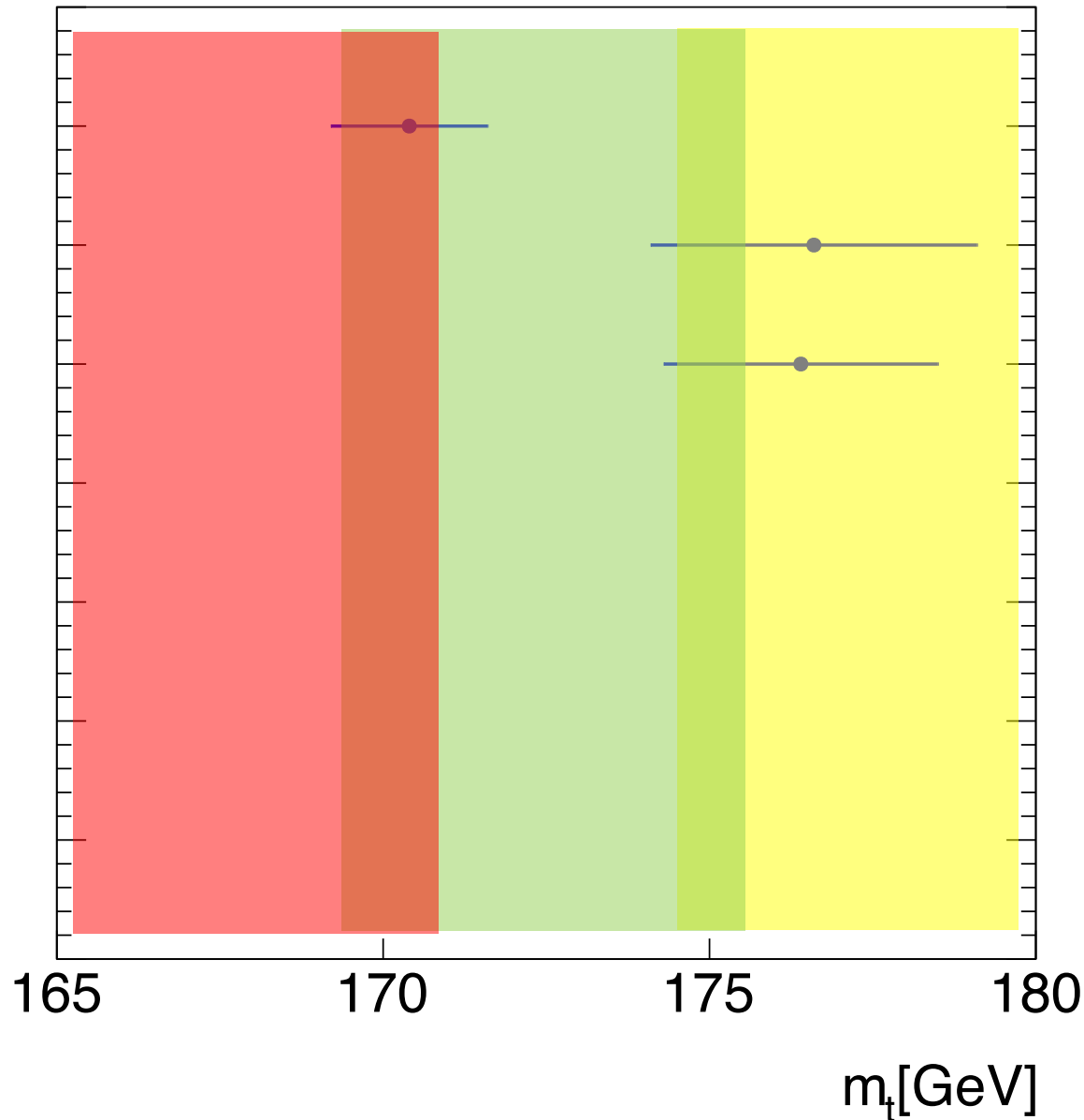
New Physics at Low EW Energy Scales  
(Heinemeyer et al.)

Need for New Physics @ Large Energy Scales  
**Vacuum Stability:** Meta-stable Universe

Need for New Physics @ Large Energy Scales  
**Vacuum Stability:** Unstable Universe

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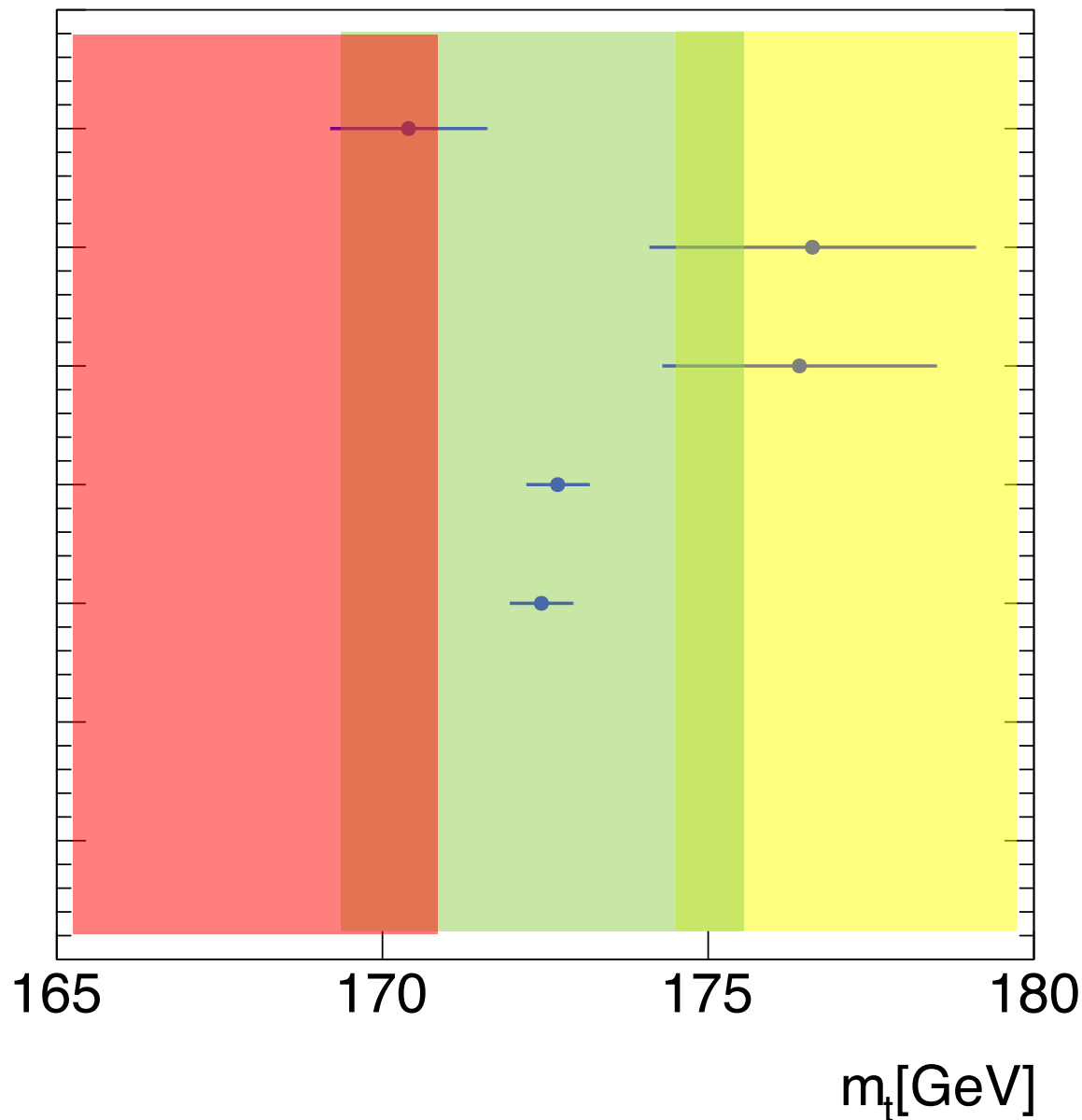
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**Vacuum Stability: Unstable**

**Global Fit – NNLO QCD**

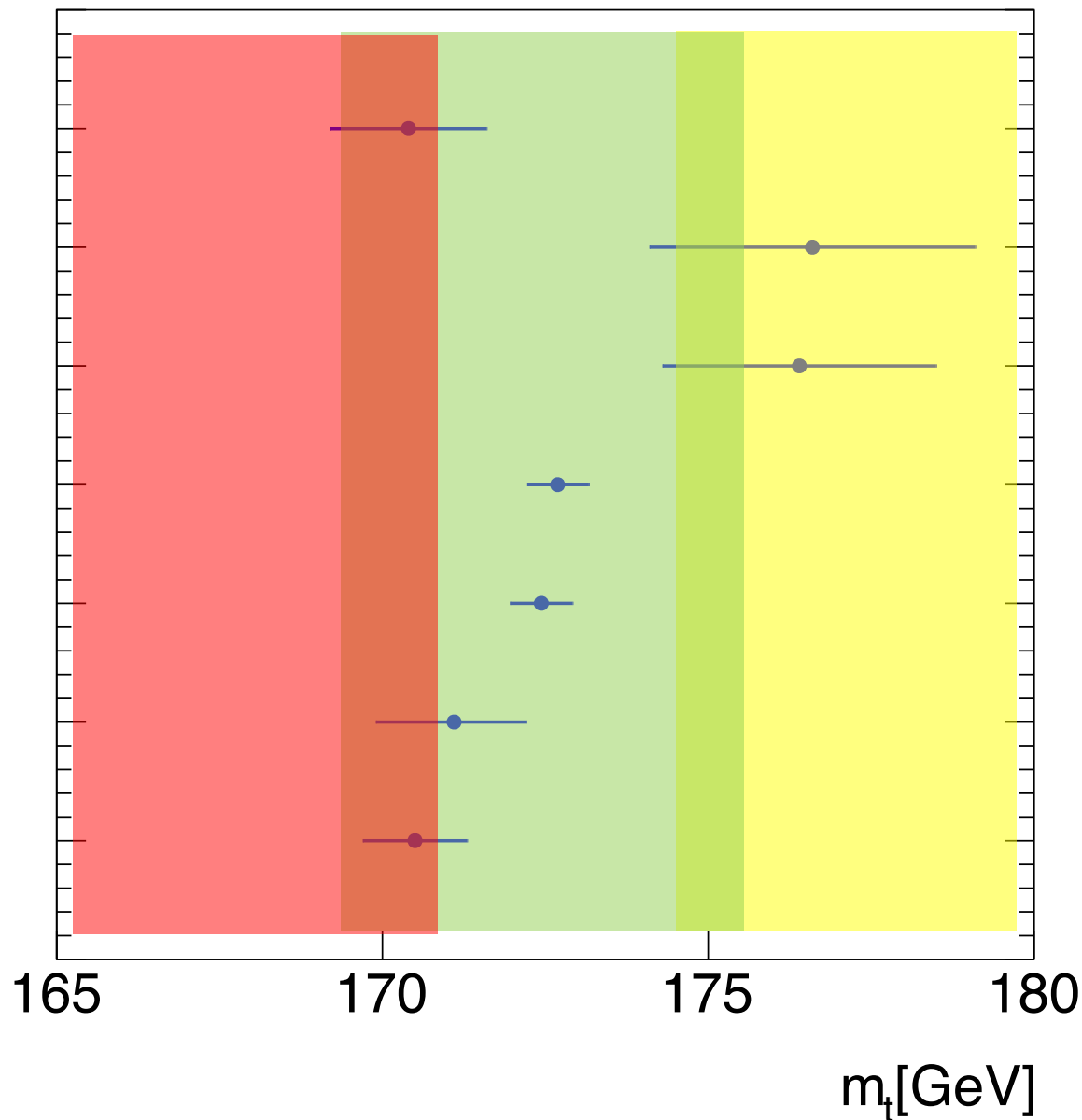
**EW- Fits: HEPfit. & Gfitter**

**ATLAS & CMS**

Combinations of “Direct Measurements”

# The ups and downs of the top-quark mass: discussion (qualitative)

Keeping present values/uncertainties of  $m_h$  and  $m_W$



New Physics at Low EW Energy Scales  
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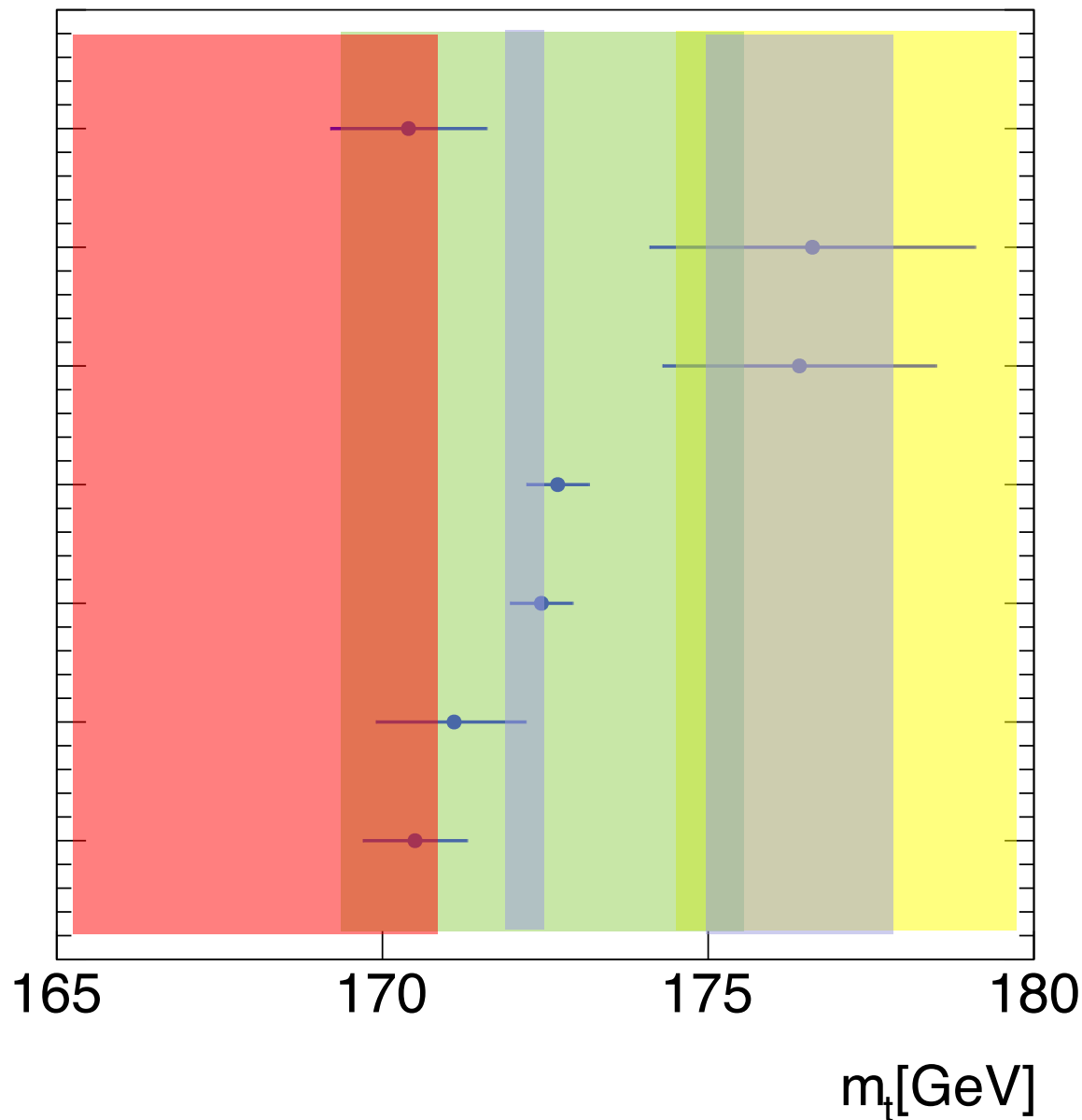
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Total & Differential Cross-Sections  
3D and tt+1jet



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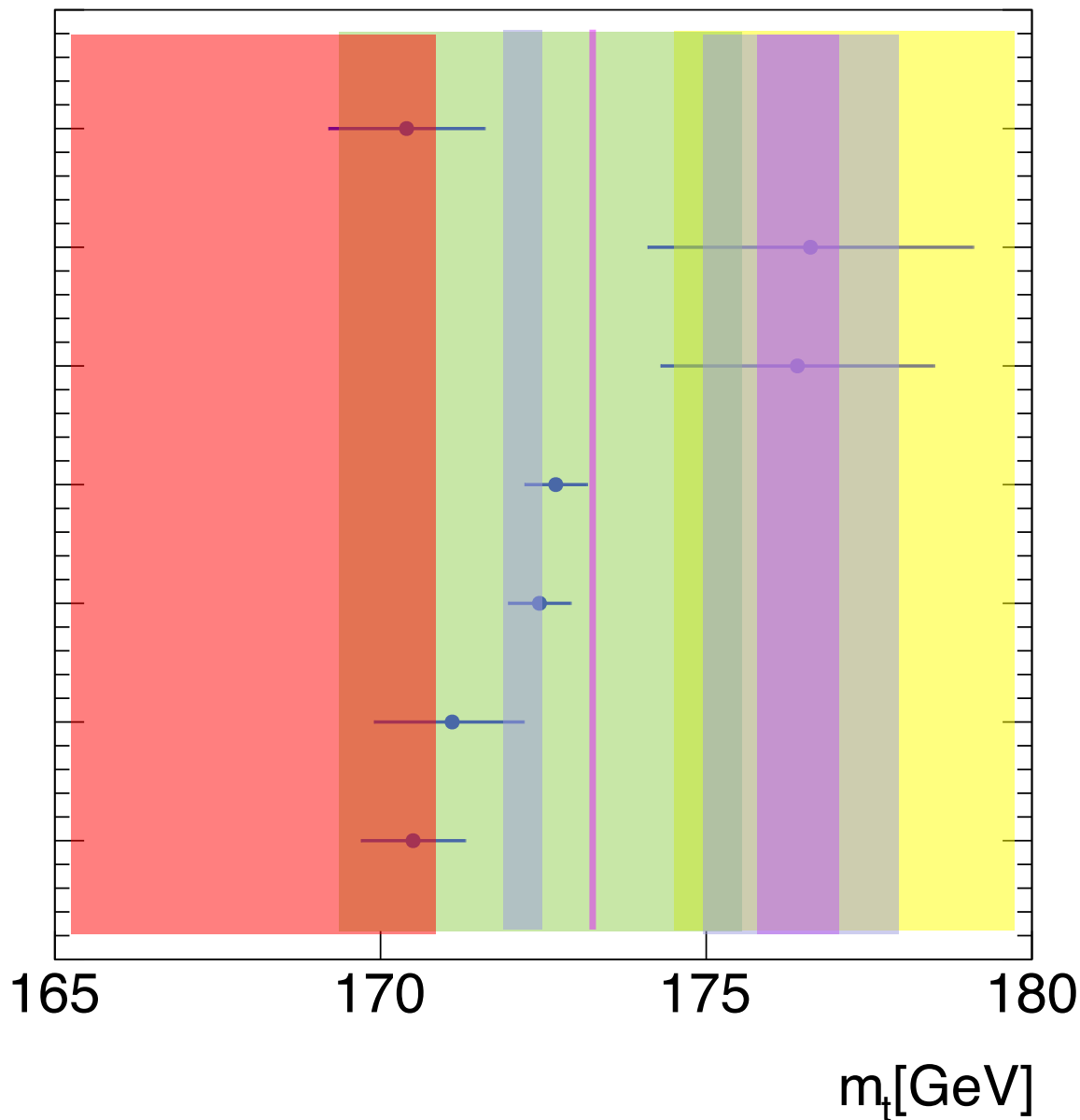
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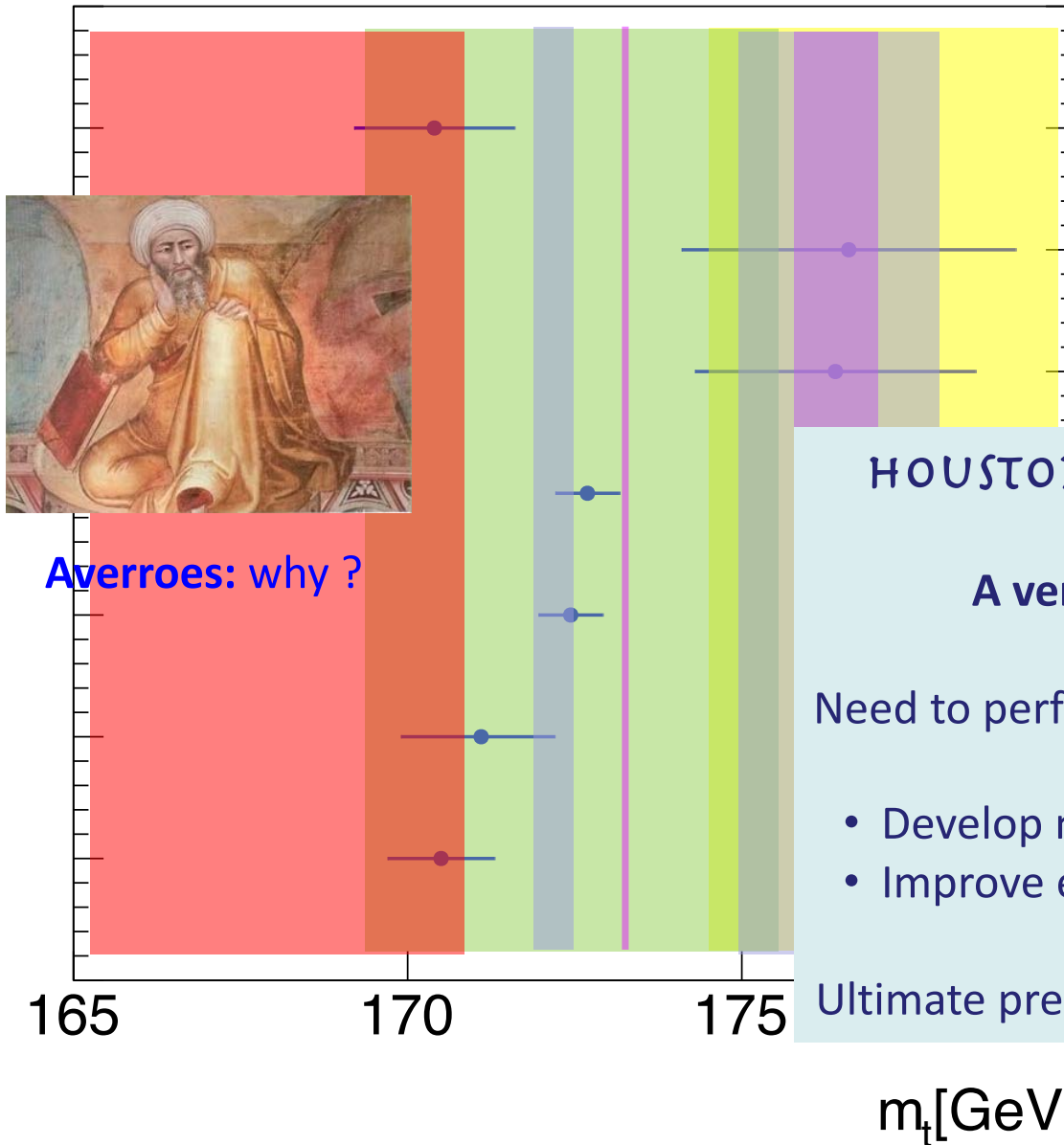
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Future e+e- Collider

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**HOUSTON WE HAVE A CHALLENGE !!!**

**A very interesting challenge in fact**

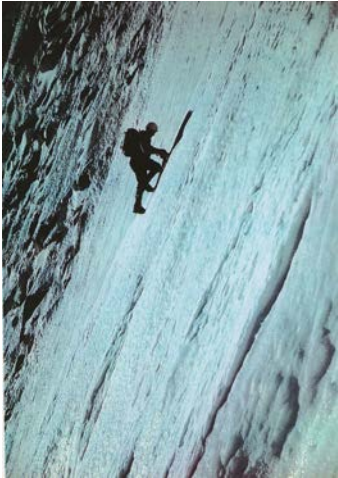
Need to perform highly precise  $m_t$  measurements:

- Develop new calculations/observables
- Improve experimental methods

Ultimate precision at e+e- future colliders

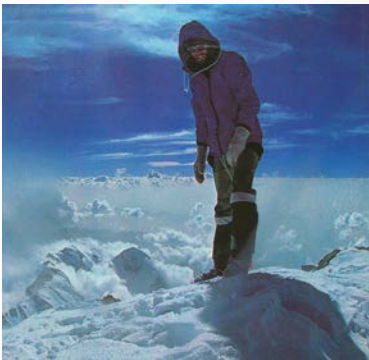
Future e+e- Collider

# The top quark is *beautiful and charming and still very strange*



Long and successful scientific programme, many studies, resources, and investigations during years of research in theory and experiment (PETRA, PEP, Babar, Belle, HERA, SLD, LEP, Tevatron, LHC, etc..) have led to build up the Standard Model

Reinhold Messner



Culminated with the discovery of H(125)



But.. this is just one more “step” which allows us to have a “better view” of what is coming next.

- One question answered, H(125)
- Still some old questions remain
- New questions open

# In much nicer words and better English

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“It is the great beauty of our science that advancement in it, whether in a degree great or small, instead of exhausting the subject of research, opens the doors to further and more abundant knowledge, overflowing with beauty and utility”, (M. Faraday)