

The road to the ElectroWeak Symmetry Breaking

14th - 15th February 2012
Seminar at DESY

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Outline

- ❑ The Higgs boson and the **ElectroWeak Symmetry Breaking (EWSB)**
- ❑ **Status: CMS results** in a nutshell, focusing on **high mass H->VV**
- ❑ **What's next ?**
 - move to **larger mass (>600 GeV beyond SM)**
 - control of V+jets background
 - jet merging
 - signal characterization (angular analysis)
 - improve sensitivity to **smaller xsec -> Vector Boson Fusion**
 - control of VV background
- ❑ The final arbiter: **VV scattering**

Why we need the Higgs

The Higgs boson provides

- 1) an EXPLICATION of the **W,Z mass (ie EWSB)**
 - 2) a DESCRIPTION of the **fermions masses**
- **1 is really fundamental to make the SM “working”** (next slides)
... even if not less arbitrary!
 - 2 is just another way of formulating the same question:
 - why the fermions have those particular masses?
 - why the fermions have those particular Higgs couplings?(SM works well without 2: just the **fermio-phobic Higgs**)

EWSB and the W, Z mass

$$L_{gauge} = -\frac{1}{4}W_{\mu\nu}^i W^{\mu\nu i} - \frac{1}{4}B_{\mu\nu} B^{\mu\nu} + \frac{1}{2}m^2 W_{\mu\nu}^i W^{\mu\nu i} + \frac{1}{2}m^2 B_{\mu\nu} B^{\mu\nu}$$

}

$W_{\mu\nu}^i = \partial_\nu W_\mu^i - \partial_\mu W_\nu^i + g\varepsilon^{ijk}W_\mu^j W_\nu^k$

$B_{\mu\nu} = \partial_\nu B_\mu - \partial_\mu B_\nu$

SU(2)
×
U(1)

Gauge invariance



scalar potential ($\lambda > 0, \mu < 0$) $V(\Phi) = \mu^2 |\Phi^+ \Phi| + \lambda (\Phi^+ \Phi)^2$

complex scalar
doublet of SU(2) $\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$

with minimum (= empty state) at
(v = empty expectation value)

$\langle \Phi \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$ or ~~$\langle \Phi \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} v \\ 0 \end{pmatrix}$~~ \rightarrow ~~SU(2)~~

generic Gauge

$$\Phi = \frac{1}{\sqrt{2}} e^{i\frac{\omega^i \tau^i}{2v}} \begin{pmatrix} 0 \\ v+h \end{pmatrix}$$

1 physical scalar field $\langle h \rangle = 0$

3 Goldstone bosons ω_j

4 Gauge fields W_μ^i, B_μ

Gauge unitario

$$\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v+h \end{pmatrix}$$

1 physical scalar field -> Higgs

4 Gauge fields combined into known vector bosons:
W,Z with mass, photon massless

$$W_\mu^\pm = \frac{1}{\sqrt{2}} (W_\mu^1 \mp iW_\mu^2) \quad M_W^2 = \frac{1}{4} g^2 v^2$$

$$Z^\mu = \frac{-g'B_\mu + gW_\mu^3}{\sqrt{g^2 + g'^2}} \quad M_Z^2 = \frac{1}{4} (g^2 + g'^2) v^2$$

$$A^\mu = \frac{-gB_\mu + g'W_\mu^3}{\sqrt{g^2 + g'^2}} \quad M_A^2 = 0 \quad \mathbf{U(1)_{EM}}$$

Higgs and unitarity in VBF

W,Z mass (-> longitudinal degrees of freedom) arise from the Higgs mechanism:

without Higgs, $W_L^+ W_L^- \rightarrow W_L^+ W_L^-$
would break unitarity

$$A(W_L^+ W_L^- \rightarrow W_L^+ W_L^-) \approx \frac{1}{v^2} \left(-s - t + \frac{s^2}{s - m_H^2} + \frac{t^2}{t - m_H^2} \right)$$

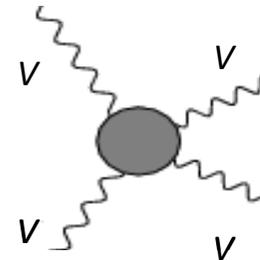
Same behavior for all VV amplitudes

$$A(Z_L Z_L \rightarrow W_L^+ W_L^-) \approx \frac{1}{v^2} \left(-s + \frac{s^2}{s - m_H^2} \right) \text{ (s channel only)}$$

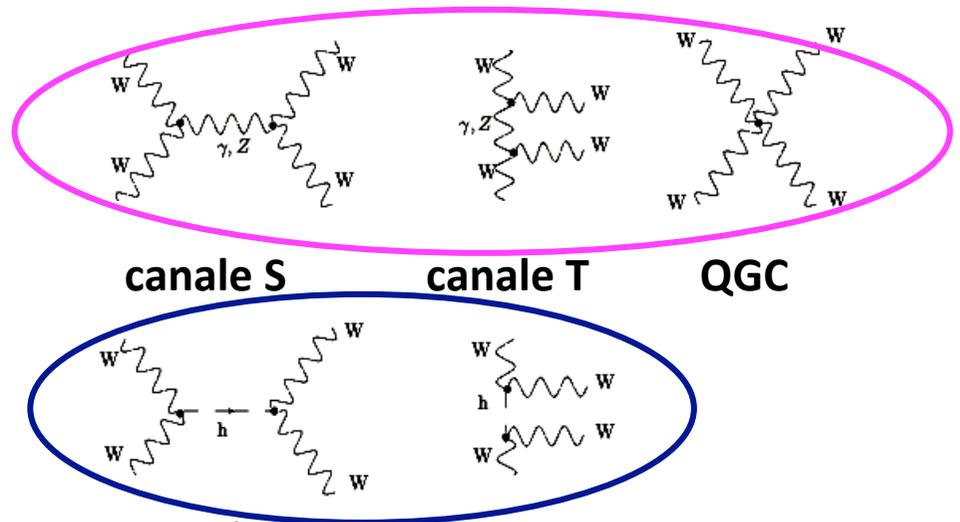
$$A(Z_L W_L^- \rightarrow Z_L W_L^-) \approx \frac{1}{v^2} \left(-t + \frac{t^2}{t - m_H^2} \right) \text{ (t channel only)}$$

$$A(W_L^- W_L^- \rightarrow W_L^- W_L^-) \approx \frac{1}{v^2} \left(-u - t + \frac{u^2}{u - m_H^2} + \frac{t^2}{t - m_H^2} \right) \text{ (t and u channels)}$$

Vector Boson Fusion (VBF)



$VV \rightarrow VV$

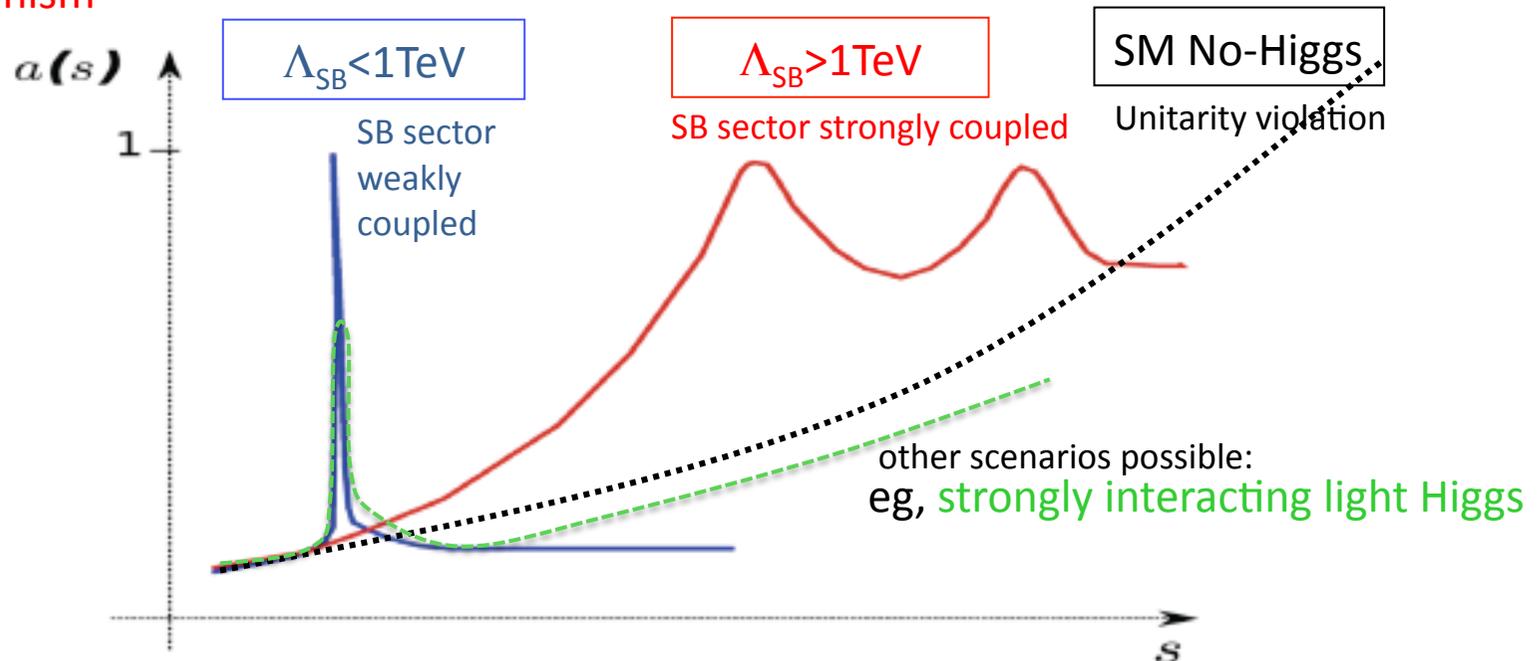


VBF is the smoking gun of the EWSB !

VBF and VV scattering

□ VV scattering spectrum $\sigma(VV \rightarrow VV)$ vs $M(VV)$

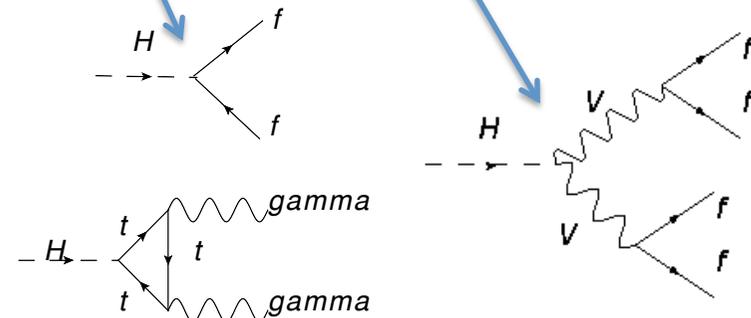
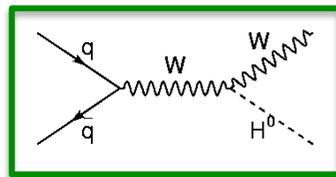
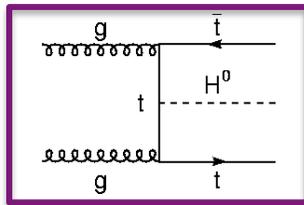
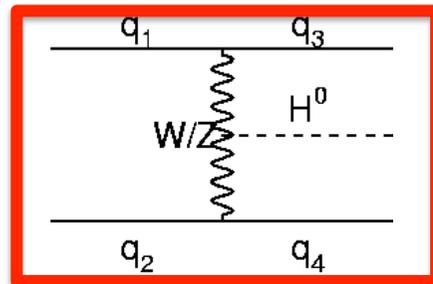
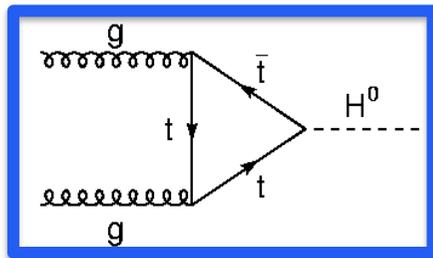
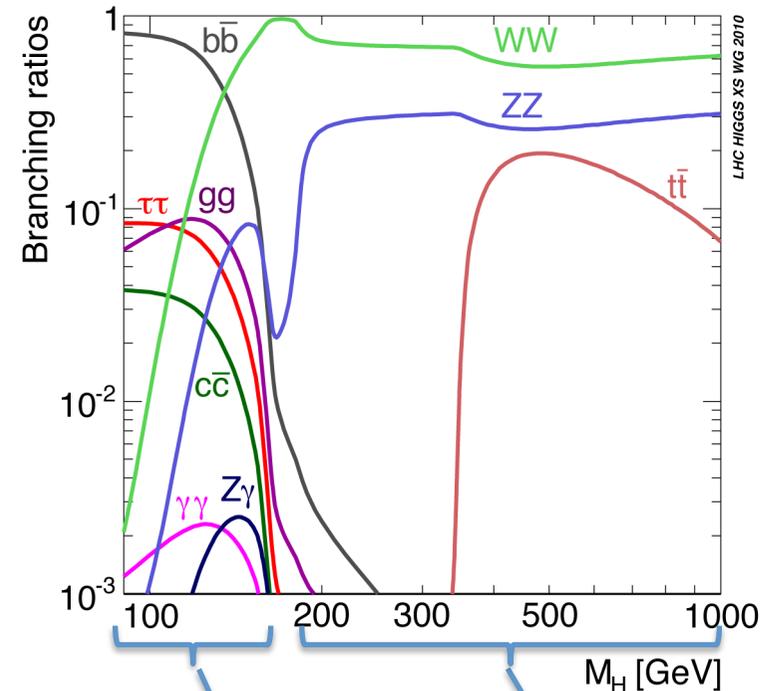
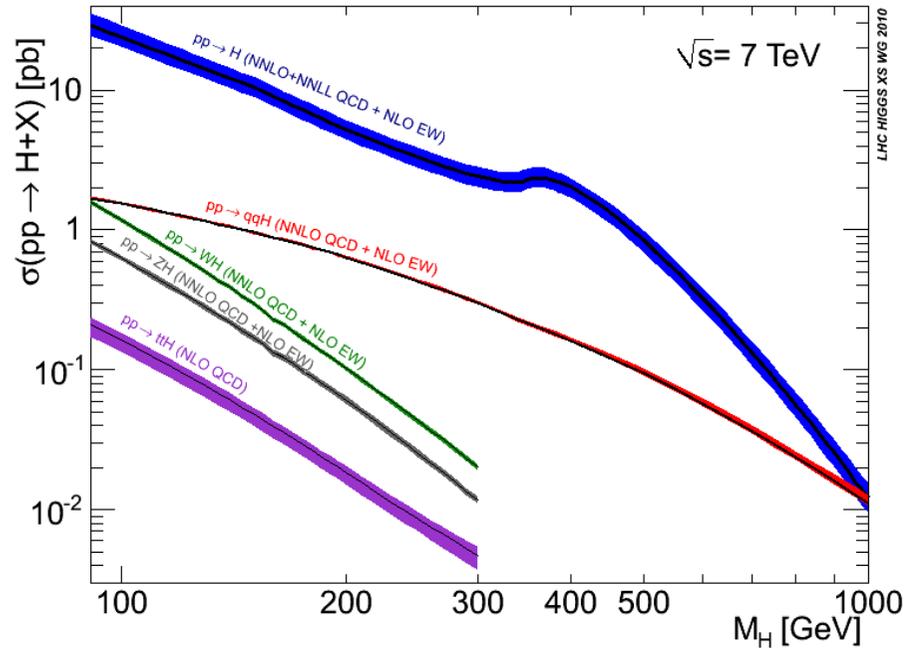
is the fundamental probe to test **nature of Higgs boson** or to find **alternative EWSB mechanism**



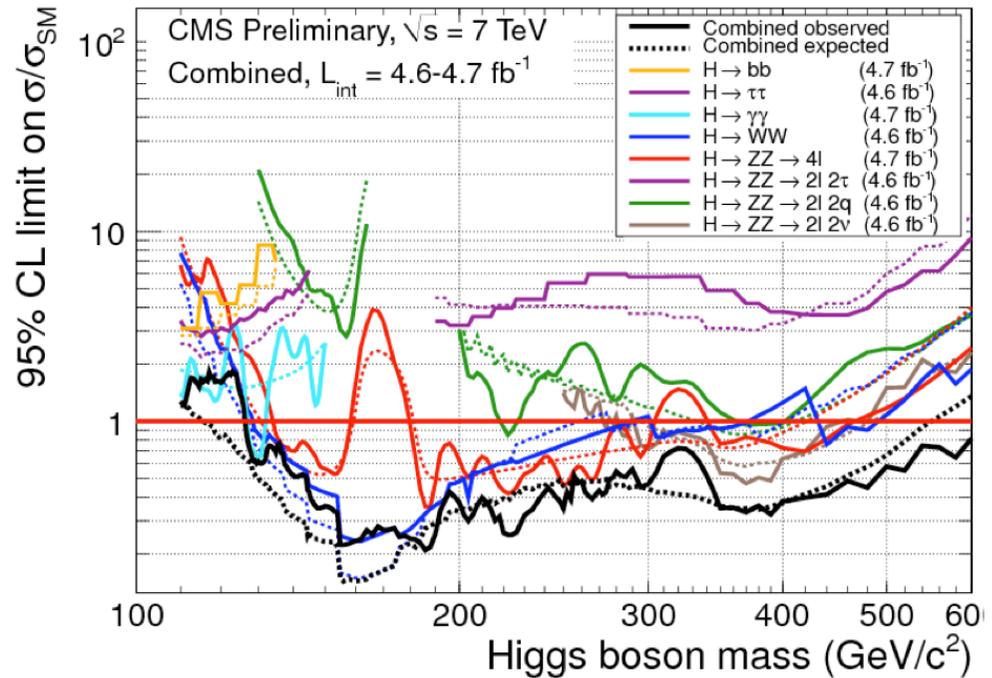
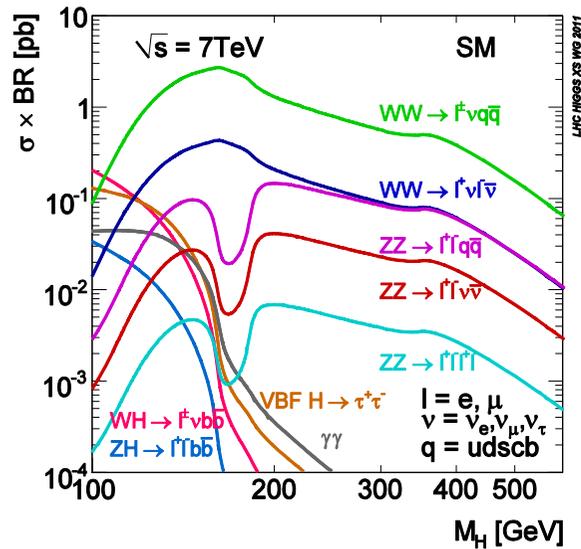
□ Whatever we will see or not see at low mass ($< 2 \times m_W$), the EWSB mechanism must be probed in the VV final state

- search for possible **resonances in VBF**
- measurement of **VV scattering spectrum**

Higgs production and decay



CMS results



- **WW→lnqq, ZZ→llqq** limited by huge **V+jets background**, taken from simu/data with large theoretical/statistical error
- **WW→lnln** at high mass limited by signal << **WW background** ($\Delta\phi$ not effective)
- **ZZ→lnn**:
 - 200-400 GeV limited by **non-Z background** (top, W+jets, WW)
 - >400 GeV limited by **Z+jets tail at high MET**: not large but not well known (controlled with γ +jets → statistical error+met uncertainty)

➡ drives the UL for $m_H > 300-400$

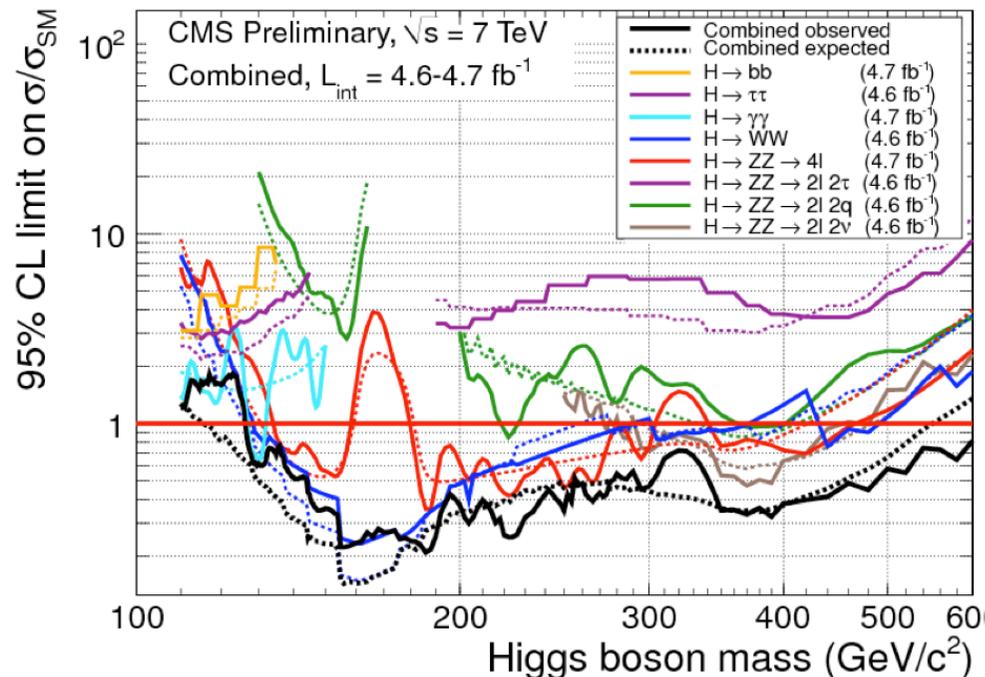
- **ZZ→4l** limited by **statistics** (only ZZ background: small and well known)

➡ drives the UL for m_H 200-300

Future improvements

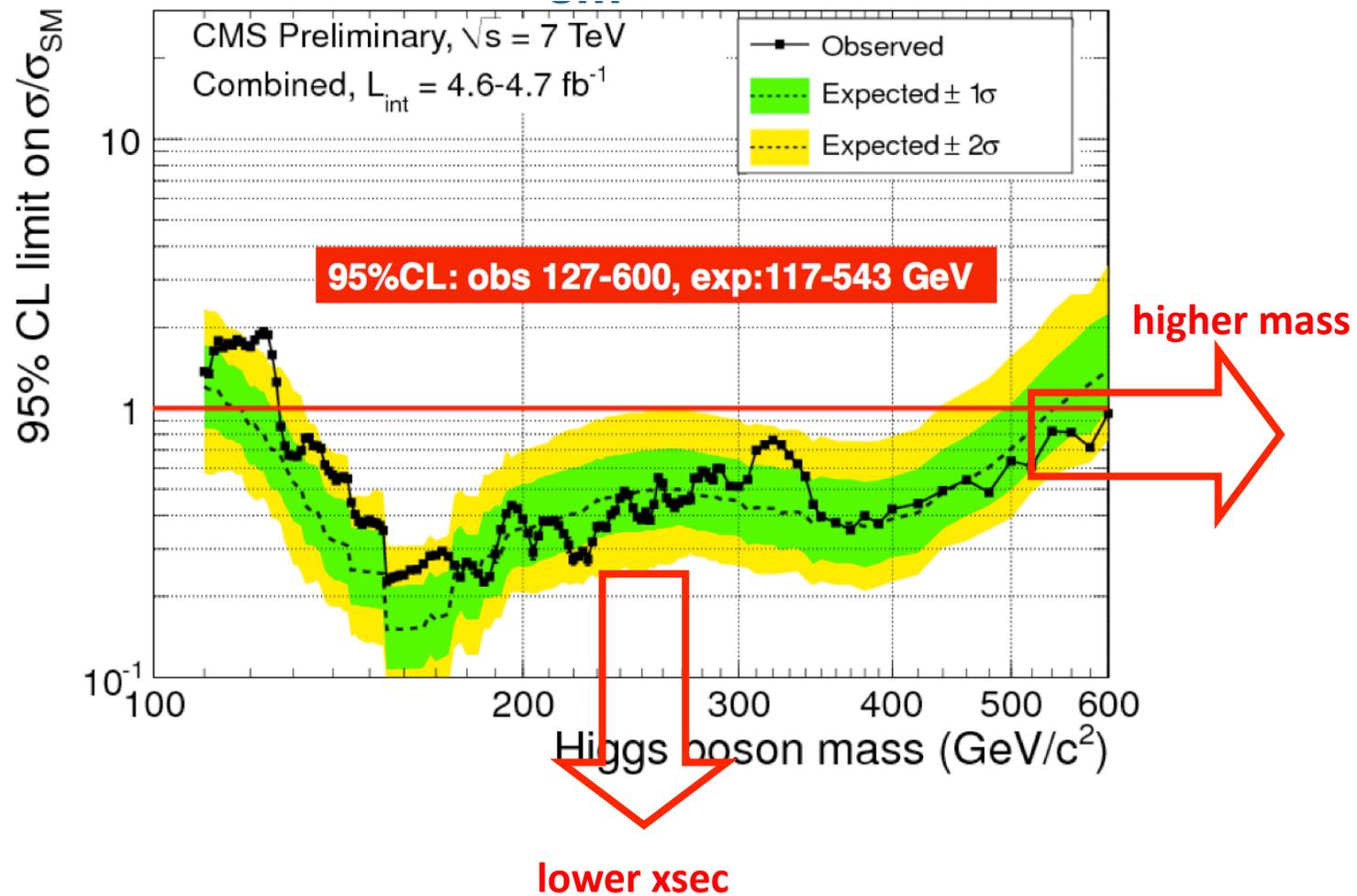
?

- Combination of **>5 different channels** (ele, mu, btag, ...) Robust!



- **Very optimized analyses**, some space for further improvement. With higher lumi:
 - use shape analyses (where not yet done)
 - extract signal with multidimensional fit (now only mZZ fit)
 - extract background (norm and shape) from data with lower uncertainty

What's next ?



1 TeV masses: not anymore “the” Higgs

→ **General search for $X \rightarrow VV \rightarrow 4f$:** exotic models (eg, Technicolor, ExtraDimension, ...)

RS Graviton vs SM Higgs:

Signal Channels				
	140 GeV	250 GeV	500 GeV	1 TeV
$\sigma(H_{SM} \rightarrow ZZ^{(*)})$	9.26×10^2	1.14×10^3	2.40×10^2	-
$\sigma(G_{RS} \rightarrow ZZ)$ at $c = 0.01$	-	1.94×10^3	6.83×10^1	1.41
$\sigma(G_{RS} \rightarrow ZZ)$ at $c = 0.05$	-	4.83×10^4	1.69×10^3	3.53×10^1
$\sigma(G_{RS} \rightarrow ZZ)$ at $c = 0.1$	-	1.94×10^5	6.76×10^3	1.41×10^2
$\mathcal{B}(H_{SM} \rightarrow ZZ^{(*)})$	0.068	0.295	0.260	-
$\mathcal{B}(G_{RS} \rightarrow ZZ)$	-	0.052	0.049	0.046
$\mathcal{B}(ZZ \rightarrow 4l)$	$\sim 0.067 \times 0.067 \simeq 4.48 \times 10^{-3}$			
$\mathcal{B}(ZZ \rightarrow 2l2j)$	$\sim 2 \times 0.067 \times 0.699 \simeq 8.96 \times 10^{-2}$			
$N(H_{SM} \rightarrow ZZ^{(*)} \rightarrow 4\mu)$	0.52	0.64	0.13	-
$N(H_{SM} \rightarrow ZZ^{(*)} \rightarrow 4e)$	0.52	0.64	0.13	-
$N(H_{SM} \rightarrow ZZ^{(*)} \rightarrow 2\mu 2e)$	1.04	1.28	0.27	-
$N(H_{SM} \rightarrow ZZ \rightarrow 2\mu 2j)$	-	22.4	38.9	-
$N(H_{SM} \rightarrow ZZ \rightarrow 2e 2j)$	-	21.5	45.5	-
$N(G_{RS} \rightarrow ZZ \rightarrow 4\mu)$ at $c = 0.01$		1.7	0.06	0.0013
$N(G_{RS} \rightarrow ZZ \rightarrow 4e)$		1.7	0.06	0.0013
$N(G_{RS} \rightarrow ZZ \rightarrow 2\mu 2e)$		3.4	0.12	0.0025
$N(G_{RS} \rightarrow ZZ \rightarrow 4\mu)$ at $c = 0.05$		42.9	1.4	0.028
$N(G_{RS} \rightarrow ZZ \rightarrow 4e)$		42.9	1.4	0.028
$N(G_{RS} \rightarrow ZZ \rightarrow 2\mu 2e)$		85.8	2.8	0.055
$N(G_{RS} \rightarrow ZZ \rightarrow 4\mu)$ at $c = 0.1$		162	5.3	0.11
$N(G_{RS} \rightarrow ZZ \rightarrow 4e)$		162	5.3	0.11
$N(G_{RS} \rightarrow ZZ \rightarrow 2\mu 2e)$		324	10.6	0.22

CMS AN-2010-35:
Angular Analysis of Resonances $pp \rightarrow X \rightarrow ZZ$
 A. Bonato, A.V. Gritsan, Z.J. Guo, N.V. Tran, A. Whitbeck
 Johns Hopkins University, Baltimore, MD, USA

xsec larger than Higgs:

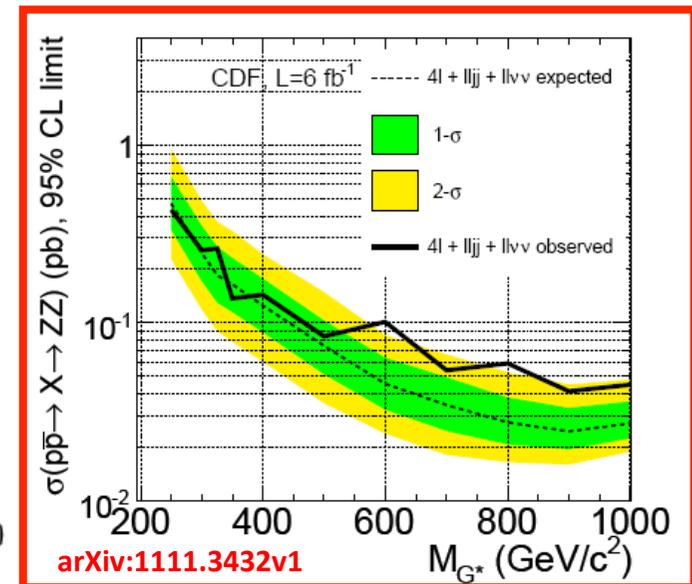
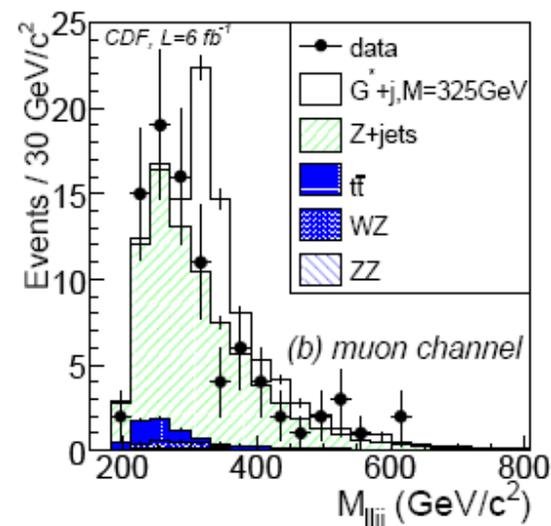
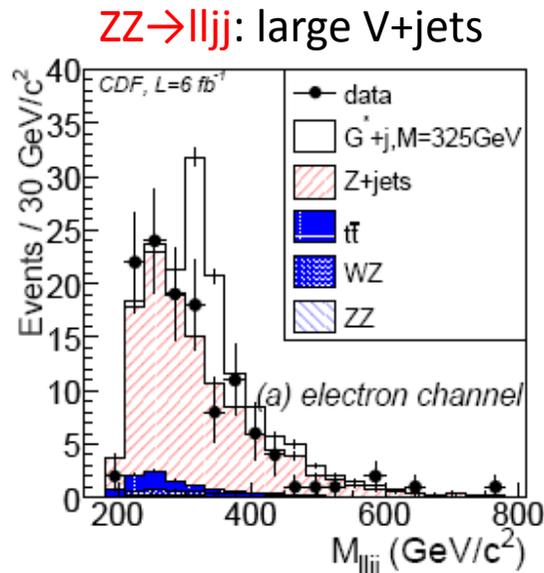
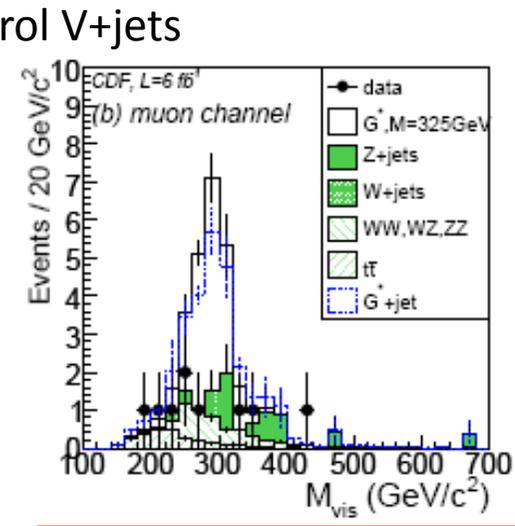
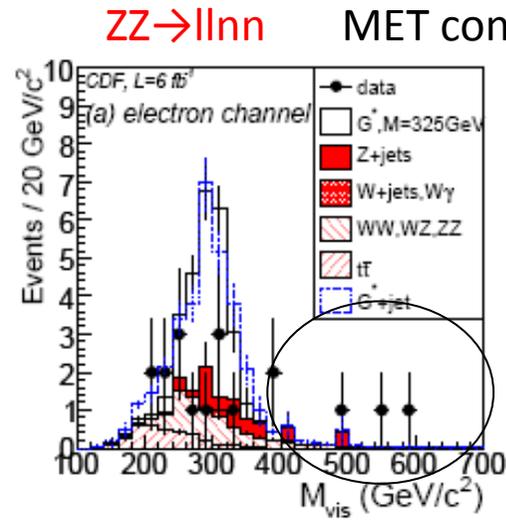
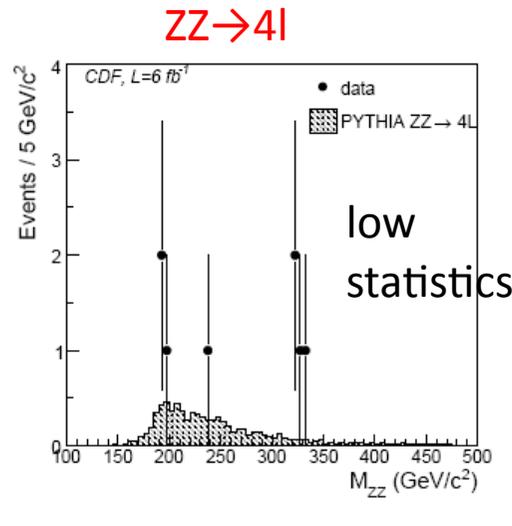
first, repeat “Higgs” search for different spin, width resonance

at high mass still very low number of events per fb⁻¹

→ importance of semileptonic final states

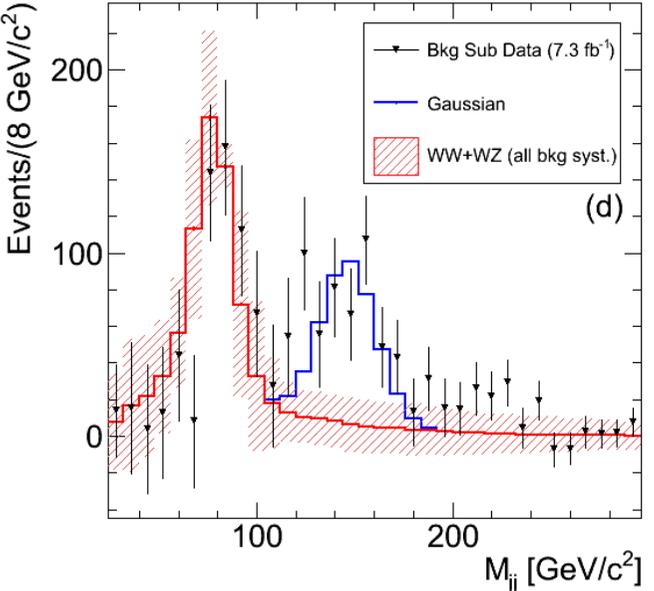
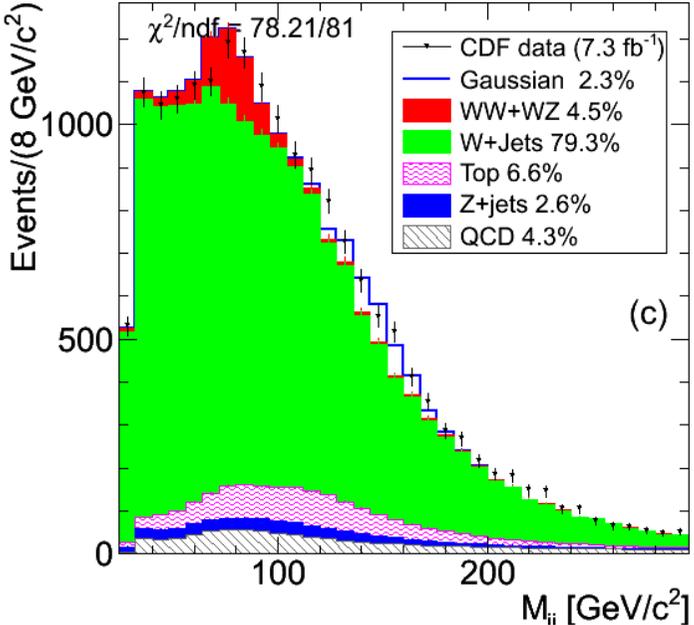
Available results: ZZ

- **CDF search for $G \rightarrow ZZ$** : same features discussed for high mass Higgs @ LHC

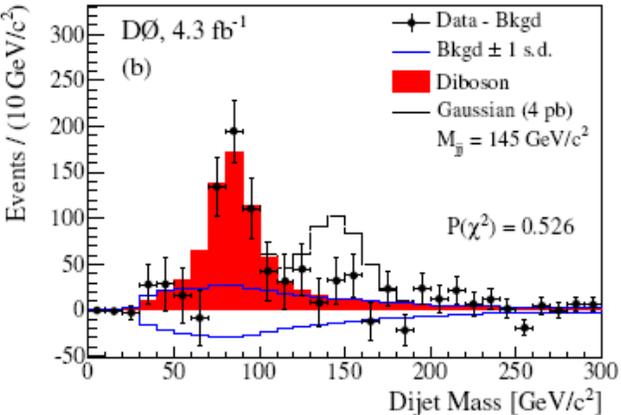
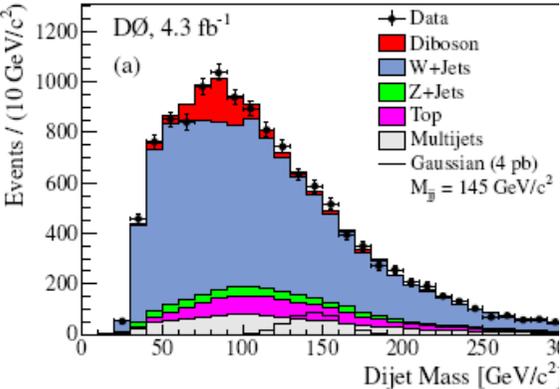
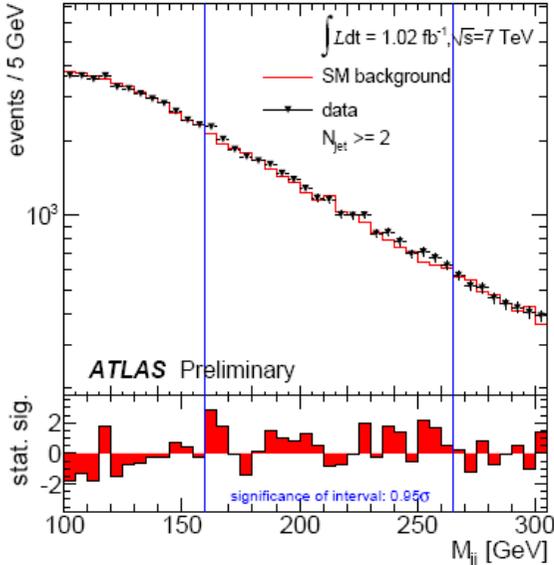


Available results: W+2jets

CDF "bump"



ATLAS & D0 xchecks



Control of V+jets

- ❑ **Control region** (eg, $Z \rightarrow jj$ sidebands) has very low stat for $M(l\bar{l}jj) \sim 1$ TeV

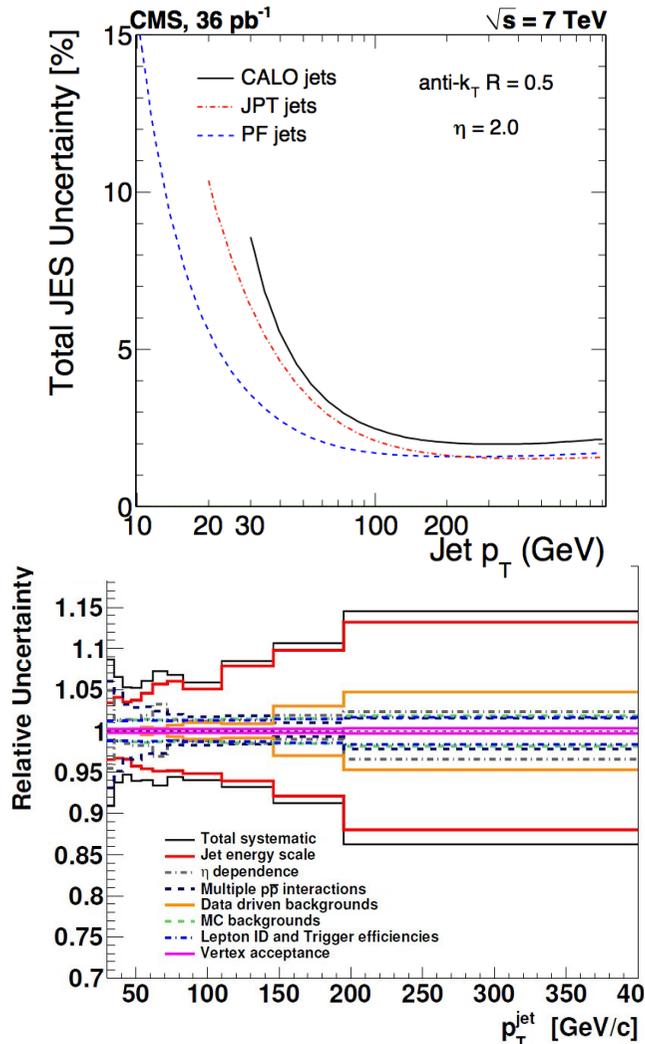
- ❑ Improving **theoretical tools** (Blackhat, Madgraph, ...)
 - test them where we have statistics
 - rely on them to extrapolate at higher energy/multiplicity

V+jets

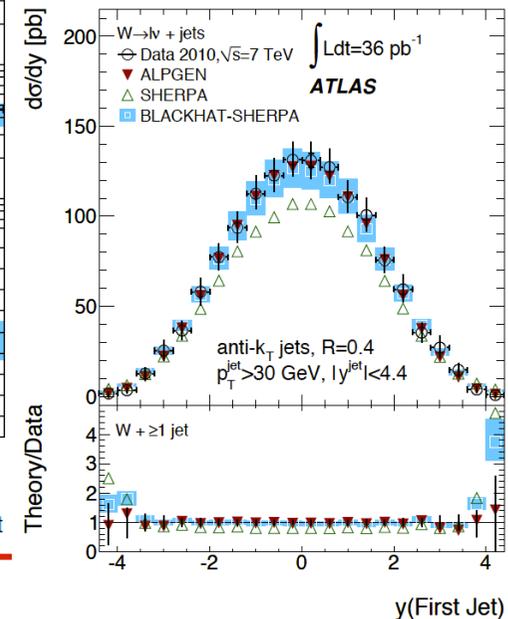
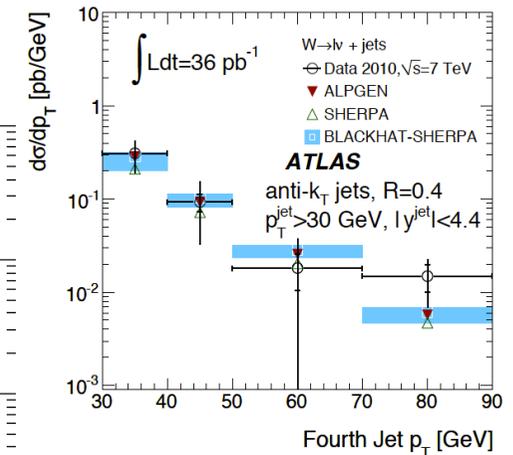
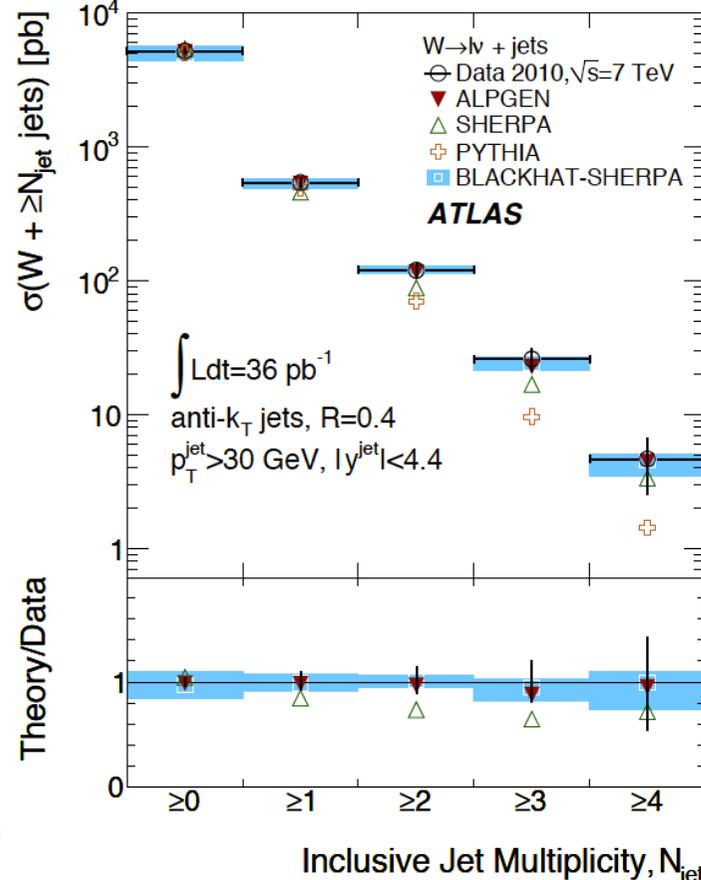
QCD measurement (jet $p_T > 20-30$ GeV):

→ syst. dominated by jet scale, PileUp removal

Data unfolded for detector effects → compared to NLO (“hadron level”)



ATLAS:

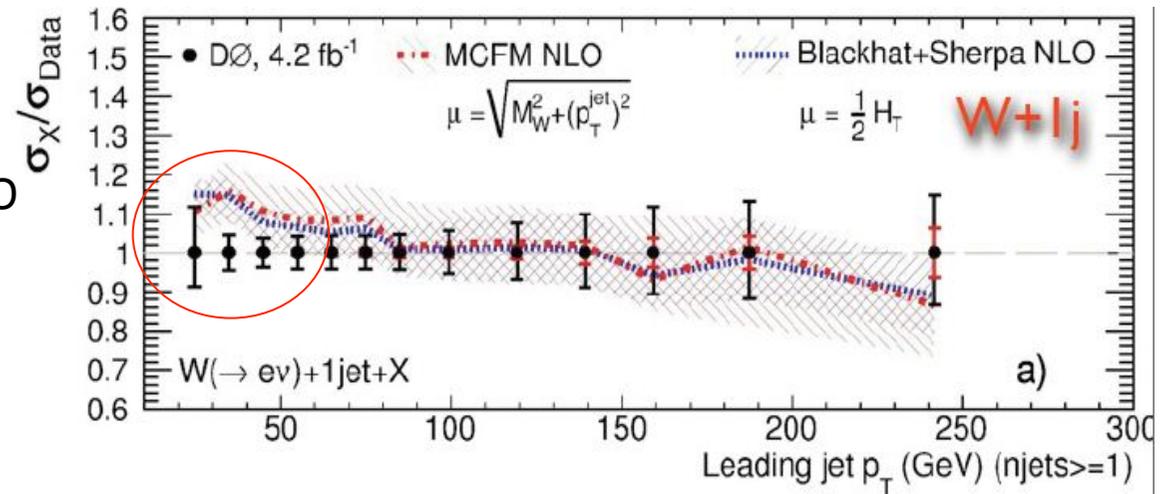


V+jets at Tevatron

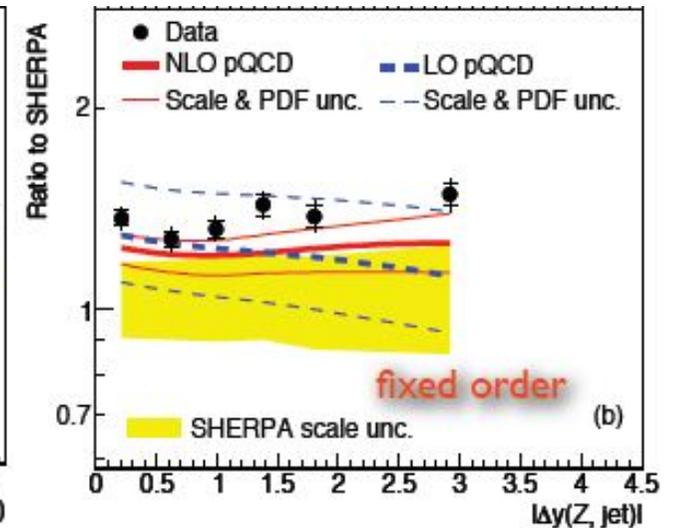
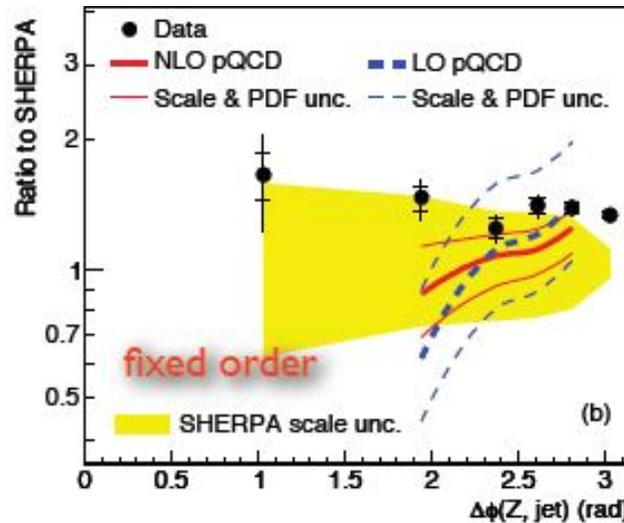
At low p_T , low multiplicity:
interesting discrepancy data-NLO
observed

but results limited by systematics

→ new variables



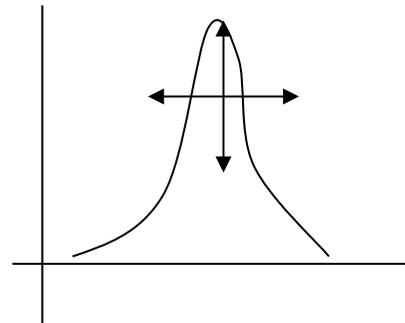
D0 novel measurement:
angular correlations have
much lower systematics



High mass: what's new ?

- ❑ Can we simply keep the same Higgs analysis strategy? Not at very high masses!
- ❑ New experimental issues at very high mass (1 TeV and above)
 - X → boosted VV → **jet merging (and nearby leptons)**
- ❑ Unknown signal and very small background → no point in pushed optimization! Keep **model independent approach** as much as possible

- ❑ How to disentangle the various models?
 - peak → mass and width, xsec and BR
 - spin! → **angular analysis**

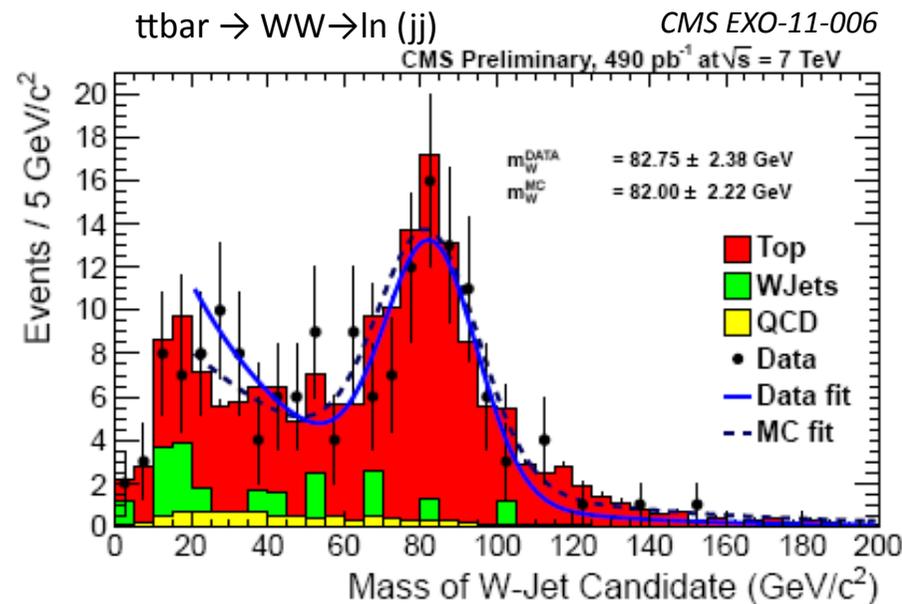


>1 TeV M(ZZ)→4f : jet merging (1)

□ **Jet merging:** $\Delta R_{qq} \approx \frac{1}{\sqrt{z(1-z)}} \frac{M_Z}{p_Z^T} \xrightarrow{\text{approx}} \Delta R \text{ 0.8 (CA)} \rightarrow M_X > 600 \text{ GeV}$
 $(z = |\vec{p}_q| / |\vec{p}_Z|)$ $\Delta R \text{ 0.5 (Akt)} \rightarrow M_X > 900 \text{ GeV}$

Handles to distinguish wrt to jets from QCD (eg, X→ZZ→2l2j VS Z+jets):

▪ **jet mass**

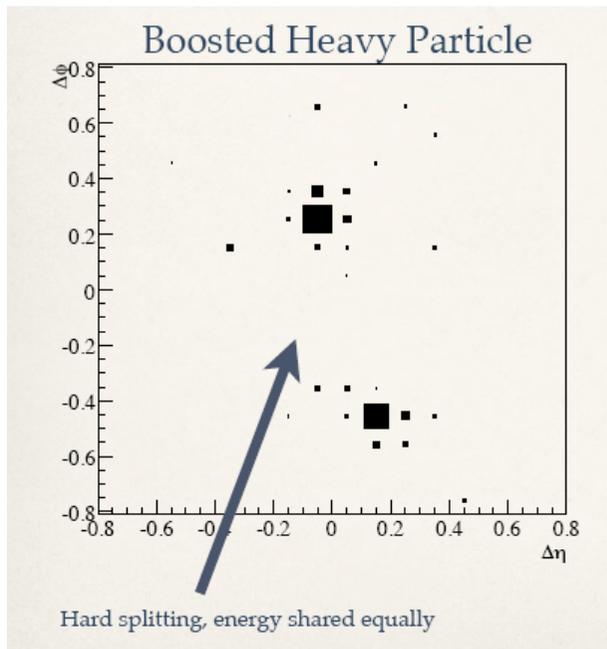


Jet merging (2)

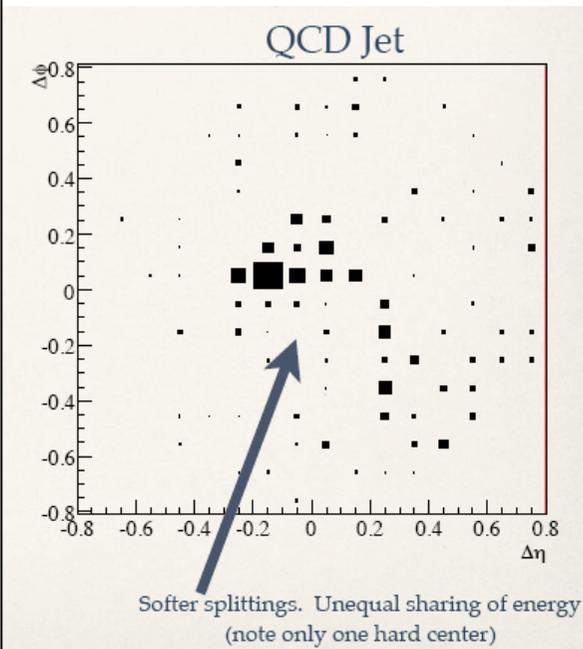
Handles to distinguish wrt to jets from QCD (eg, $X \rightarrow ZZ \rightarrow 2l2j$ VS Z +jets):

- **jet radiation:**

no singularity, just decay!



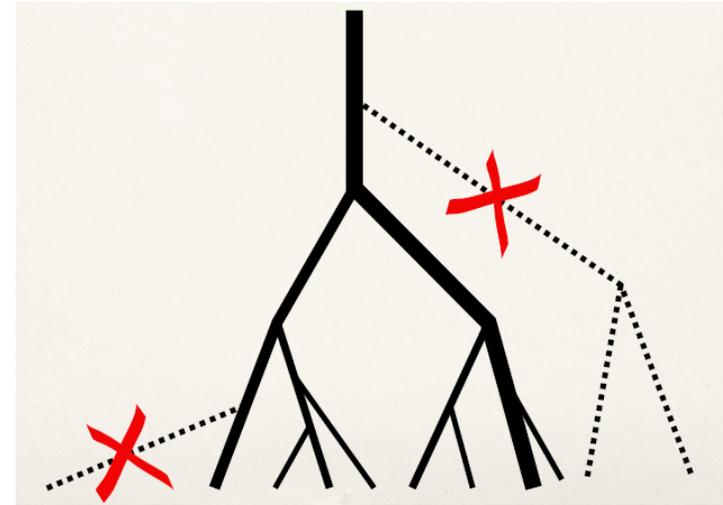
soft/collinear singularity in QCD



JHU seminar: **Path-Integral Jets** by David Krohn (Harvard)
www.pha.jhu.edu/groups/particle-theory/seminars/talks/F11/talk.khron.pdf

Jet pruning

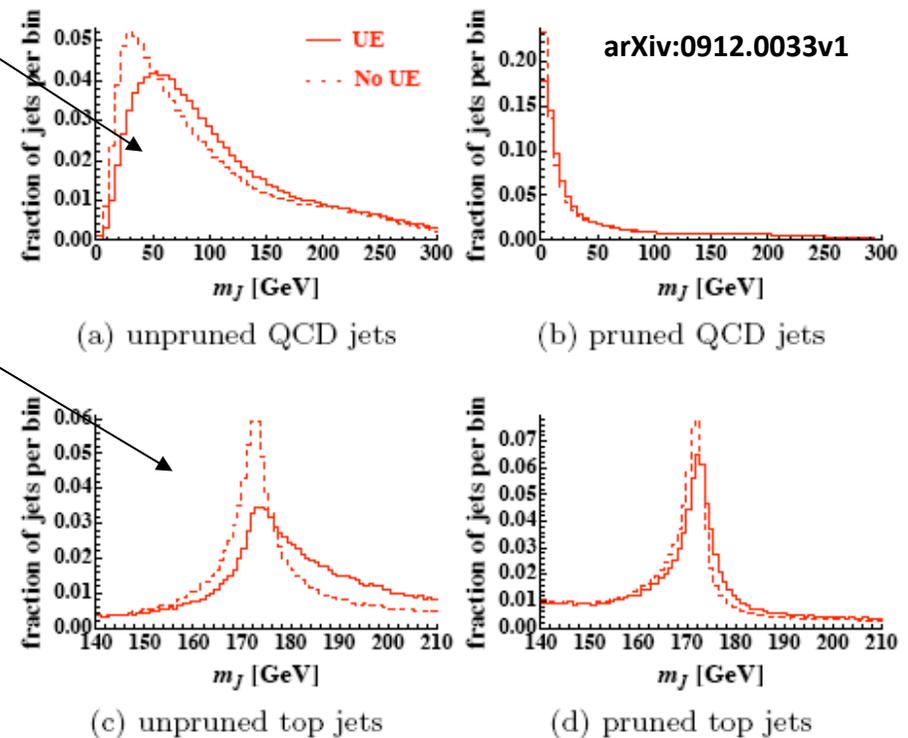
Remove all parts of the jet which are soft and wide angle



QCD jets mass substantially decreased -> lower backgrounds

Boosted objects mass reconstruction improved

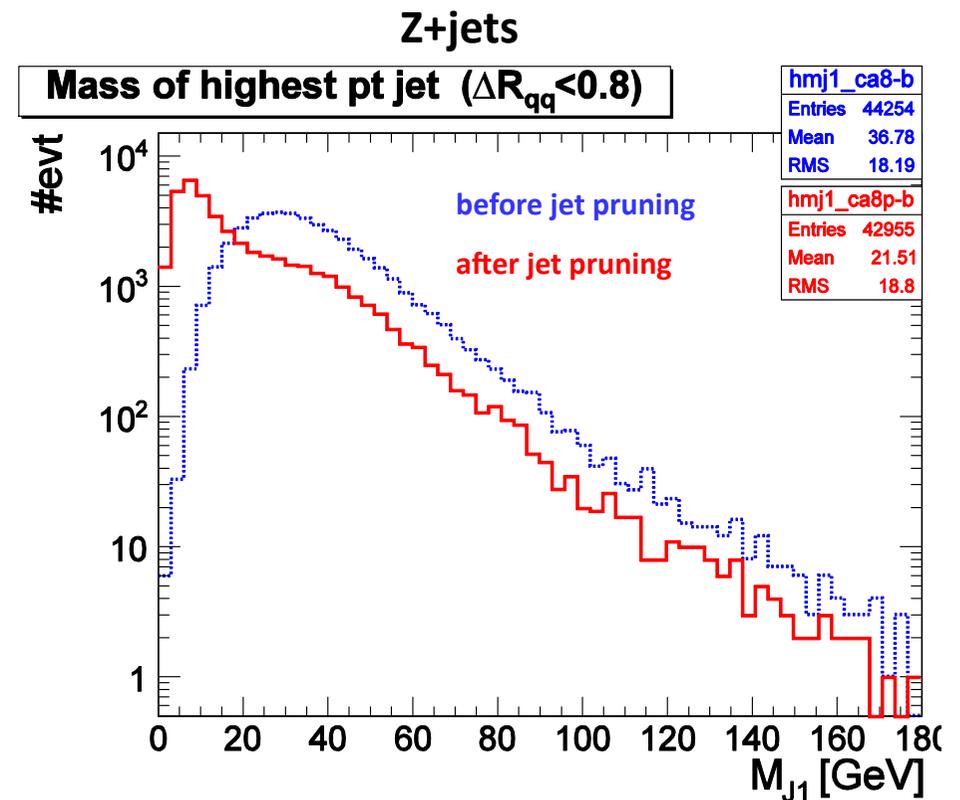
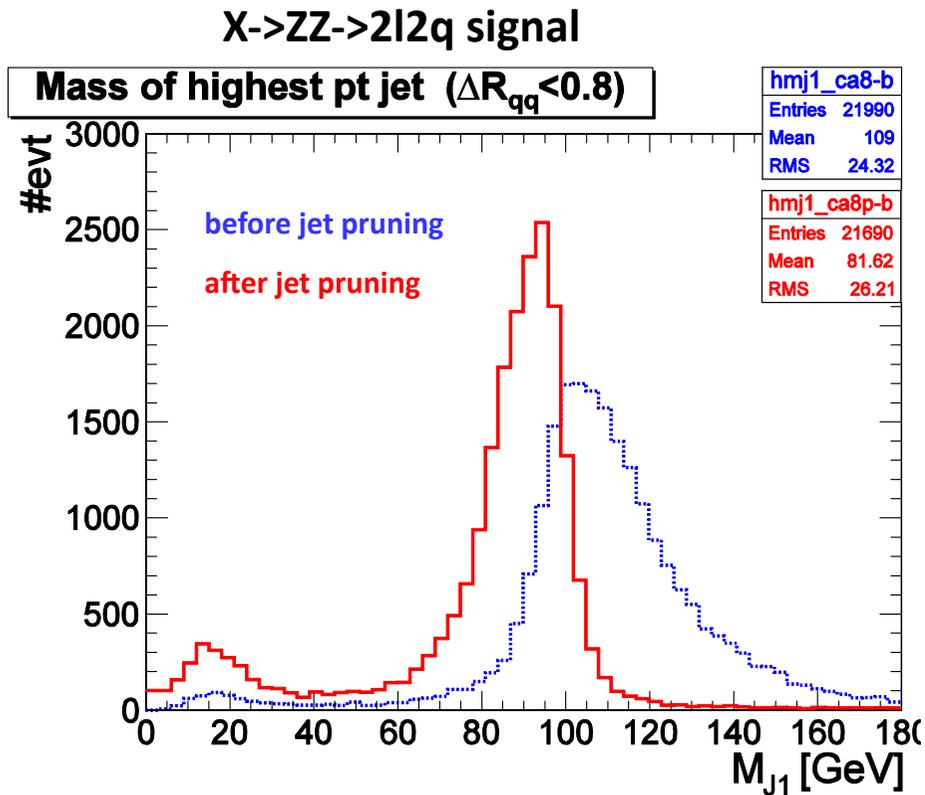
Typically used for boosted top or boosted $H \rightarrow bb \dots$



Example in $X \rightarrow ZZ \rightarrow 2l2j$

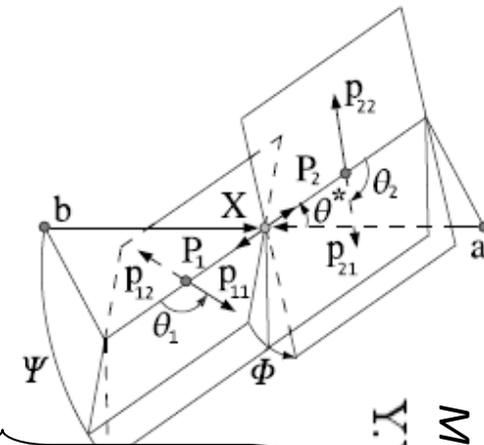
First look at Z boosted (no numbers yet) ...
preliminary, A.Bonato, R.Covarelli

- RS Graviton
- M_G 1500 GeV
- CA 0.8



Angular analysis (1)

□ $X \rightarrow ZZ \rightarrow 4f$ decay kinematic fully defined by 5 angles

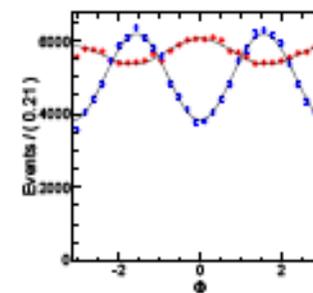
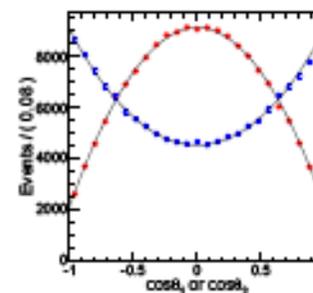
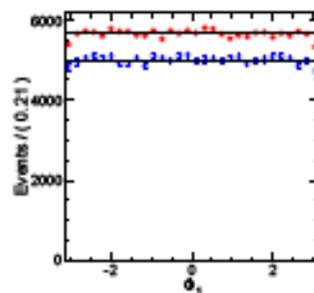
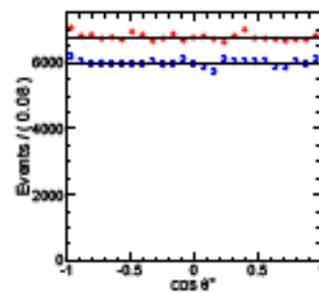


signal (M_X 250):

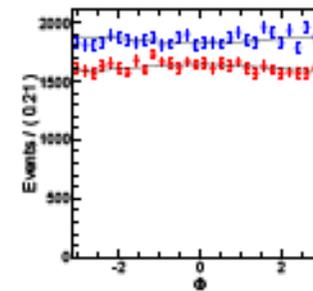
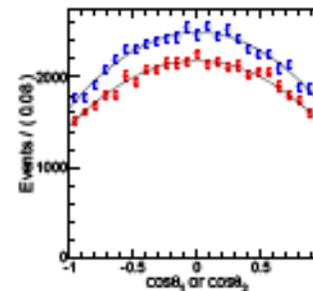
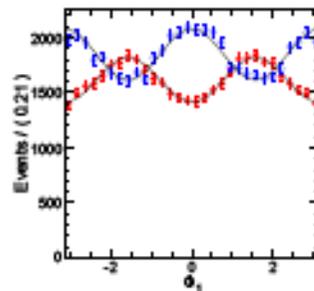
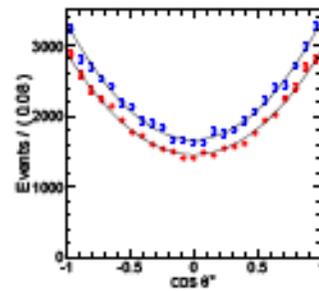
$X \rightarrow ZZ$

Z decays

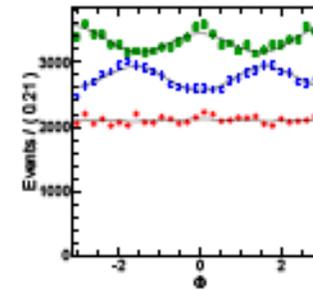
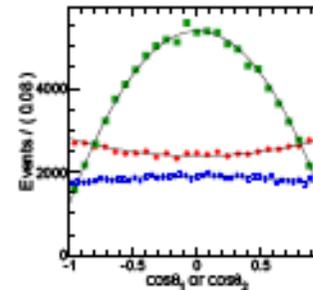
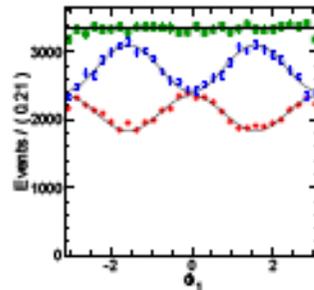
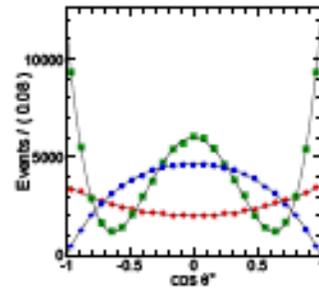
$0^+, 0^-$



$1^+, 1^-$



$2^+_{m=0}, 2^+_{L}, 2^-$



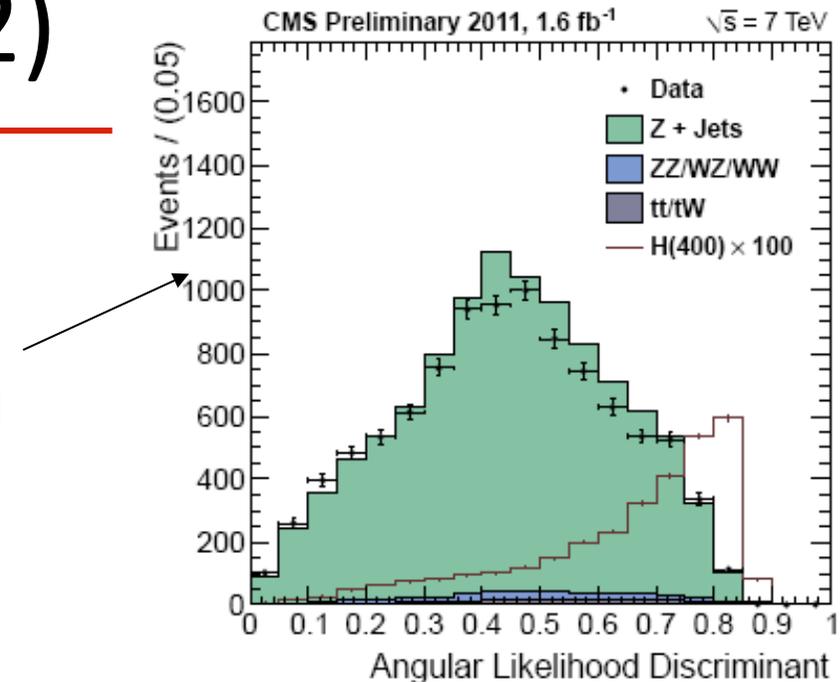
MC from Johns Hopkins
 Y.Y. Gao *et al.*, Phys. Rev. D **81**, 075022 (2010).

Angular analysis (2)

❑ Can be clearly used to disentangle different signals... but **what about background?**

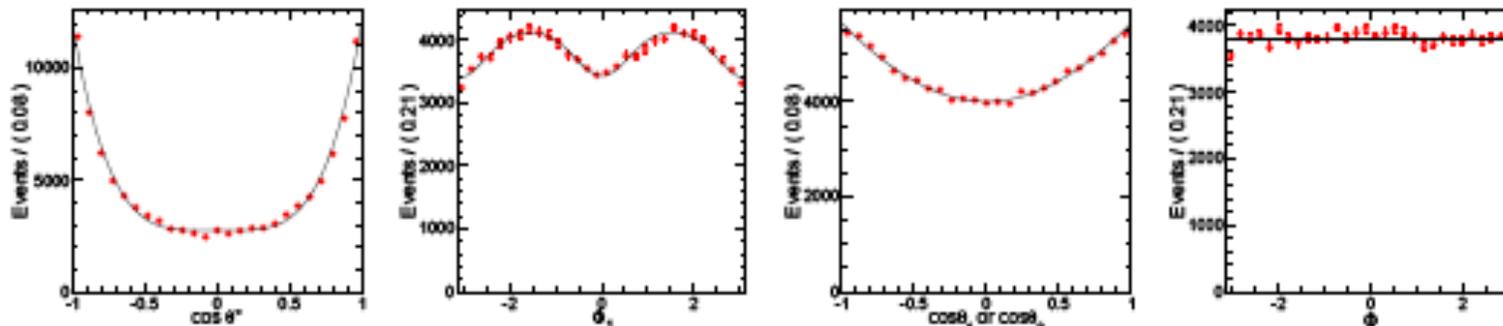
❑ Already used in $H \rightarrow ZZ \rightarrow 2l2q$: cut on likelihood

- signal: ideal \times uncorr. accept
- Z+jets from MC: no correlations, (background from jj sidebands)



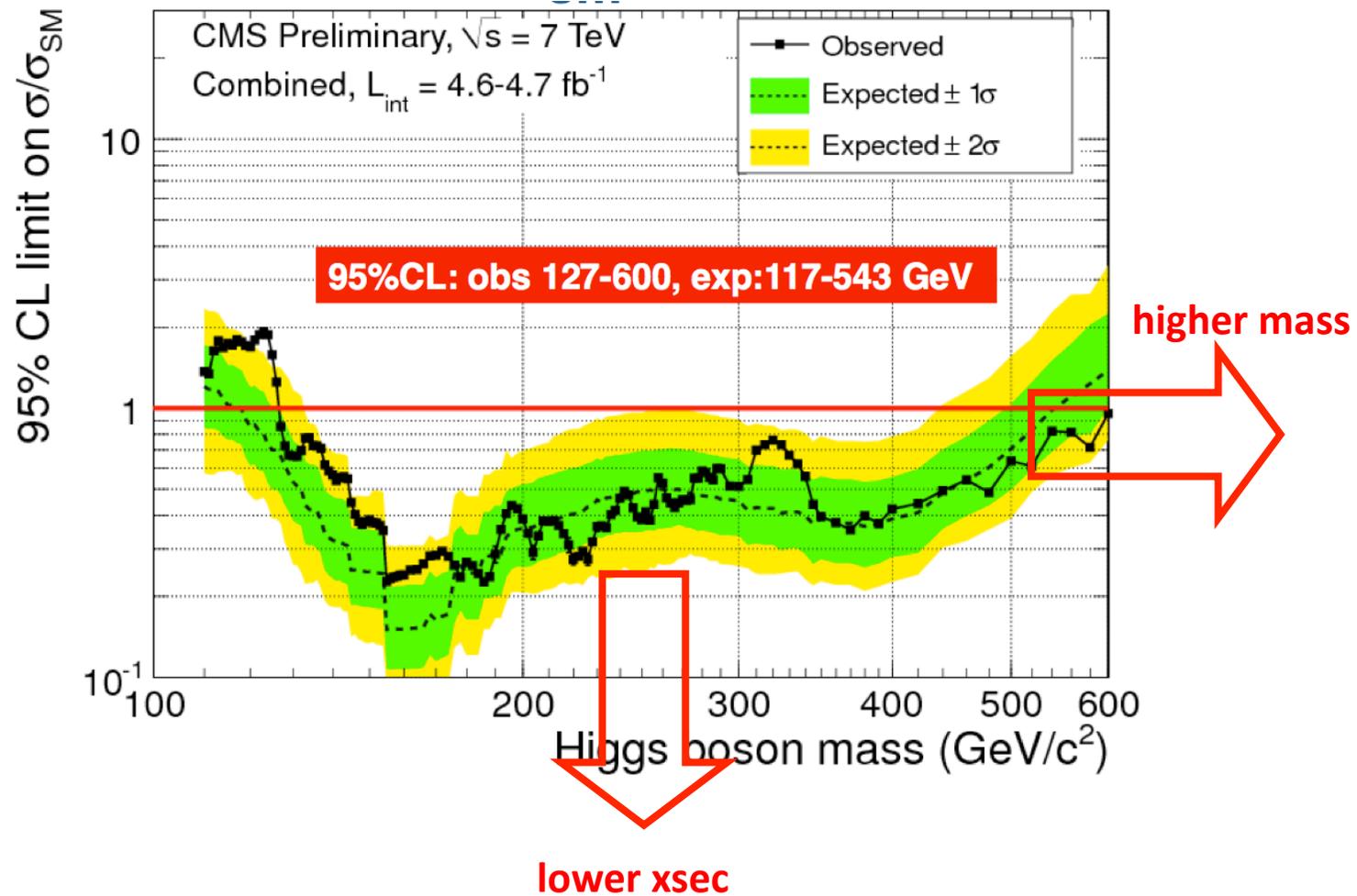
❑ To optimize further (**multidimensional fit**), need full theoretical description of background:

▪ $qq \rightarrow ZZ$:



▪ gg also available \rightarrow can be used to disentangle qq-gg!!

What's next ?



Improve sensitivity

First LHC to Terascale Workshop (Sept 2011):
LCH at LHC by J.R. Espinoza

WHY? Models with lower xsec

Ex of (light) composite higgs:

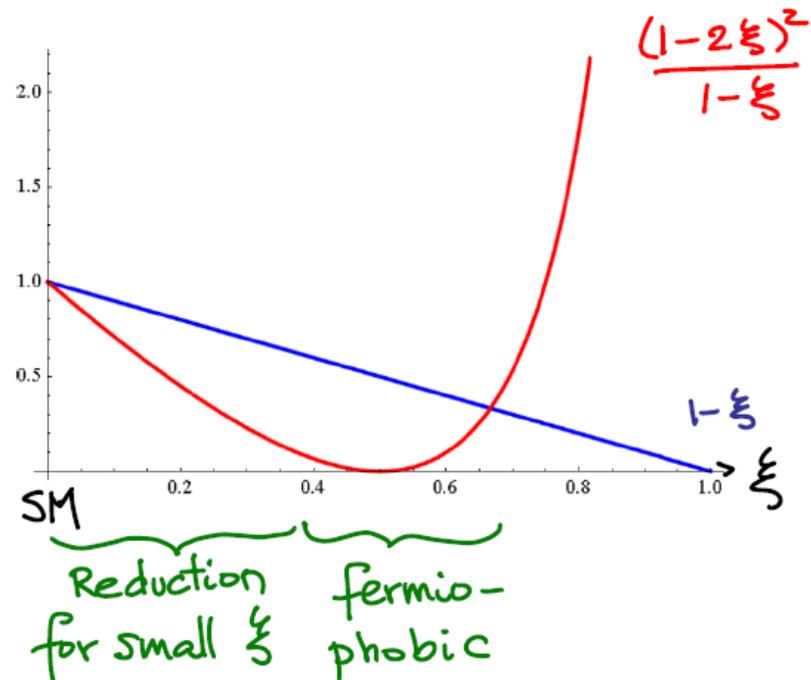
$$\frac{1}{p^2 - m_H^2} \rightarrow \frac{1}{1 + \xi_{CH}} \frac{1}{p^2 - m_H^2}$$

$$h \dots \begin{matrix} \nearrow V \\ \searrow V \end{matrix} \xrightarrow{\sqrt{1-\xi} \times (SM)} \Rightarrow \sigma \left[\begin{matrix} \nearrow \dots \\ \searrow \dots \end{matrix} \right] = (1-\xi) \sigma_{SM}$$

$$h \dots \begin{matrix} \nearrow f \\ \searrow f \end{matrix} \xrightarrow{\frac{1-2\xi}{\sqrt{1-\xi}} \times (SM)} \Rightarrow \begin{matrix} h \dots \begin{matrix} \nearrow g \\ \searrow g \end{matrix} \\ h \dots \begin{matrix} \nearrow g \\ \searrow g \end{matrix} \end{matrix} \xrightarrow{\frac{1-2\xi}{\sqrt{1-\xi}} \times (SM)}$$

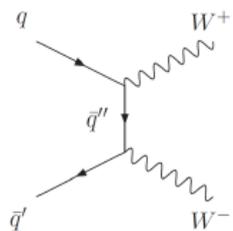
HOW?

- Factor 5 in luminosity wrt to present results
- Improve theoretical control of
 - signal: \rightarrow NNLO&NNLL effects, precise mass shape prediction, signal-background interference (back-up)
(studied in the Higgs Xsec WG and documented in 2 Yellow Report)
 - background: \rightarrow control of ZZ, WW ewk continuum

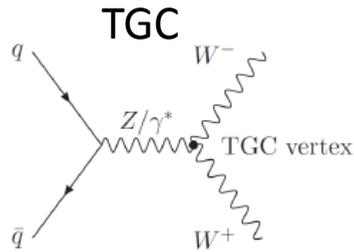


Diboson production (WW,WZ,ZZ)

qqbar \rightarrow VV



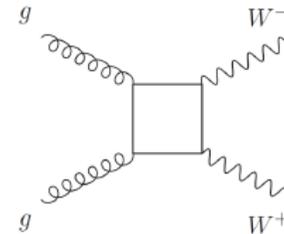
WW,WZ,ZZ



forbidden for ZZ

+ NLO qqbar +

gg \rightarrow VV



forbidden for ZW

(LHC: few % of xsec with ~50% uncertainty)

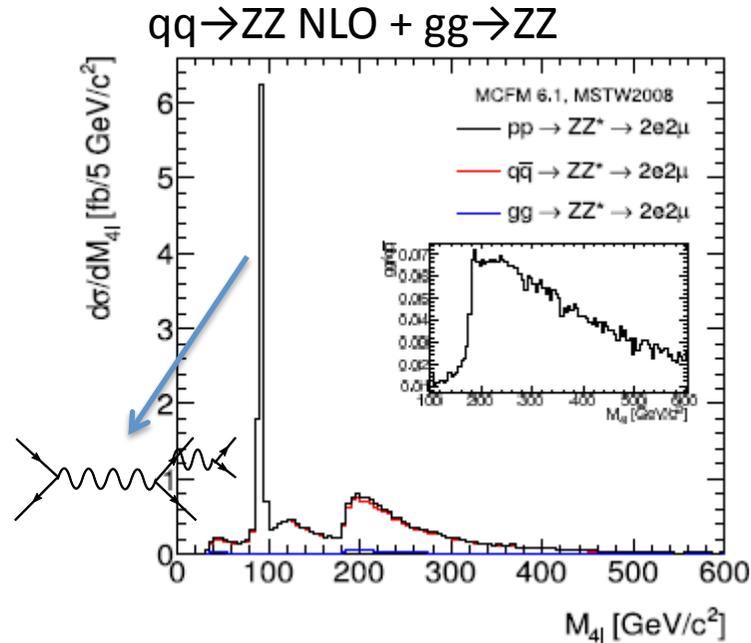
□ SM test: TGC fixed by ewk gauge structure

\rightarrow any deviation from SM in VV xsec is direct hint of NP in bosonic sector

□ Backgrounds for high mass Higgs \rightarrow VV

LHC focused on leptonic final state, Tevatron looked at semileptonic but limited by systematics (V+jets)

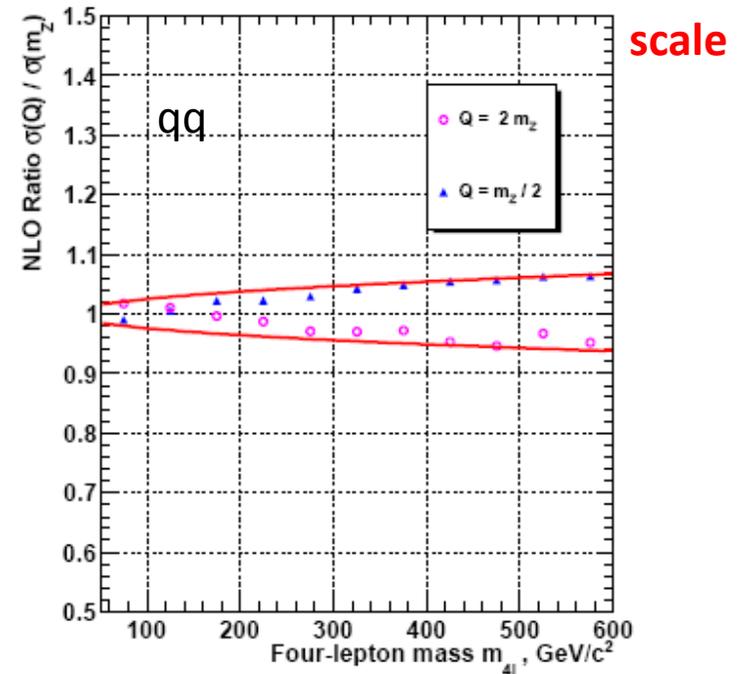
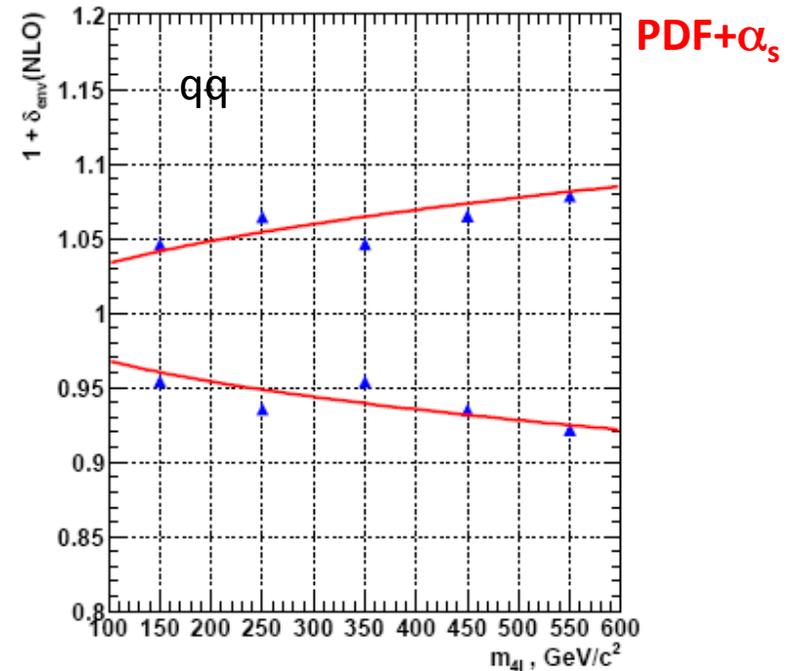
VV: theoretical prediction



□ Uncertainty dominated by QCD part

□ WW in jet bins: uncertainty on $\sigma(\geq N)$ + modeling: MC@NLO vs ALPGEN

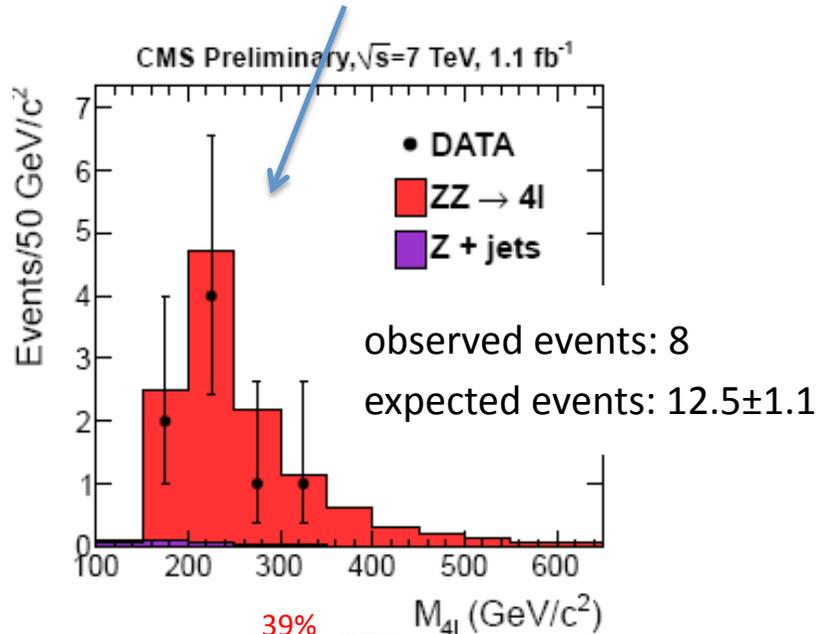
$\Delta\sigma_{\geq 0}$ (%)	$\Delta\sigma_{\geq 1}$ (%)	$\Delta\sigma_{\geq 2}$ (%)	$\Delta\sigma_{\geq 3}$ (%)
3	6	42	100



ZZ → 4l: measurement

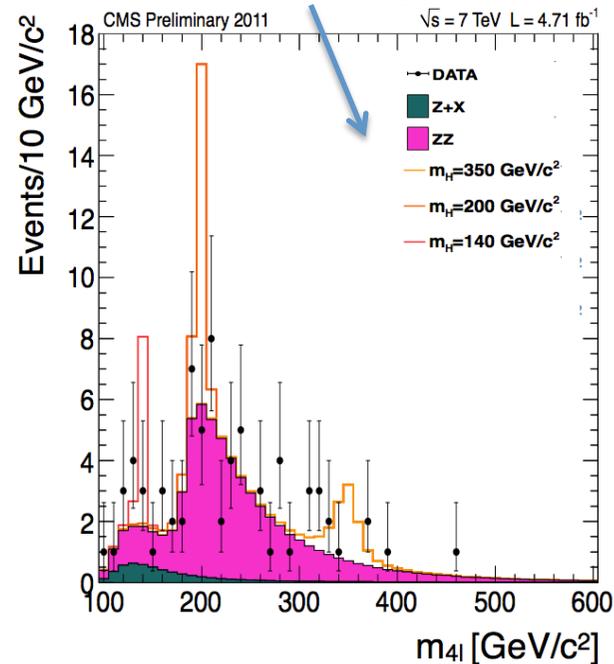
- 4l is 0.5% of ZZ xsec but **very clean**

Dedicated EWK analysis with very low luminosity, Higgs results much beyond that



$$\sigma(ZZ) = 3.8^{+1.5}_{-1.2} \text{ (stat)} \pm 0.2 \text{ (syst)} \pm 0.2 \text{ (lumi)} \text{ pb}$$

39%



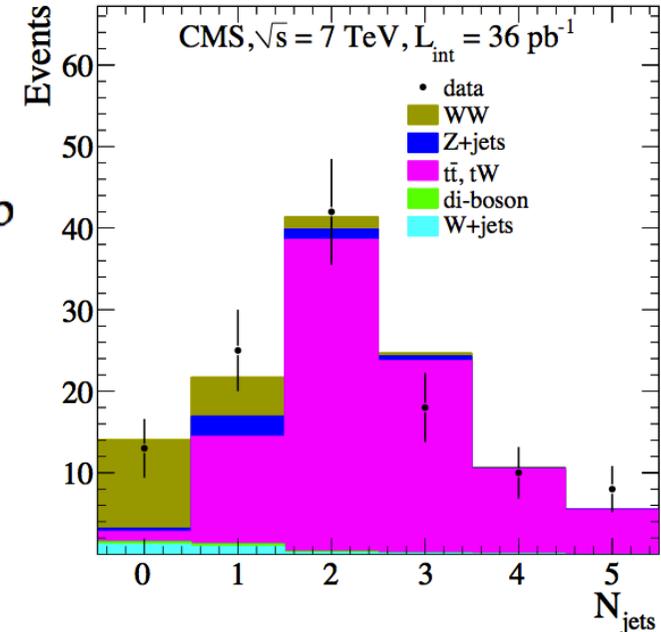
$$\sigma(pp \rightarrow ZZ + X) \times \mathcal{B}(ZZ \rightarrow 4\ell) = 28.1^{+4.6}_{-4.0} \text{ (stat.)} \pm 1.2 \text{ (syst.)} \pm 1.3 \text{ (lumi.)} \text{ fb}$$

16%

WW->lnln: measurement

- Dedicated EWK results only with very low luminosity,

$$\sigma_{W+W^-} = 41.1 \pm 15.3 \text{ (stat)}^{37\%} \pm 5.8 \text{ (syst)} \pm 4.5 \text{ (lumi)} \text{ pb}$$



- Higgs analysis much beyond that:

	data	all bkg.	$qq \rightarrow W^+W^-$	$gg \rightarrow W^+W^-$	$t\bar{t}+tW$	W + jets
0-jet	1359	1364.8 ± 9.3	980.6 ± 5.2	58.8 ± 0.7	147.3 ± 2.5	99.3 ± 5.0
1-jet	909	951.4 ± 9.8	416.8 ± 3.6	23.8 ± 0.5	334.8 ± 3.0	74.3 ± 4.6
2-jet	703	714.8 ± 13.5	154.7 ± 2.2	5.1 ± 0.2	413.5 ± 2.7	37.9 ± 3.6

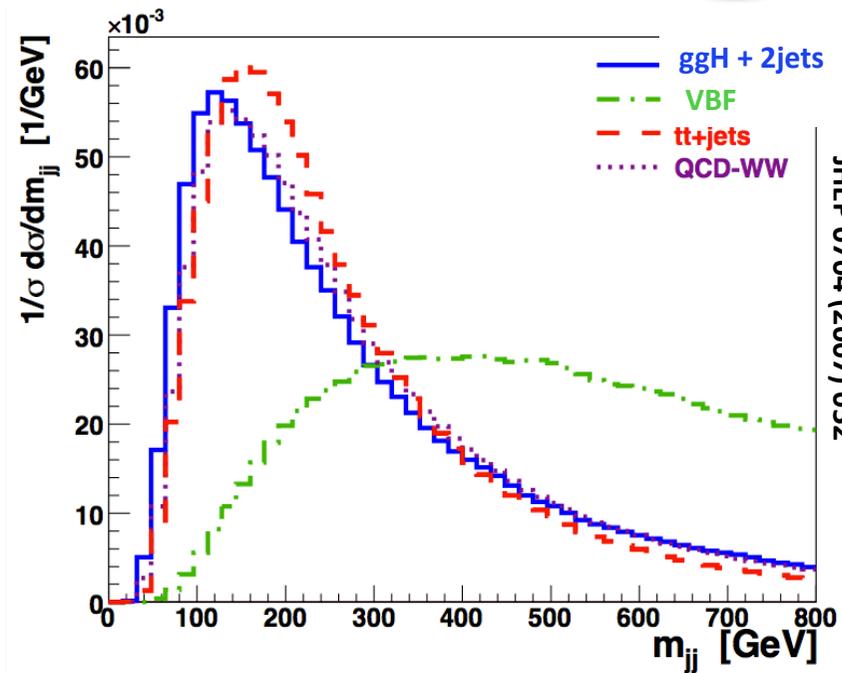
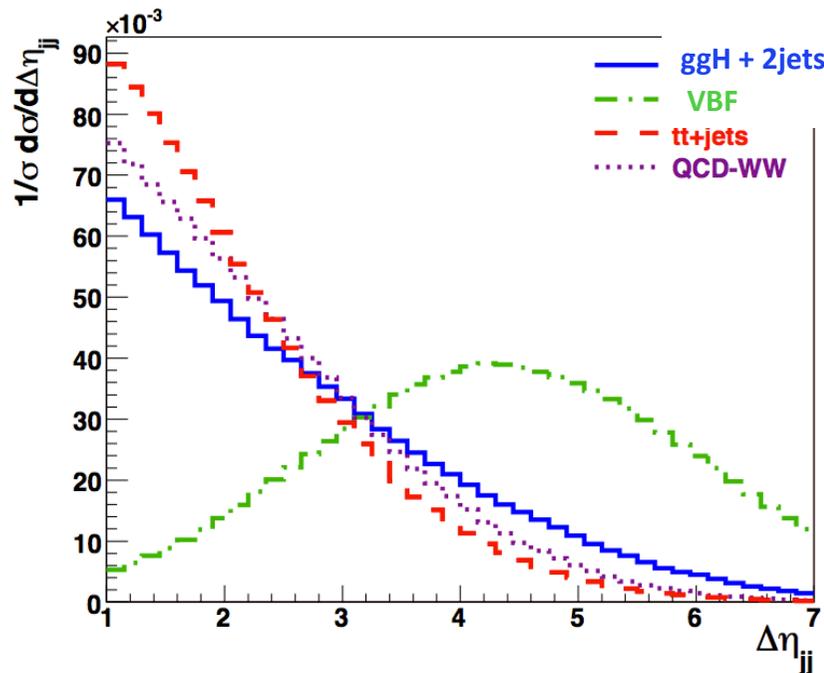
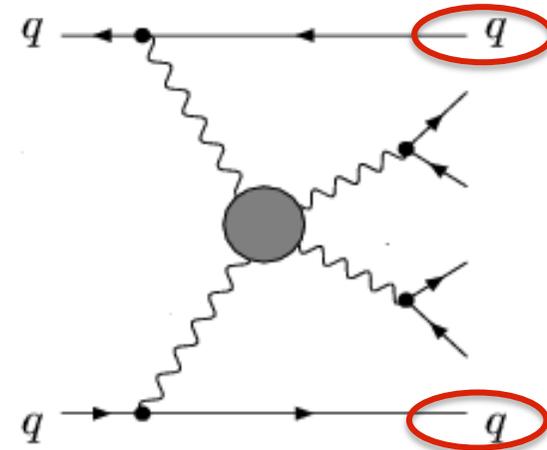
	WZ/ZZ	$Z/\gamma^* \rightarrow \ell^+\ell^-$	$W\gamma$	$Z/\gamma^* \rightarrow \tau^+\tau^-$
0-jet	33.0 ± 0.5	16.6 ± 4.0	26.8 ± 3.5	2.4 ± 0.5
1-jet	28.7 ± 0.5	39.4 ± 6.4	13.0 ± 2.6	20.6 ± 0.4
2-jet	15.1 ± 0.3	56.1 ± 11.7	10.8 ± 3.6	21.6 ± 2.1

stat. and syst. errors included

From VBF to VV scattering

- ❑ First search for a VBF resonance, feasible in 2012
- ❑ Measurement of VV scattering spectrum with higher lumi ($>50 \text{ fb}^{-1}$)

Typical signature: forward-backward “spectator” jets with very high energy



JHEP 0704 (2007) 052

Higgs-like resonance in VBF

□ RE-DO all the analyses in VBF mode

□ Today only WW→lnln. **Expectations for next year:**

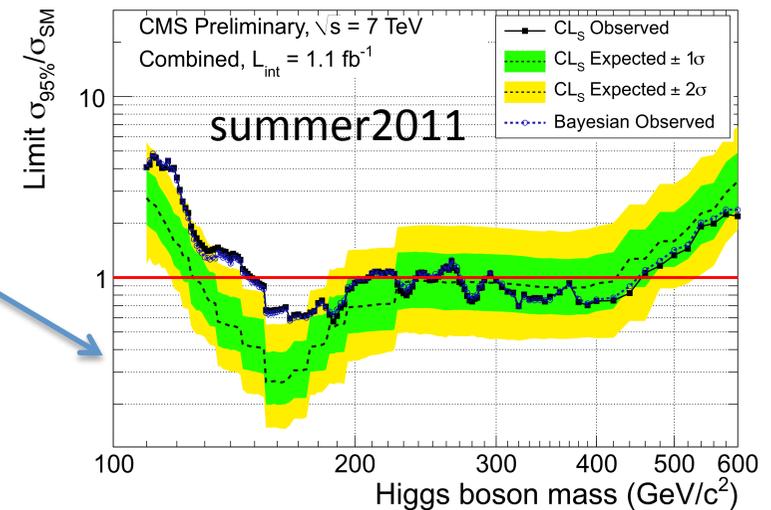
- lumi > 10 fb⁻¹
- $\sigma(\text{vbf}) \sim 0.1 \times \sigma(\text{gg})$
- 0.5 effic. VBF cuts

VBF yields in 2012 ~
0.5 gg yields of 2011
summer results,

with much less background:

- ZZ→4l will be still limited by statistics
- WW→lnln will improve S/B (signal/10, WW* α_s^2)
- **semileptonic final states** will have **reasonable signal yields + much lower background** than inclusive analysis

- eg, ZZ→lljj :
- signal yields for m_H 300-500 ~ 15 – 5 events
 - V+(N+1)jets/V+N jets ~ 0.15 → asking 2 jets reduces **background to 2%!**
 - S/B may increase of a factor 2 (eff 0.5 × σ 0.1 / 0.02)

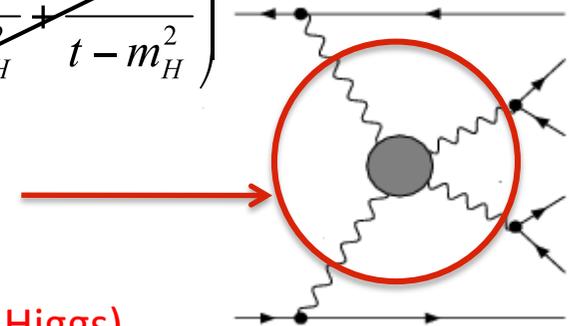


VV scattering spectrum

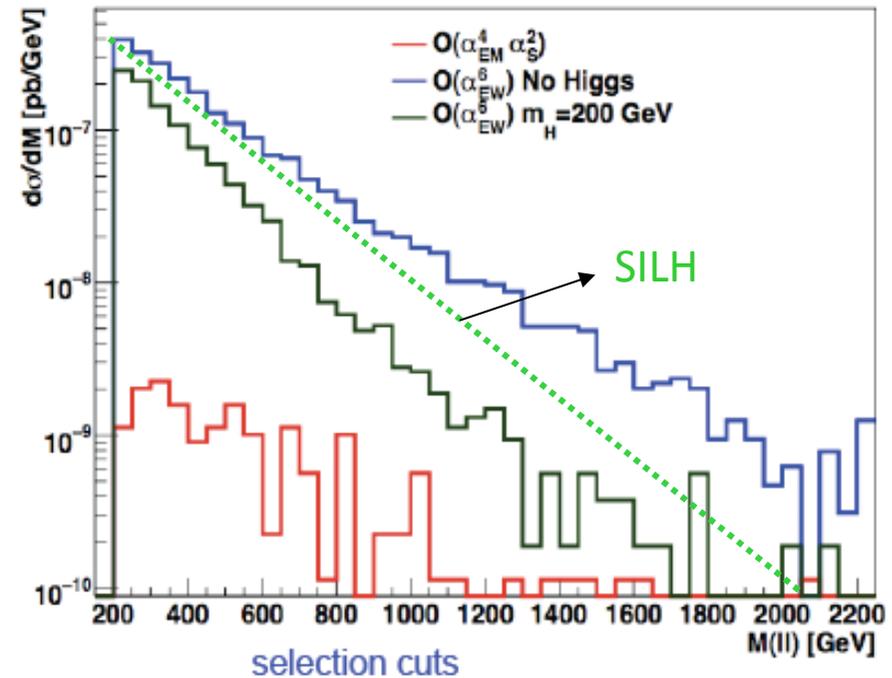
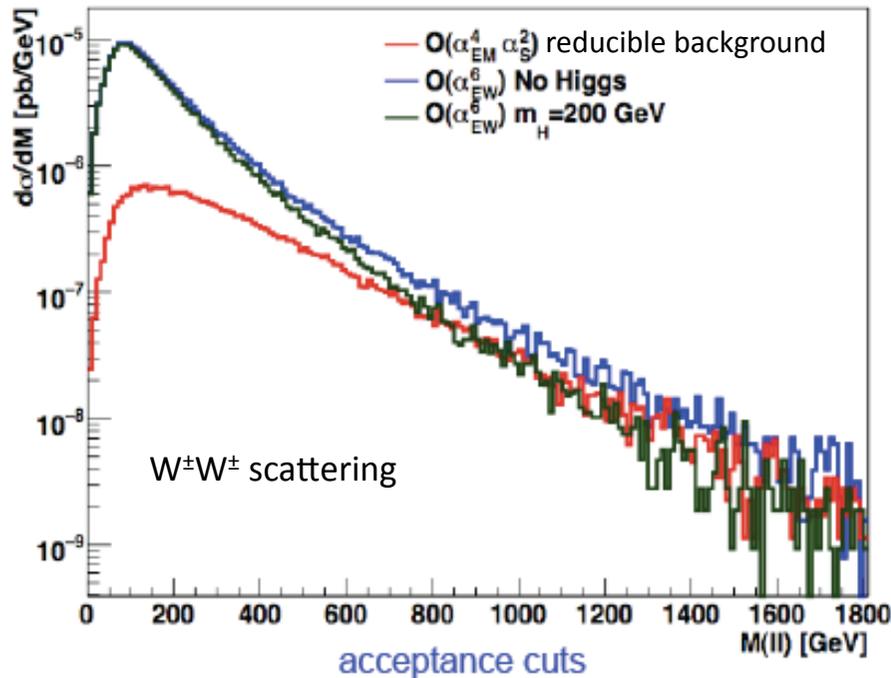
□ In no Higgs case: $A(W_L^+ W_L^- \rightarrow W_L^+ W_L^-) \approx \frac{1}{v^2} \left(-s - t + \frac{s^2}{s - m_H^2} + \frac{t^2}{t - m_H^2} \right)$

BUT increasing of xsec at high VV is suppressed by

- PDF
- offshell bosons
- unpolarized bosons

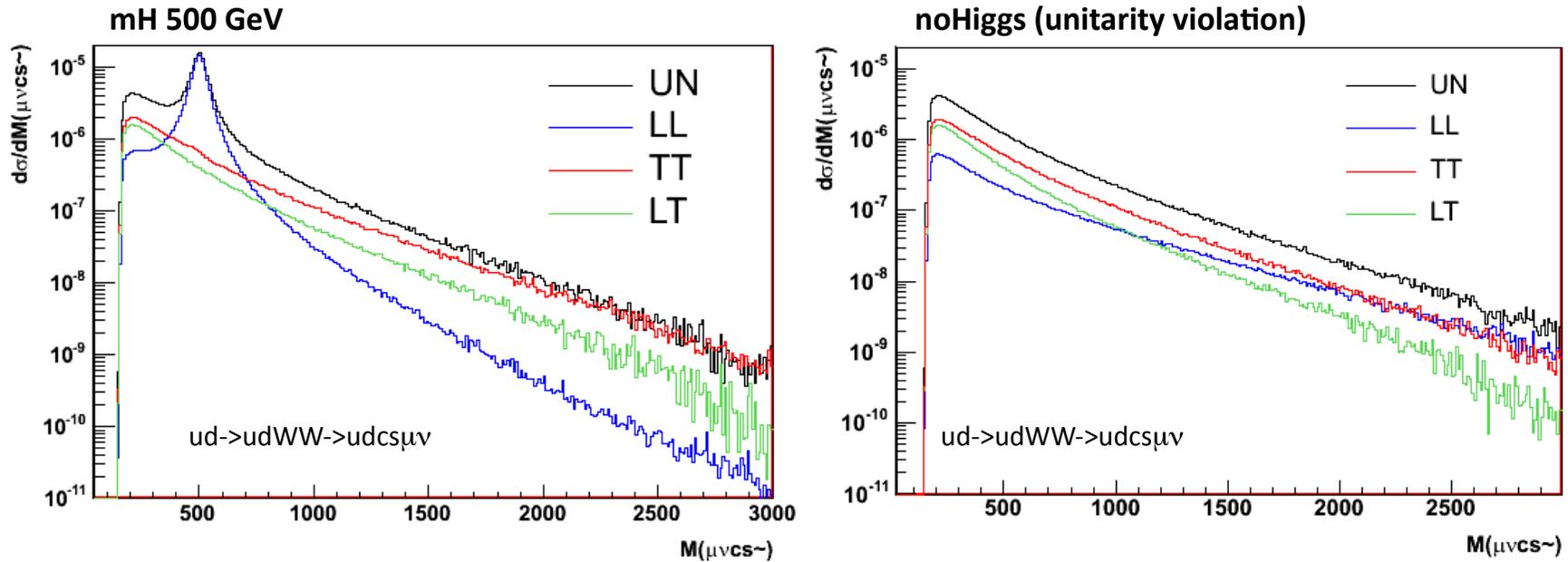


→ small difference btw SM and violation of unitarity (no Higgs)



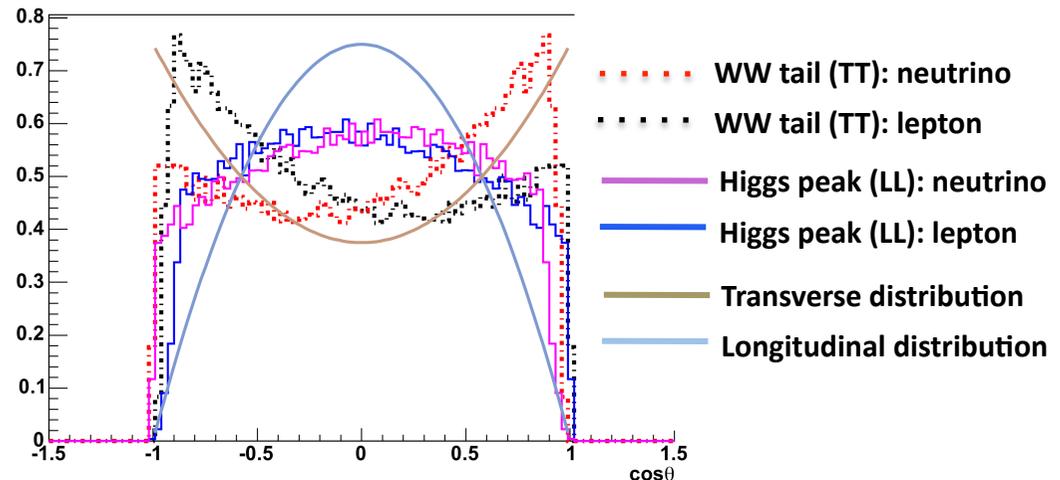
→ with proper cut (eg $\Delta\eta$ jets) can be enhanced → selection of the longitudinal W

Longitudinal polarization



Angular analysis can boost LL-TT separation (new!):

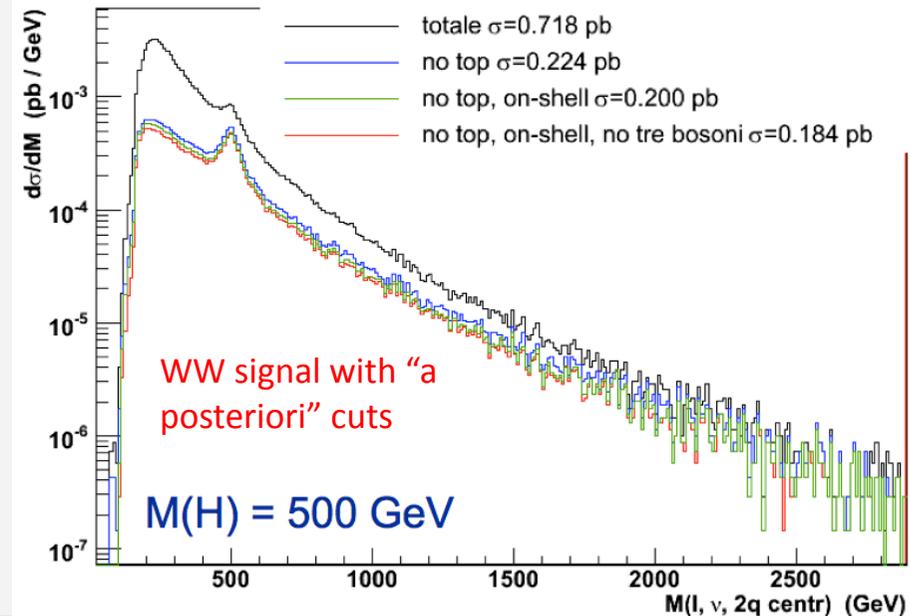
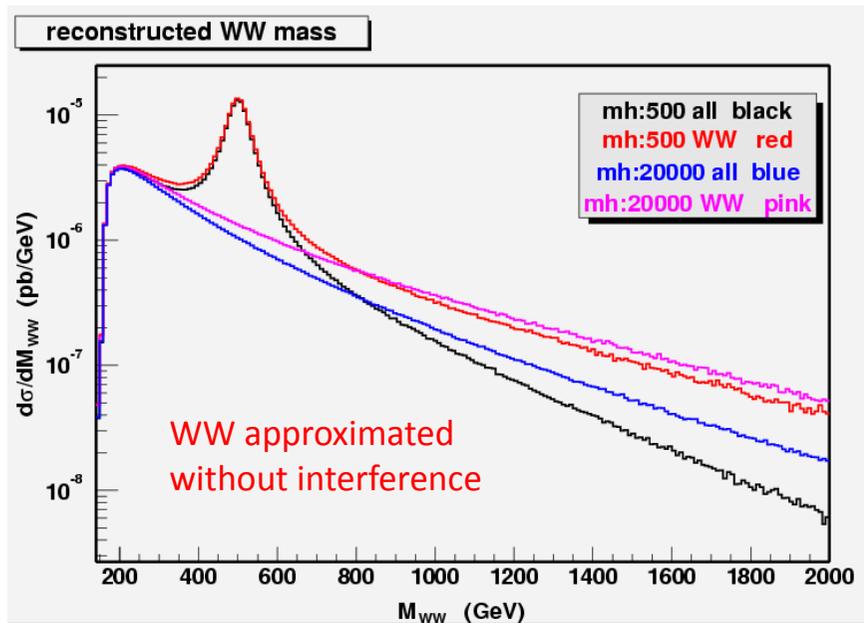
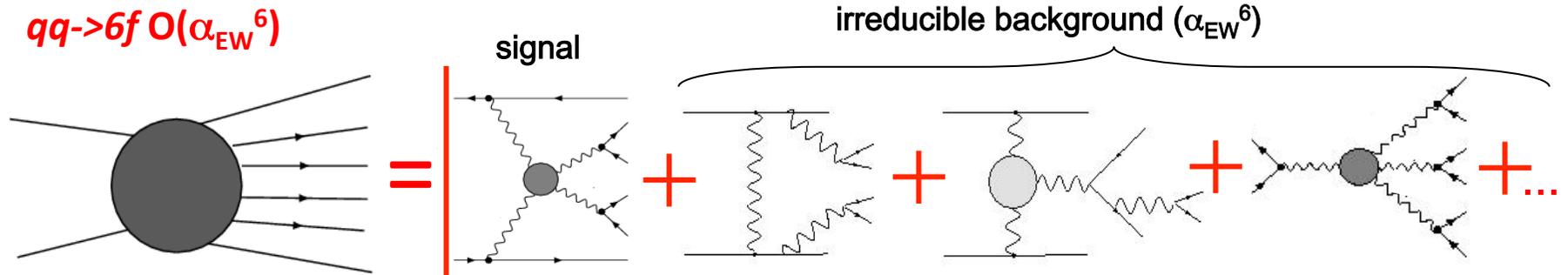
partonic study in the center of mass of W



VV scattering: interference effects

Big interference effects considered only in **Phantom**

2



Summary

- ❑ The first aim of the Higgs search is **the understanding of the EWSB**
-> focus on H->**VV final state**

- ❑ Main steps:

- search for **generic resonance $X \rightarrow ZZ \rightarrow 4f$**
 - search for **VBF resonance** and measuring of **VV scattering spectrum**
- } angular analysis

- ❑ Ingredients along the EWSB road:

- **control of V+jets (jet pruning)**
- **control uncertainties on VV EWK continuum**

The road to the ElectroWeak Symmetry Breaking

BACK-UP

18th January 2012

Seminar at Johns Hopkins University

S.Bolognesi (Johns Hopkins University)

Sources

CMS AN-2010-35: Angular Analysis of Resonances $pp \rightarrow X \rightarrow ZZ$

A. Bonato, A.V. Gritsan, Z.J. Guo, N.V. Tran, A. Whitbeck
Johns Hopkins University, Baltimore, MD, USA

—————→ *Y.Y. Gao et al.*, *Phys. Rev. D* **81**, 075022 (2010).

JHU seminar: **Path-Integral Jets** by David Krohn (Harvard)

www.pha.jhu.edu/groups/particle-theory/seminars/talks/F11/talk.khron.pdf

First LHC to Terascale Workshop (Sept 2011):

LCH at LHC by J.R. Espinoza

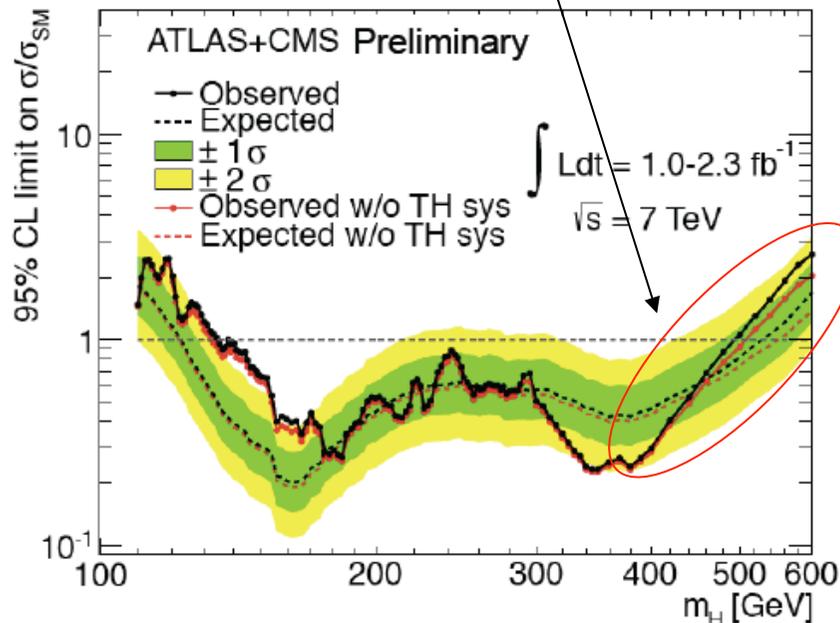
Boson Boson scattering analysis by A.Ballestrero (INFN Torino)

Mass shape

- Present approx:
 - xsec for **on-shell** Higgs production and decay in **zero width** approx
 - acceptance from MC with **ad-hoc BW distribution**



10-30% uncertainty on xsec for m_H 400–600 GeV



From Passarino talk at last LHC to Terascale WS

The off-shell Higgs production

is currently computed according to

$$\sigma_{os}(\mu_H^2) \delta(z \hat{s} - \mu_H^2) \implies \sigma_{OFS}(z \hat{s}) \text{BW}(z \hat{s}),$$

at least at lowest QCD order, where the *so-called* modified Breit–Wigner distributions is defined by

$$\text{BW}(s) = \frac{1}{\pi} \frac{s \Gamma_H^{OS} / \mu_H^2}{[s - \mu_H^2]^2 + (s \Gamma_H^{OS} / \mu_H^2)^2},$$

where now $\mu_H = M_H^{OS}$.

This ad-hoc Breit–Wigner

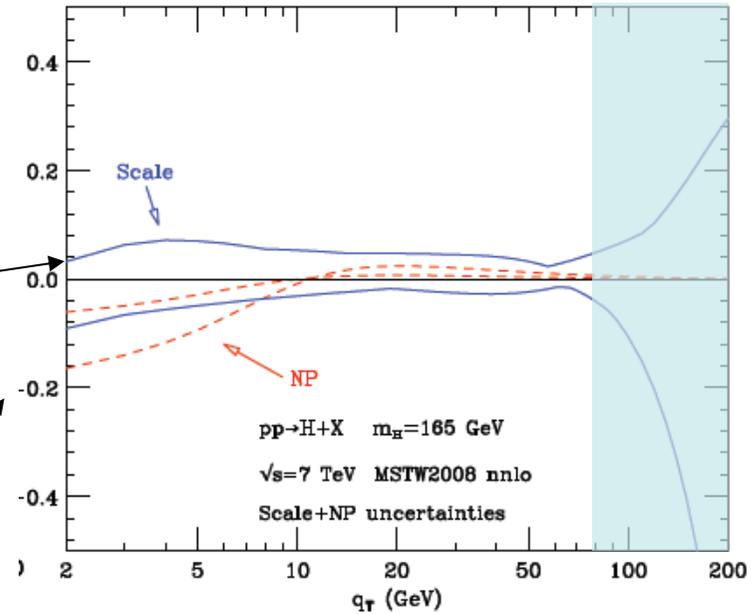
- cannot be derived from QFT and also is not normalizable in $[0, +\infty]$.
- Its practical purpose is to enforce a *physical* behavior for low virtualities of the Higgs boson but the usage cannot be justified.
- This modified Breit–Wigner cannot be derived from QFT.
- Note that this Breit–Wigner for a running width comes from the substitution of $\Gamma \rightarrow \Gamma(s) = \Gamma s / M^2$ in the Breit–Wigner for a fixed width Γ . This substitution is not justifiable.

Study with QFT-consistent Higgs propagator in the YR2

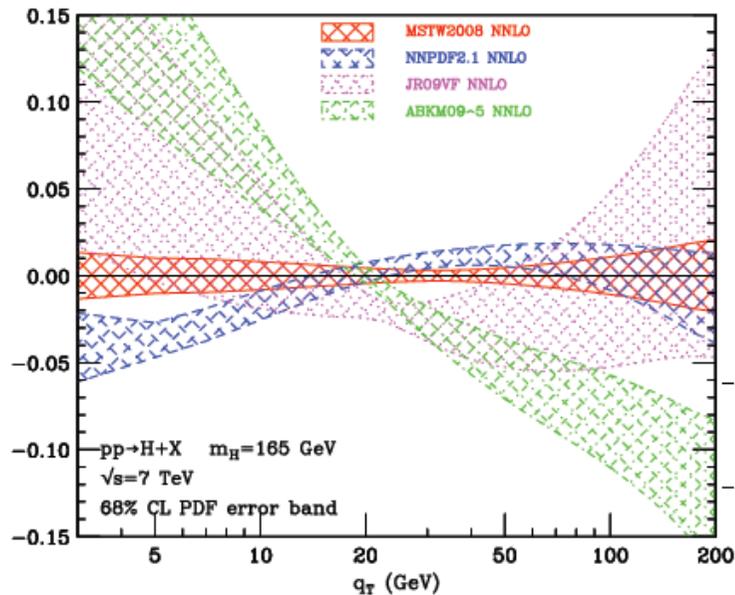
Higgs q_T

- **HqT:** $q_T > m_H$ NNLO
 $q_T \ll m_H$ NNLL (resumming $\ln(m_H^2/q_T^2)$)

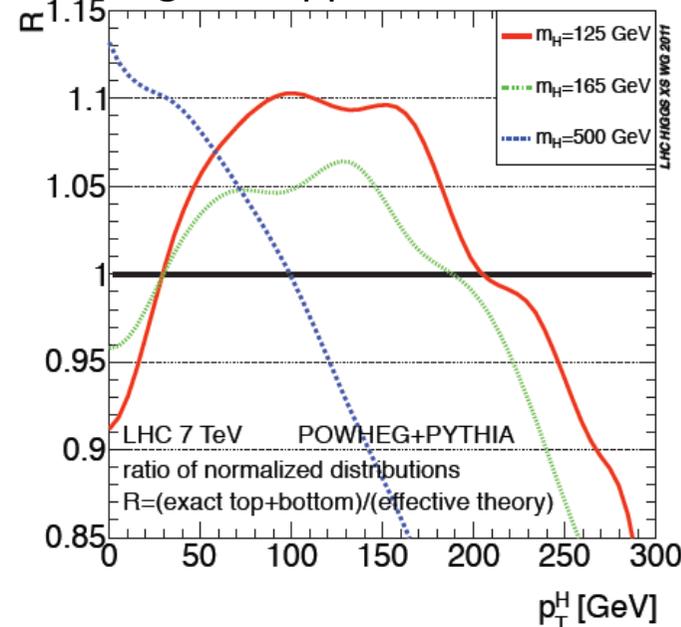
- Uncertainties:
- factor/renorm scale
 - non perturb. effects (smearing with NP form factor)



- **PDF**



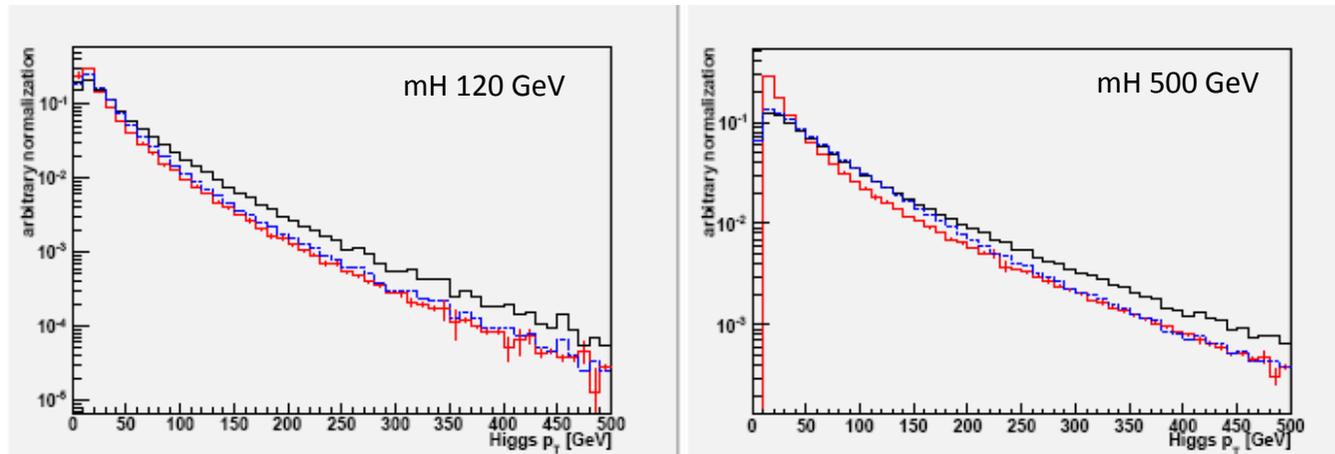
- **large m_t approximation**



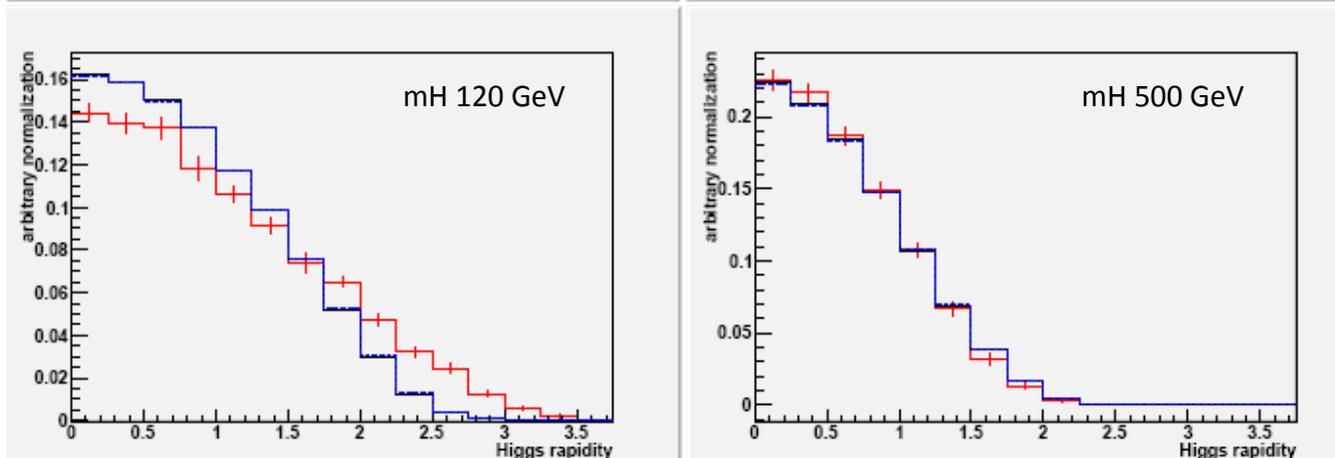
Reweight to HqT

- HqT used to reweight full event generators (POWHEG at NLO)

H p_T



H η



Powheg

Powheg re-weighted to HqT

(to be redone before PS)

HNNLO

- Very small effect on acceptance in 4l: 1-2% (larger if jet veto!)

Signal: jet counting

- Analysis in exclusive jet bins
(ex, WW+0,1,2 jets)

→ theoretical uncert in jet bins to be combined with correlations

- varying renormalization and factorization scales in the **fixed-order predictions for each exclusive jet cross section σ_N**

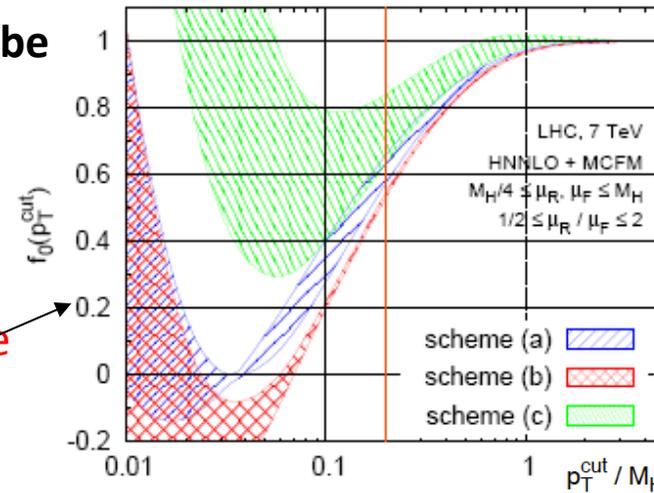
(results as 100% correlated)

- **inclusive xsec ($\sigma_{\geq N_{\text{jets}}}$), as source of perturbative uncertainties**

$$\sigma_N = \sigma_{\geq N} - \sigma_{\geq N+1}$$

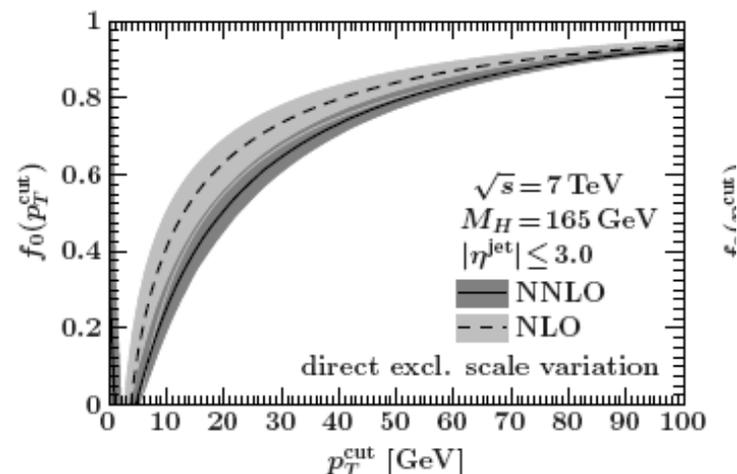
with error propagation

- if background depends on N_{jets}
- for VBF



different treatments of the uncontrolled higher-order $O(\alpha^3s)$ terms

i.e., different NNLO expansions



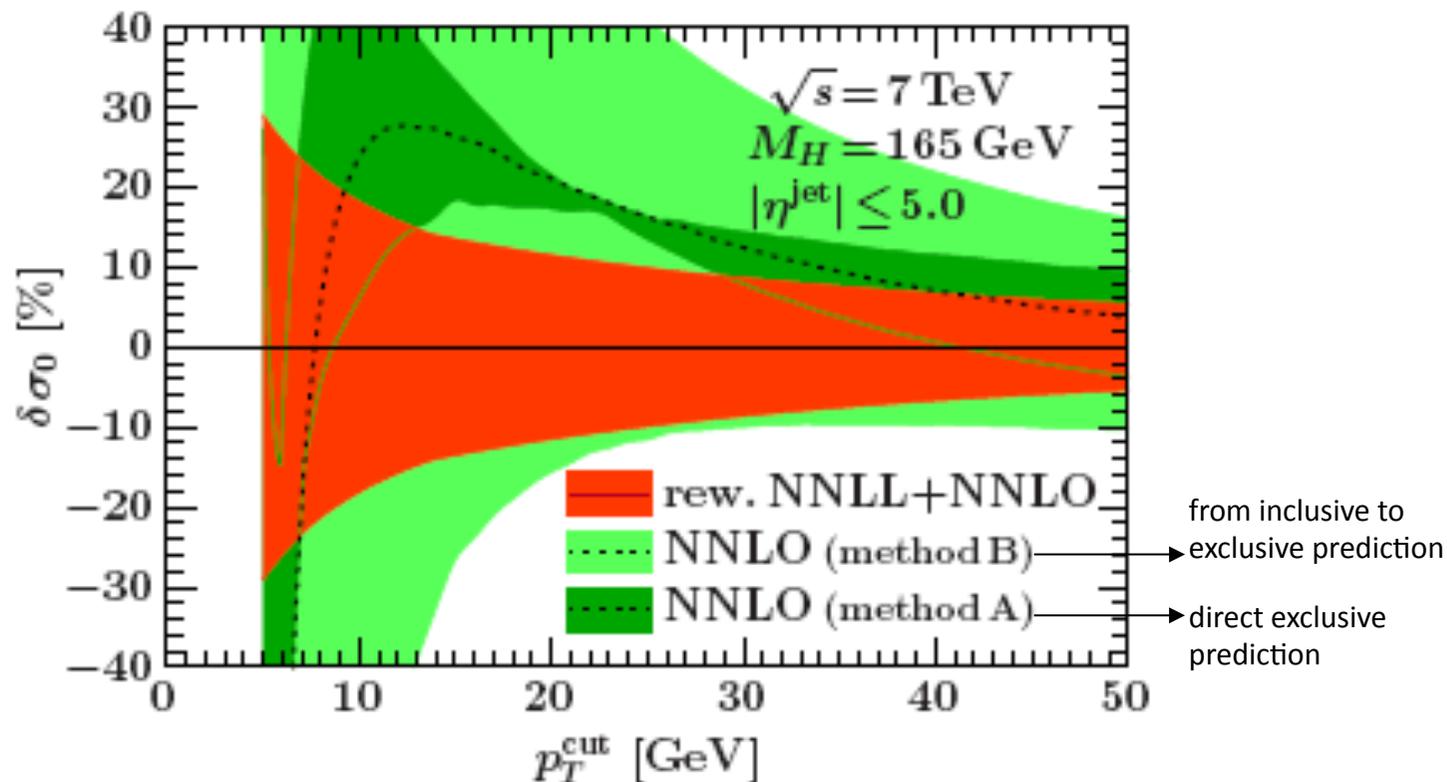
Signal: jet veto

- **Resummation of jet-veto logarithms** ($\ln(p_{\text{cut}}/m_H)$), induced by jet cut parameter p_{cut}

Presently doable only on beam thrust variable

$$\mathcal{T}_{\text{cm}} = \sum_k |\vec{p}_{Tk}| e^{-|\eta_k|} = \sum_k (E_k - |p_k^z|) \quad (\sim \text{raw approx of } p_{\text{cut}})$$

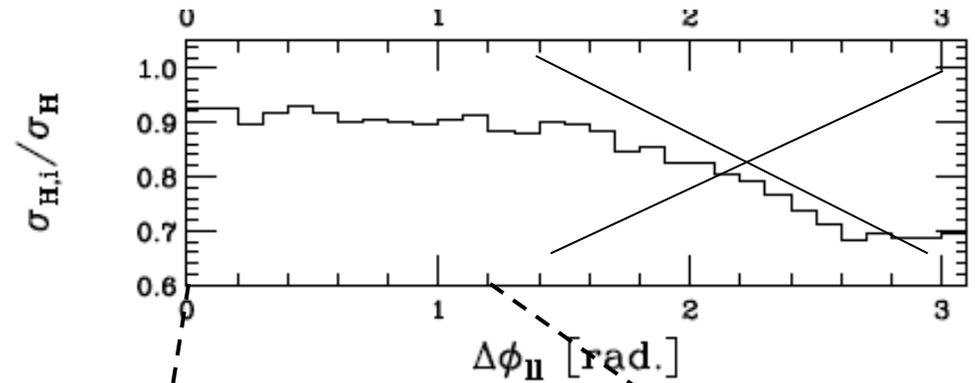
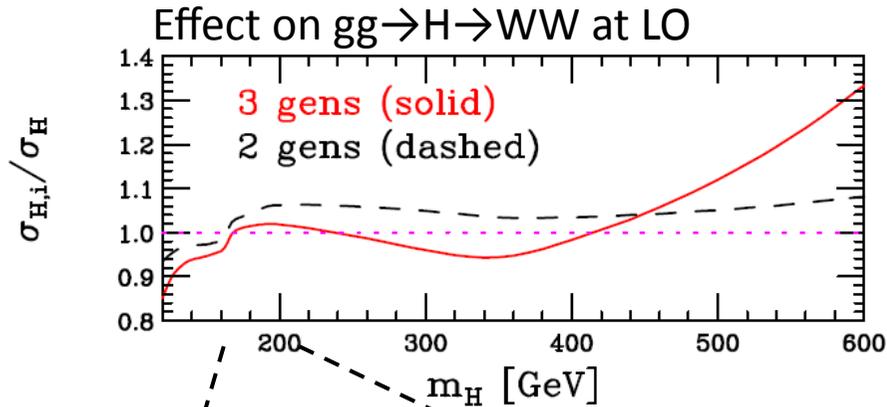
and used to reweight MC@NLO



Signal-background interference

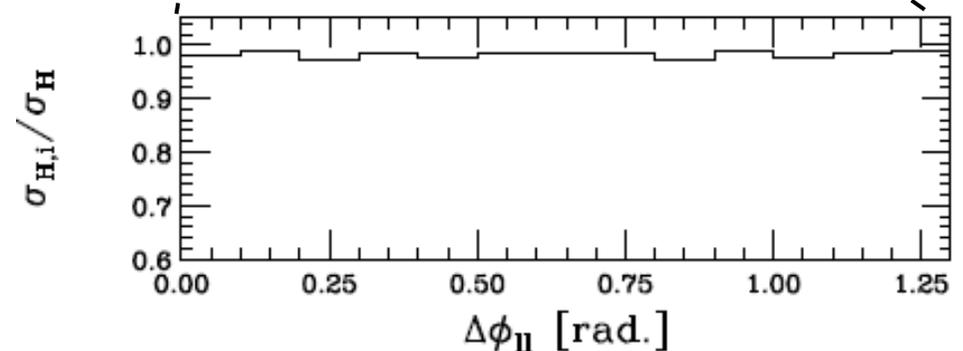
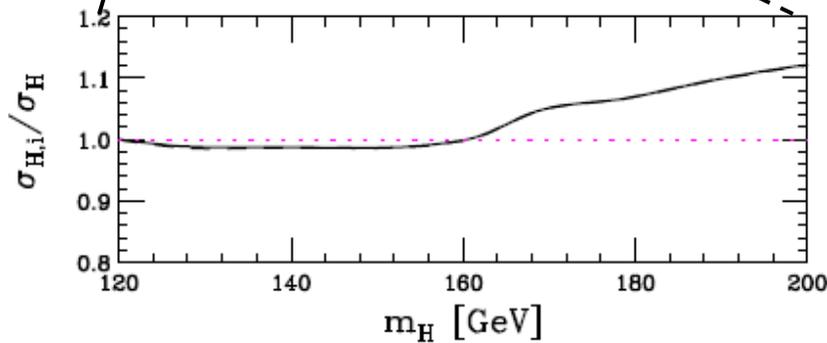
Recent results for WW, but focused on low mass

(arXiv:1107.5569v1)



$m_T < m_H$ non-resonant diagrams can be large for $m_T > m_H$

also shape effects!

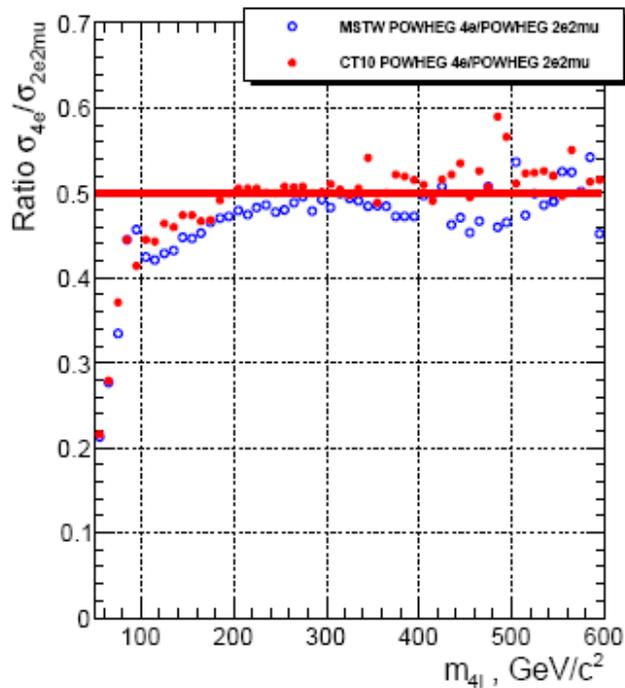


Worth to investigate further at high mass?

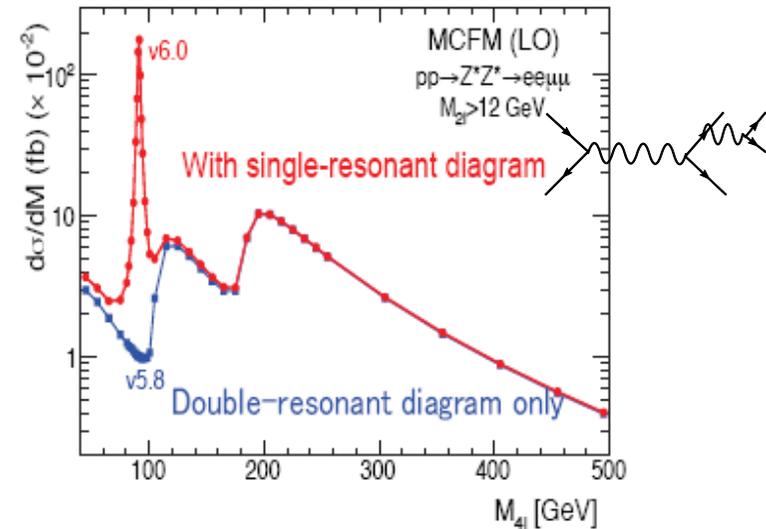
ZZ: theoretical prediction

- ZZ fully from MC, well under control

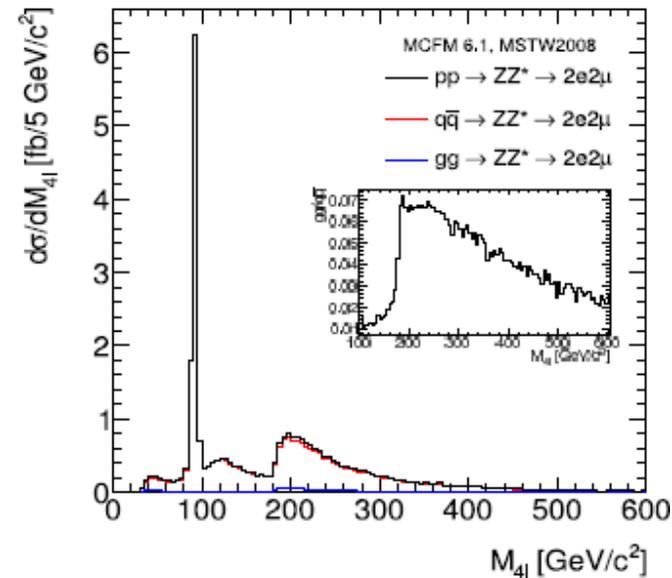
Interference in the final state with identical leptons



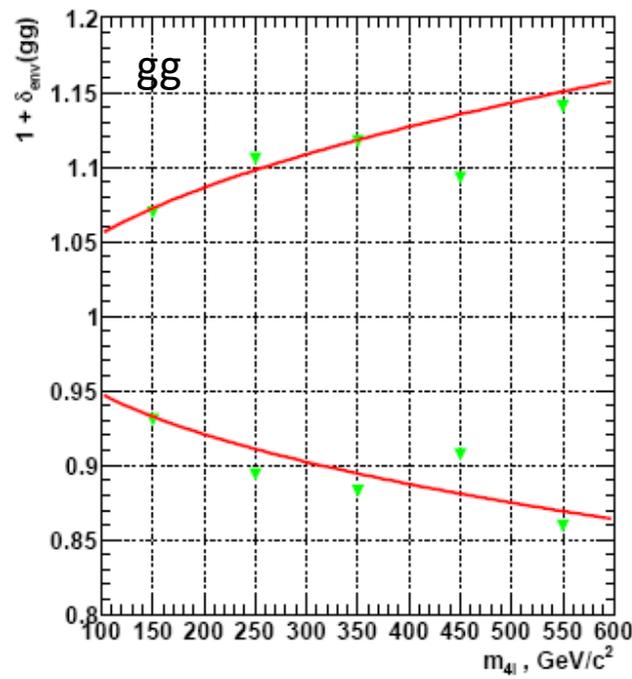
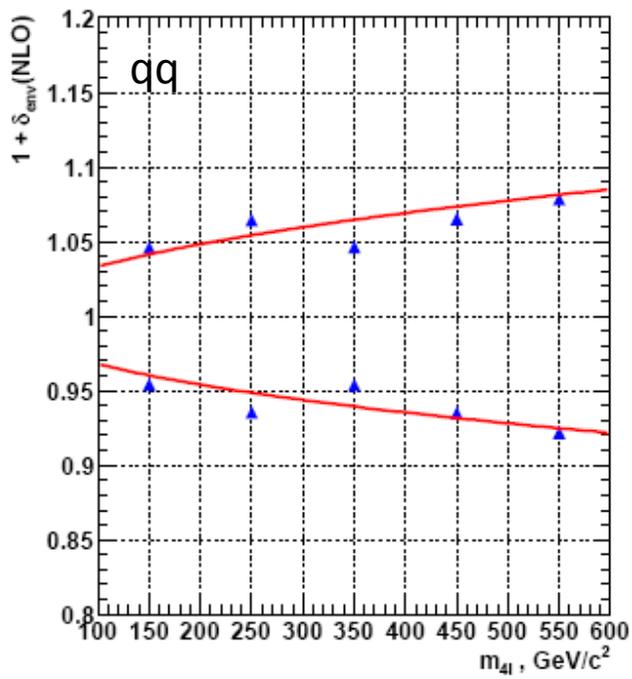
Single resonant contribution



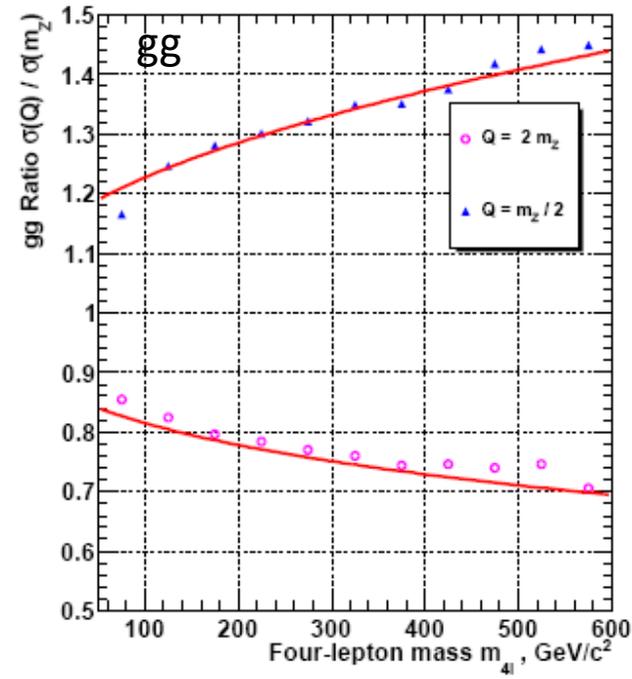
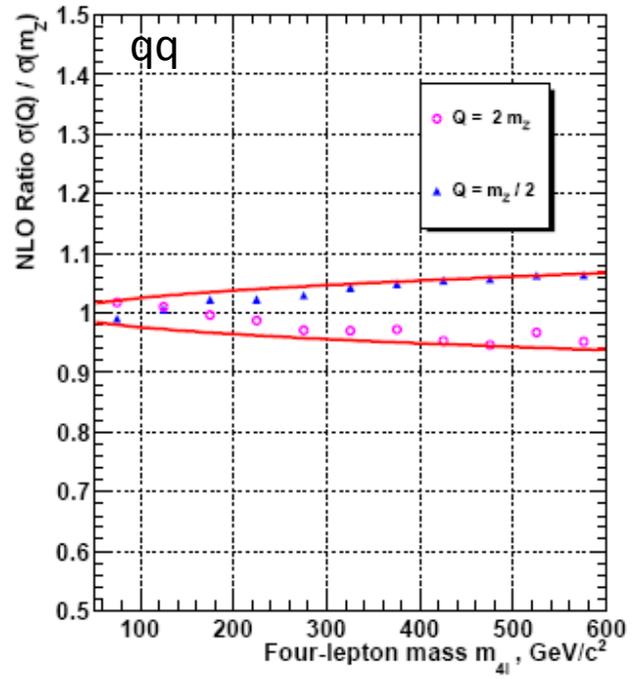
$qq \rightarrow ZZ$ NLO + $gg \rightarrow ZZ$



ZZ: theoretical uncertainties



PDF+ α_s



scale

WW: theoretical uncertainties

- WW taken from MC for large m_H

→ gg+qq NLO available (MCFM)

PDF+ α_s and scale uncertainty dominates

- in jet bins using uncert on $\sigma(>=N)$ + modeling: MC@NLO vs ALPGEN

$\Delta\sigma_{\geq 0}$ (%)	$\Delta\sigma_{\geq 1}$ (%)	$\Delta\sigma_{\geq 2}$ (%)	$\Delta\sigma_{\geq 3}$ (%)
3	6	42	100

- WW from control region for $m_H < 200$ GeV ($m_{ll}, \Delta\phi_{ll}$)

	scale	pdf CTEQ 6.6 error set	pdf central (CTEQ6.6, MSTW2008, NNPDF2.1)	Modelisation
α_{WW}^{0j}	2.5%	2.6%	2.7 %	3.5%
α_{WW}^{1j}	4%	2.5%	1.4 %	3.5 %

WW → lνlν measurement

❑ Complex analysis (no mass peak → counting experiment, many backgrounds)

❑ Main systematics:

▪ background estimate:

Background	Cuts
W+jets	tight lepton quality
top	(b-)jet veto
Drell-Yan	Z mass veto missing E_T
WZ, ZZ, Wγ	2 leptons → estimated from MC

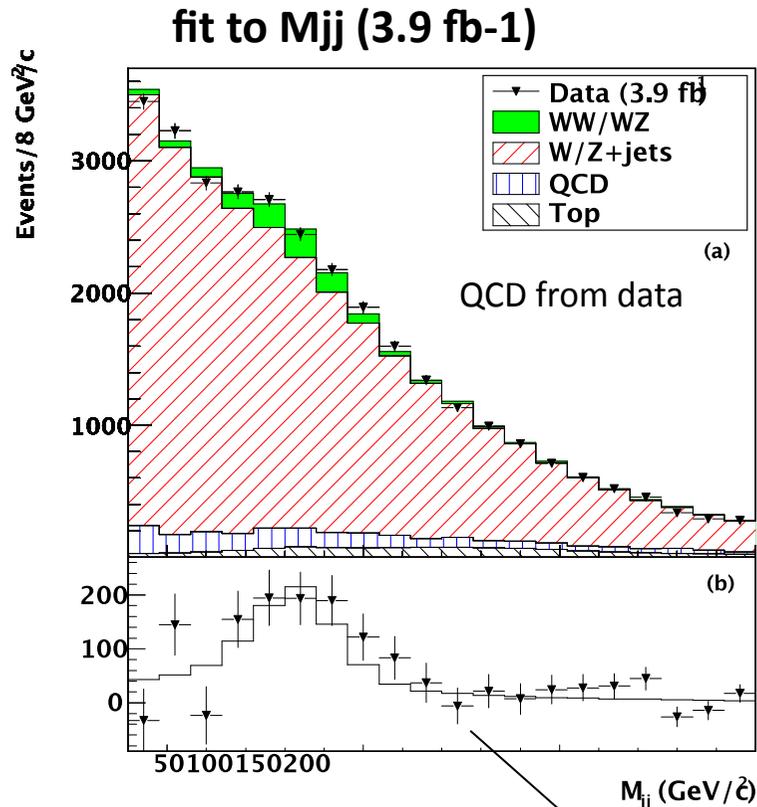
} from data
estimated

Sample	Yield
$qq \rightarrow W^+W^-$	349.7 ± 30.3
$gg \rightarrow W^+W^-$	17.2 ± 1.6
W + jets	106.9 ± 38.9
$t\bar{t} + tW$	63.8 ± 15.9
$Z/\gamma^* \rightarrow \ell\ell + WZ + ZZ$	12.2 ± 5.3
$Z/\gamma^* \rightarrow \tau\tau$	1.6 ± 0.4
WZ/ZZ not in $Z/\gamma^* \rightarrow \ell\ell$	8.5 ± 0.9
W + γ	8.7 ± 1.7
signal + background	568.6 ± 52.2
Data	626

▪ signal acceptance: jet-veto efficiency
leptonic efficiency
missing E_T uncertainty
theoretical (gg box, PDF)

WW/WZ → lν2j at CDF

- ❑ **First observation: 5.4σ** (first evidence at D0 with 4.4σ in 2008)
- ❑ Much **larger backgrounds**, no **resolution** to distinguish W/Z → jj



+ matrix element method: (2.7 fb⁻¹)

- discriminant exploiting full kinematic information, based on calculations of differential xsec of signal and background
- data-MC validation of input kinematic variables
- fit to shape of discriminant

$$\sigma(WW + WZ) = 16.0 \pm 3.3(\text{stat} + \text{syst}) \text{ pb} \quad (\text{NLO expected} \quad 16.1 \pm 0.9 \text{ pb})$$

WZ/ZZ → lν/νν + 2b at Tevatron

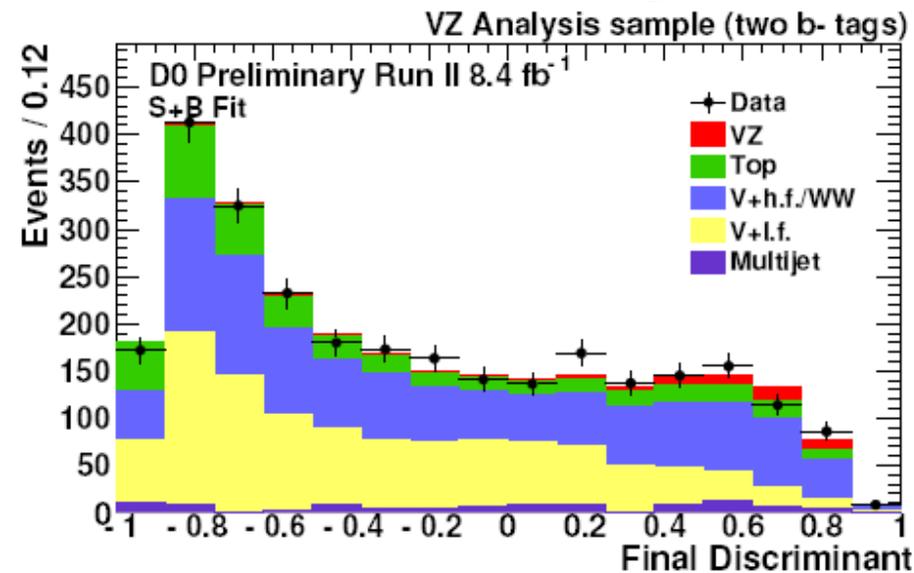
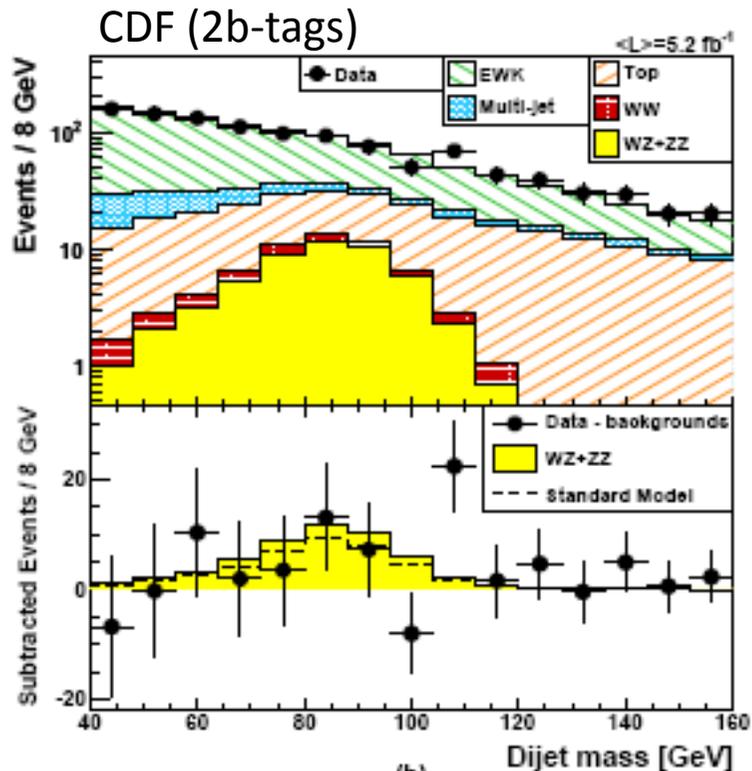
❑ Crucial for Higgs search (ZH → ννbb). Very complex analysis!

• WZ → lνbb + ZZ → ννbb → b-tag jets + missing E_T (+ topological cuts)

• No leptons! → huge background: multijets QCD, V+jets (from data)

• very sophisticated techniques: b-tag probability with Boosted Decision Tree or Neural Network which exploits much info and different variables

different channels combined (0,1,2 b-tag)



D0: $\sigma(WZ, ZZ) = 6.9 \pm 1.3(stat) \pm 1.8(syst) pb$

CDF: $\sigma(WZ, ZZ) = 5.8^{+3.6}_{-3.0} pb$ (expected 4.6 pb)

(expected 5.1 pb)

ZZ at CDF

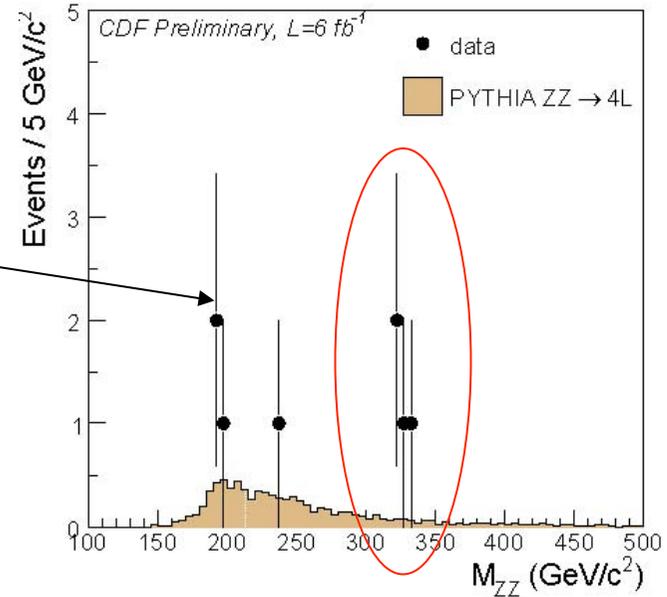
□ **ZZ→4l** (6 fb⁻¹) excess of events at high M_{ZZ} (eg Randall Sundrum Graviton)

- But xsec still compatible with SM

$$\sigma(p\bar{p} \rightarrow ZZ) = (2.8^{+1.2}_{-0.9} (\text{stat.}) \pm 0.3 (\text{syst.})) \text{ pb}$$

$$(1.4 \pm 0.1) \text{ pb}$$

- No excess in other final states



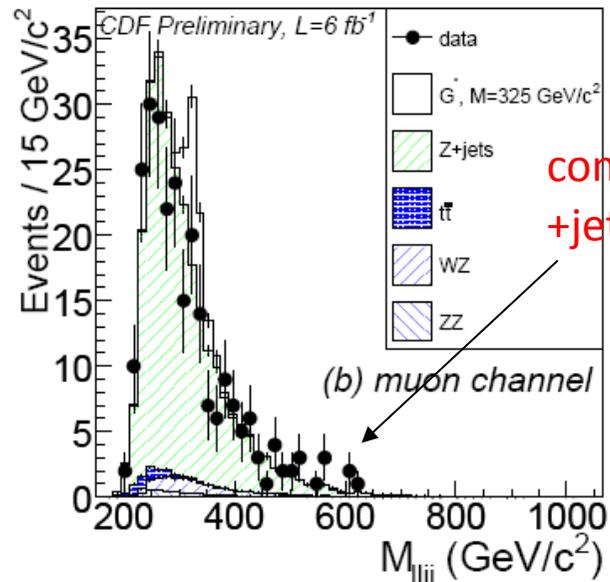
□ **ZZ→2l2ν** (5.9 fb⁻¹)

Shape analysis with fit to neural network

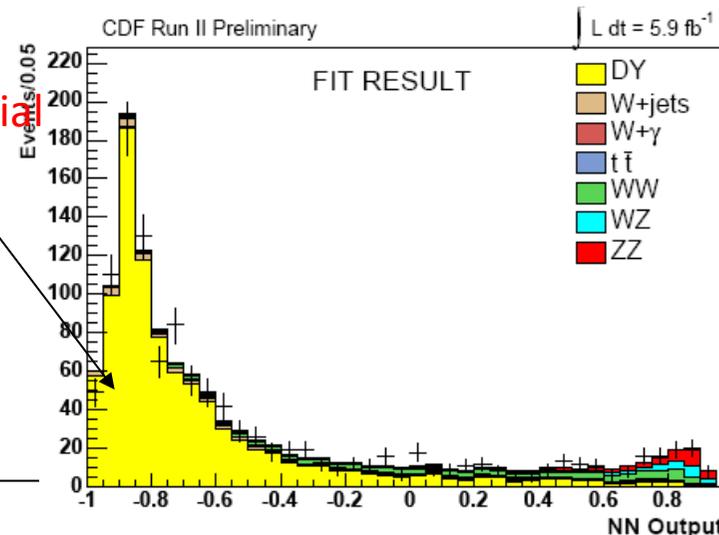
$$\sigma(p\bar{p} \rightarrow ZZ) = 1.45^{+0.45}_{-0.42} (\text{stat.})^{+0.41}_{-0.30} (\text{syst.}) \text{ pb}$$

$$\sigma(ZZ) = 1.21^{+0.05}_{-0.04} (\text{scale})^{+0.04}_{-0.03} (\text{PDF}) \text{ pb.}$$

□ **ZZ→2l2j**

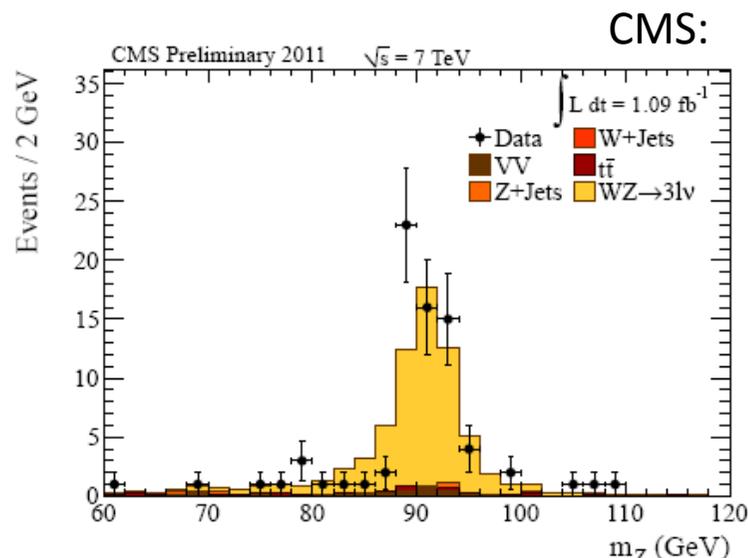


control of Z +jets is crucial



WZ → 3l + ν at LHC

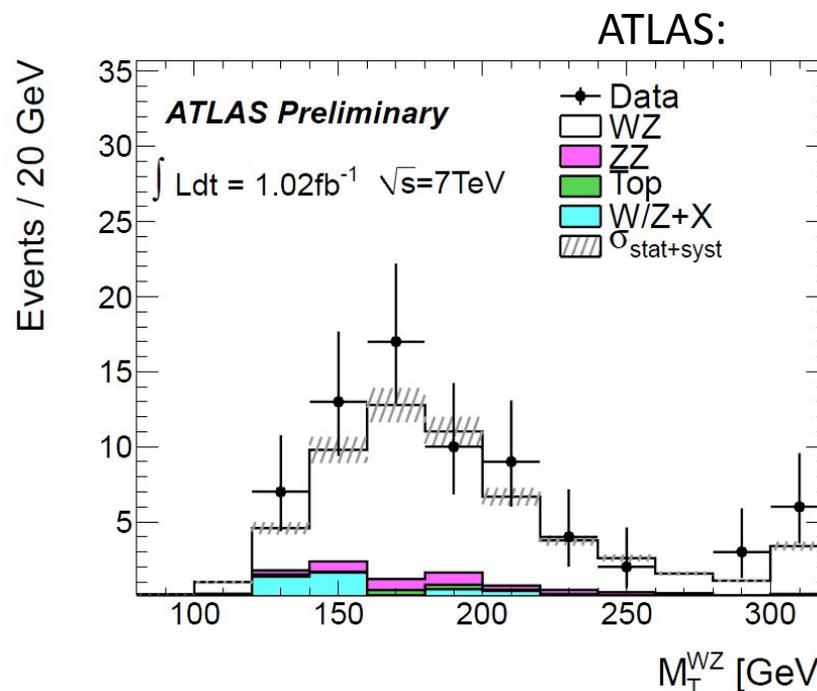
☐ Very clean, low background



Sample	3e0μ	2e1μ	1e2μ	0e3μ
Background	2.95	2.14	1.31	1.72
WZ → 3lν	14.47	17.49	13.95	18.56
AllMC	17.42	19.62	15.26	20.28
Data	22	20	13	20

$$\sigma(WZ) = 17.0 \pm 2.4(stat) \pm 1.1(syst) \pm 1.0(lumi) pb$$

$$(NLO \text{ expected } 18.75_{-0.8}^{+1.1} pb)$$



$$\sigma(WZ) = 21.1_{-2.8}^{+3.1}(stat) \pm 1.2(syst)_{-0.8}^{+0.9}(lumi) pb$$

$$(NLO \text{ expected } 17.2_{-0.8}^{+1.2} pb)$$

Recent $WZ \rightarrow 3l + \nu$ at Tevatron

□ Possible **only at hadron colliders**
(charged final state)

▪ **D0 with 4.1 fb⁻¹**

statistics 2 times smaller than LHC,
while xsec is 6 times smaller

$$\sigma(WZ) = 3.90^{+1.06}_{-0.90} pb$$

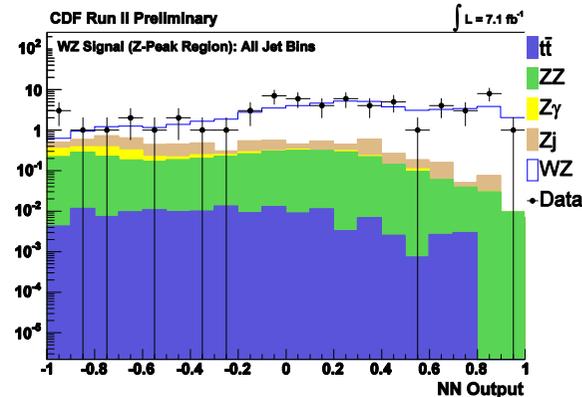
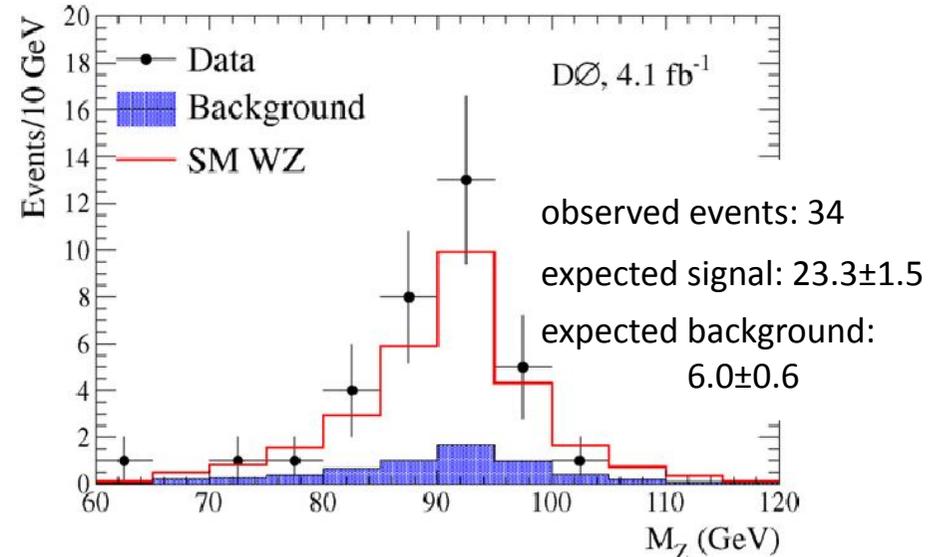
(NLO expected 3.25 ± 0.1) $9 pb$

▪ **CDF with 7.1 fb⁻¹**

Shape analysis with fit to neural
network

$$\sigma(WZ) = 3.9^{+0.6}_{-0.5} (stat)^{+0.6}_{-0.4} (syst) pb$$

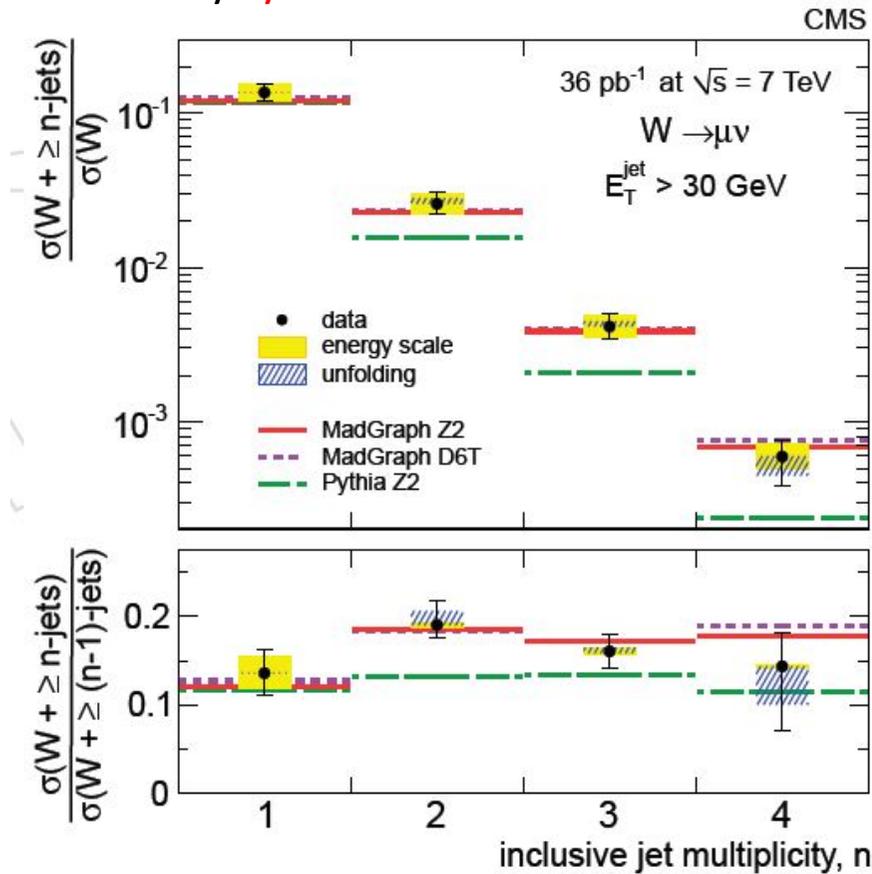
(NLO expected 3.46 ± 0.2) pb



V+jets: ratios

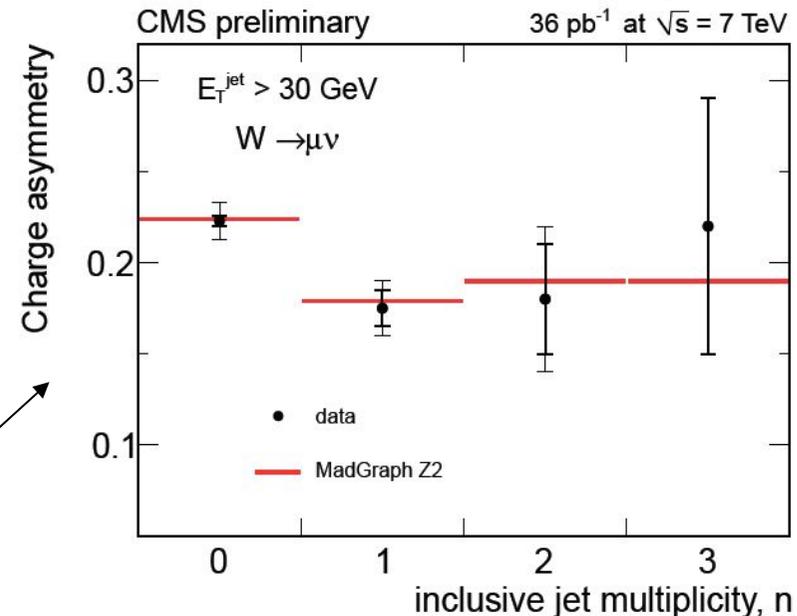
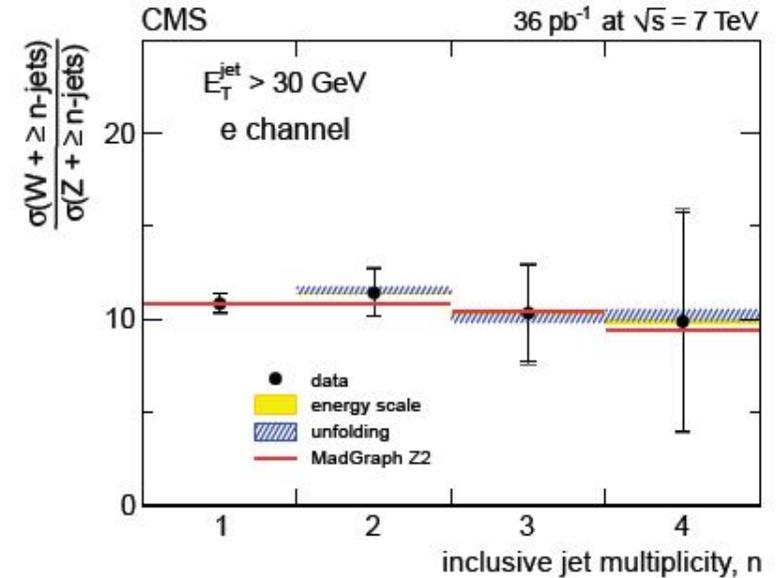
CMS PAS EWK-10-011

- Many systematics cancel out



- Very well known theoretically → any deviation is hint of NP

- Use high statistics W sample to predict Z+jets background for NP



The input: signal description

Common effort CMS-ATLAS-theoreticians:

Higgs Xsec working group

(2 yellow reports published: [arXiv:1101.0593 \[hep-ph\]](https://arxiv.org/abs/1101.0593)
[arXiv:1201.3084 \[hep-ph\]](https://arxiv.org/abs/1201.3084))

☐ Higgs p_T reweighted at NNLO+NNLL
 (Powheg+HqT+HNNLO)

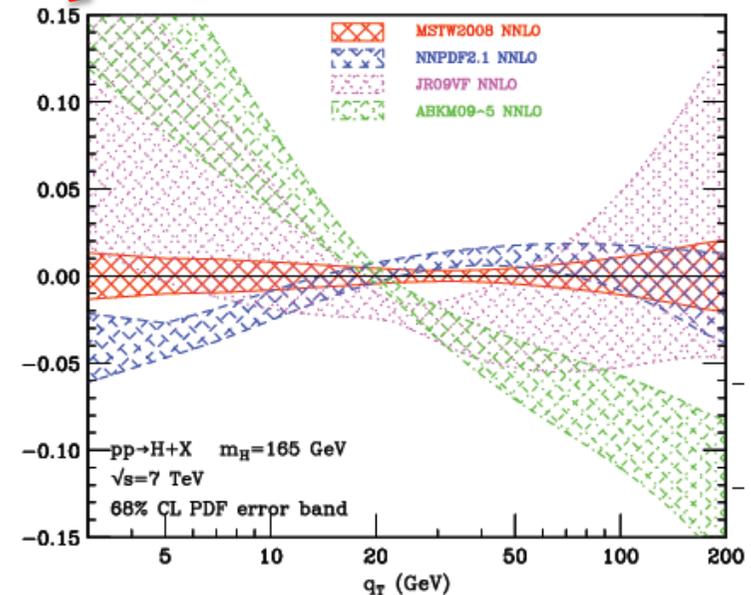
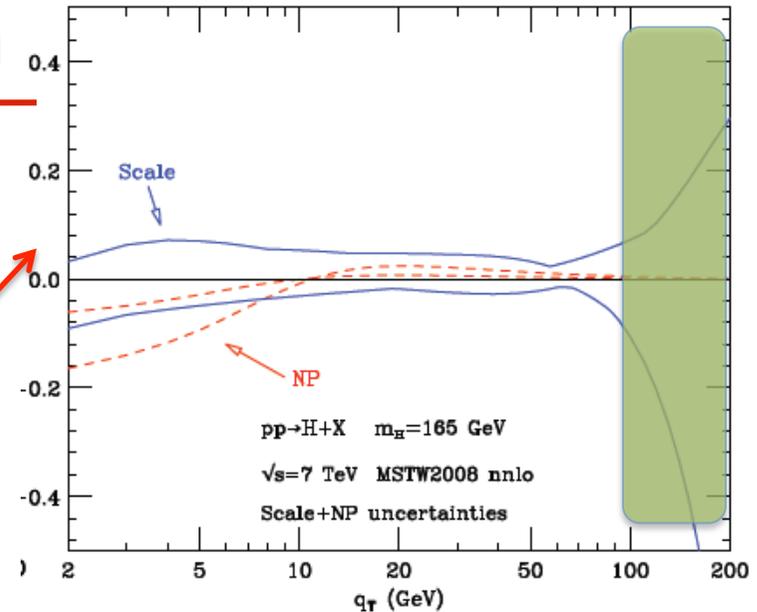
uncertainties: • factor/renorm scale
 • non perturb. effects, PDF

☐ Mass shape uncertainty at high mass:

- xsec for **on-shell** Higgs production and decay in **zero width** approx
- acceptance from MC with **ad-hoc BW distribution**

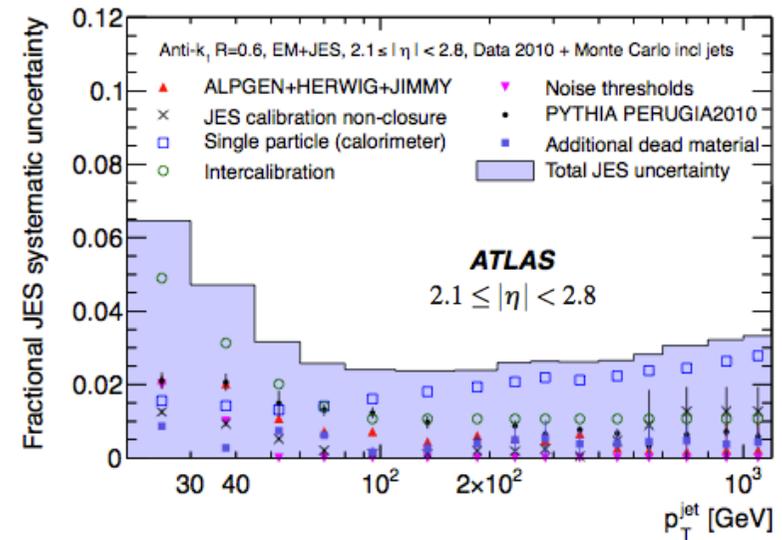
10-30% uncertainty on xsec for m_H 400–600 GeV

Now shape from proper (QFT-consistent)
 Higgs propagator available!

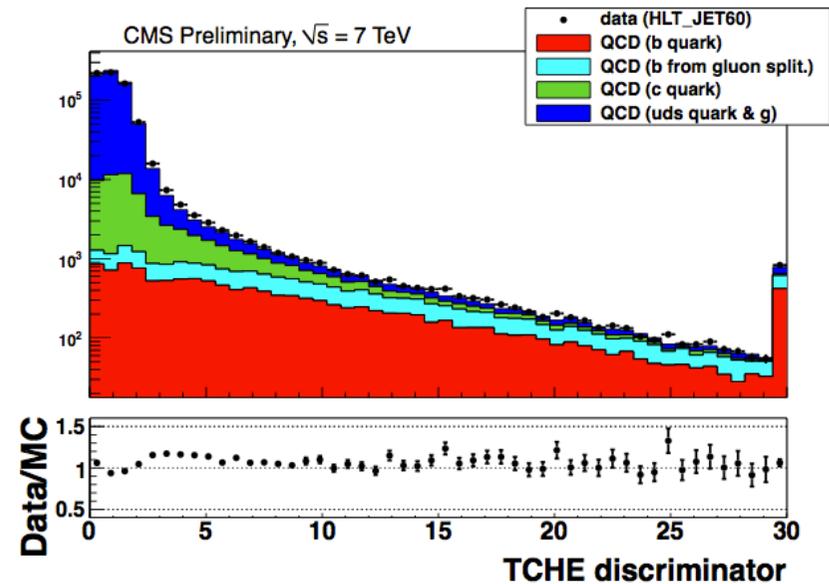
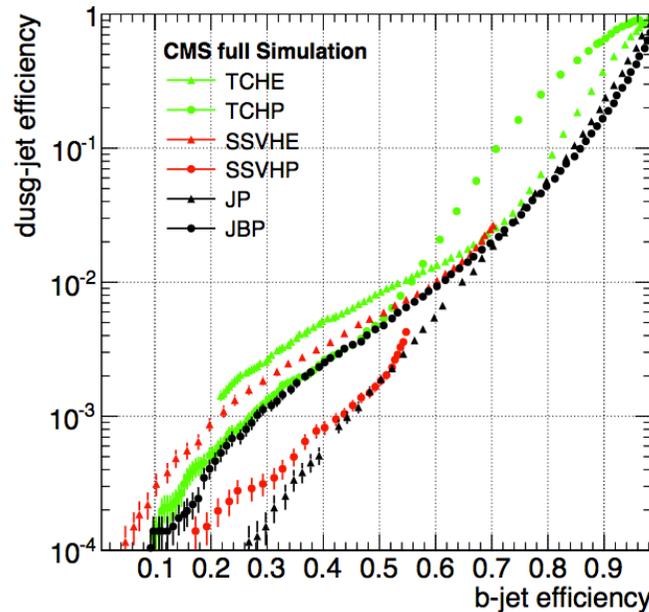


The input: object reconstruction

- **Leptons (e,μ) from V**: $p_T > 20$ GeV, isolated, if from Z then $M(\ell\ell) \sim 91$ GeV
 -> optimal performances, see previous lecturer
- **Jets**: ATLAS: clusters of calorimeter deposits, $p_T > 25$ GeV
 CMS: particle-flow (full particle identification with tracker+calo), $p_T > 30$ GeV



- **B-tagging**



- **MET** : calorimeter deposits (in CMS also part. flow)

H → ZZ → llvv

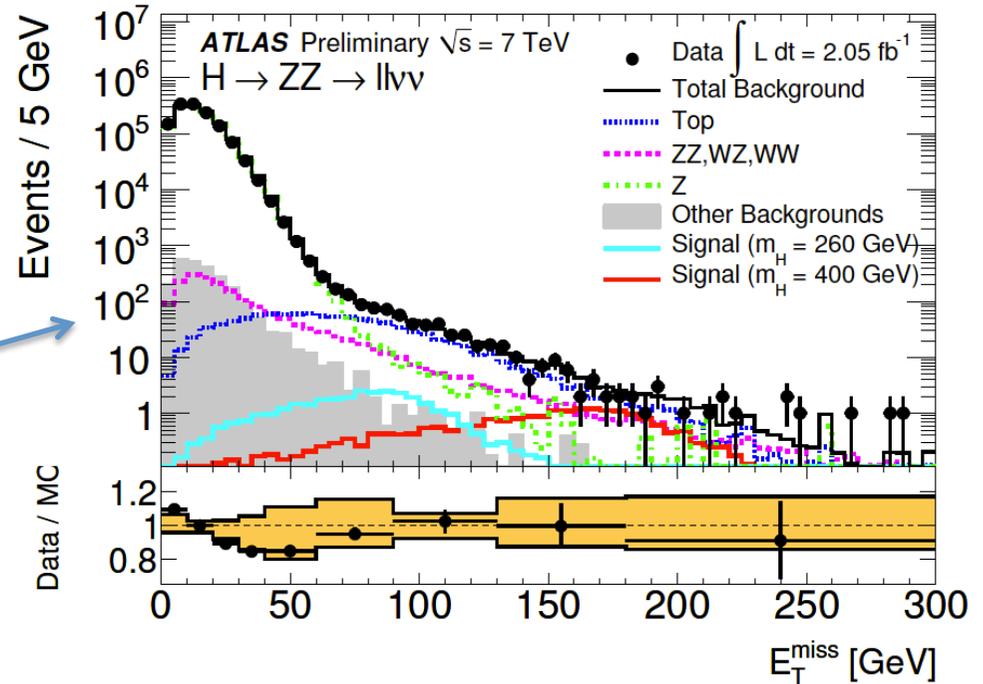
□ 2 leptons from Z + MET.

▪ **Backgrounds:**

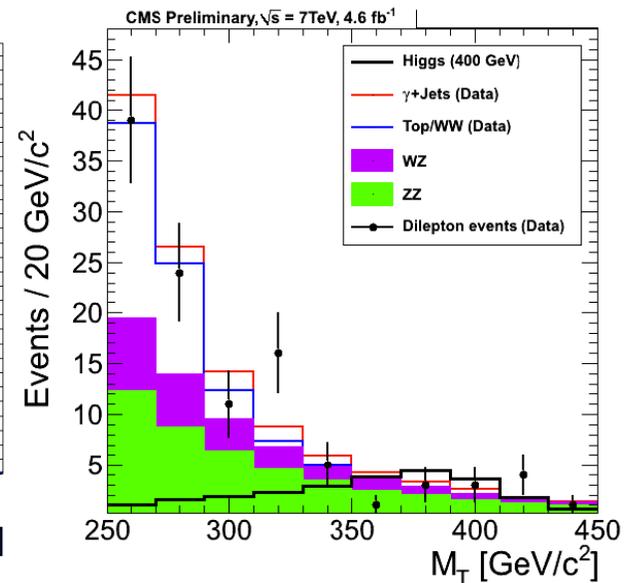
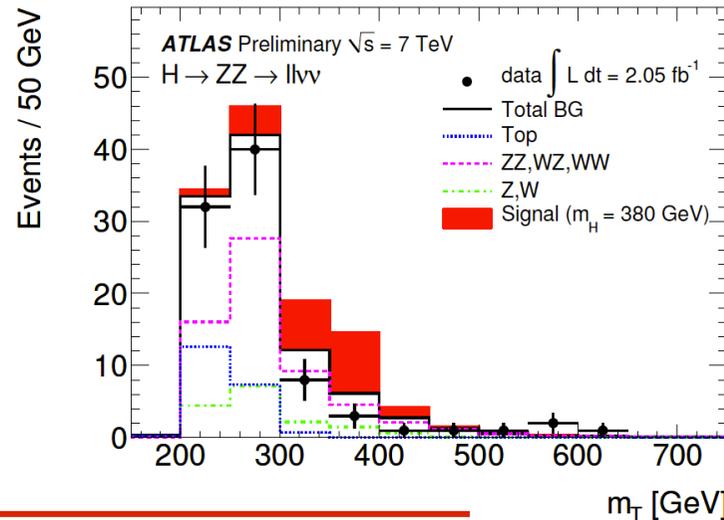
- **Z+jets** removed asking for high, “real” MET
high $\Delta\phi(\text{jet-MET}, \text{Z-MET})$
- **ttbar** removed vetoing b-tag jets
- VV mainly irreducible
- W+jets with lepton fake or from HF

▪ **Final cuts:**

- high MET
- high $\Delta\phi(\text{ll})$ or $p_T \text{Z} \rightarrow$ for high m_H
- **UL from fit to m_T**
broad distribution but pretty good S/B



$$m_T^2 \equiv \left[\sqrt{m_Z^2 + |\vec{p}_T^{\ell\ell}|^2} + \sqrt{m_Z^2 + |\vec{p}_T^{\text{miss}}|^2} \right]^2 - \left[\vec{p}_T^{\ell\ell} + \vec{p}_T^{\text{miss}} \right]^2$$



H->ZZ->llvv: Z+jets control

- **Z+jets is $10^5 \times$ signal and steeply falling** with MET and $p_T Z$
 - high MET tail dominated by **fake MET** (difficult to simulate)
 - cut on **Z boost depends on Z+jets modeling**

ATLAS

MET shape comparison
pythia-alpgen (5-20%)

JES systematics on MET

normalization data-MC
 comparison $m(\text{ll})$ (2.5%)

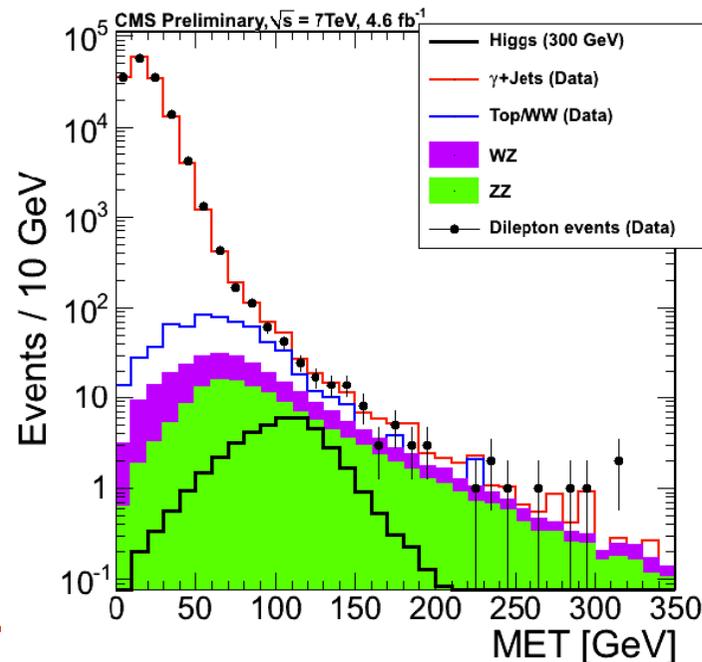


7% uncertainty

CMS: fully data-driven

γ +jets reweighted vs N jets, boson p_T , N vertices

boson mass by sampling Z line-shape from
 fit to data



Main systematics:
 contamination from
 γ + real MET



Z+jets left float btw
 0 and γ +jet prediction
(-100%,+0% uncertainty)

H->ZZ->llvv: non resonant backgrounds

□ ATLAS:

- **ttbar from MC**
 - xcheck in control regions (mll sidebands, e+mu)
 - systematics btag eff and mistag rate

-> **9% uncertainty**

- W+jets normalization from control region (same-sign ll + MET)

-> **100% uncertainty**

□ CMS:

- **non resonant from data:**
e+mu and m(ll) sidebands

$$N_{\mu\mu} = \frac{N_{\mu\mu}^{SB}}{N_{e\mu}^{SB}} \times N_{e\mu}$$

systematics from statistics in e+mu data

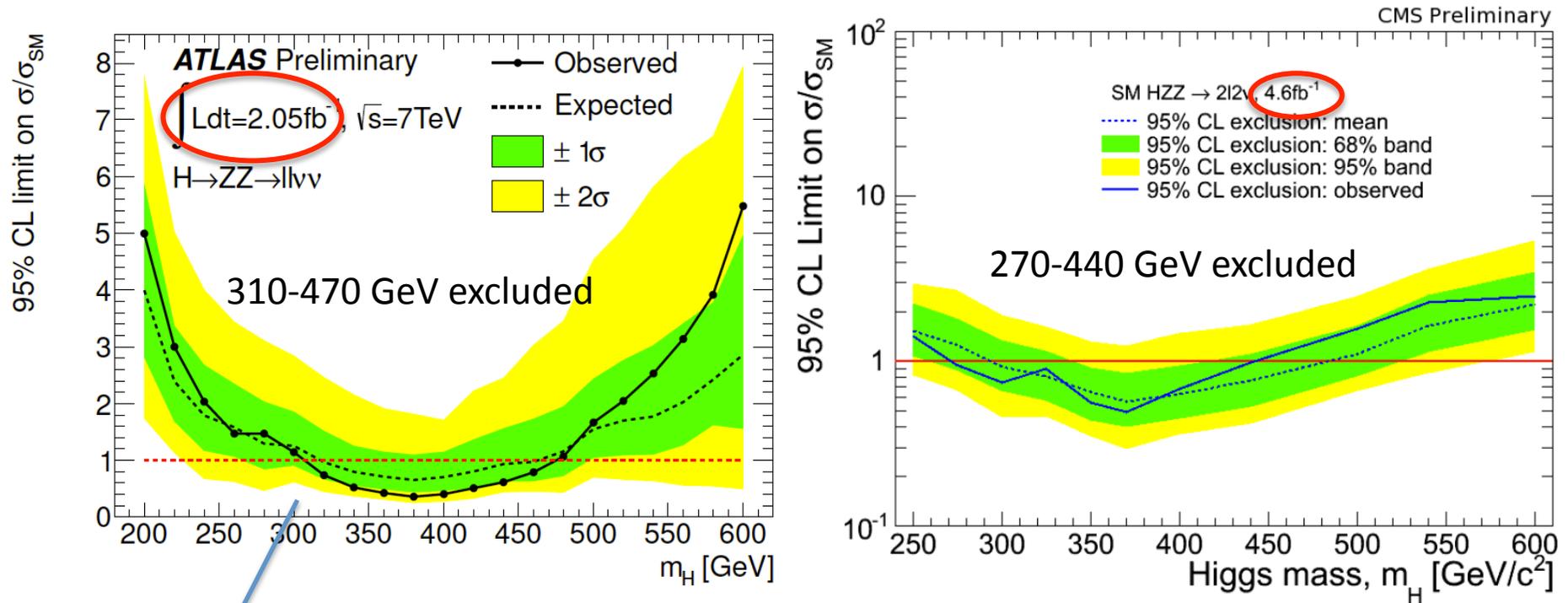
-> **15-100% uncertainty**

- VV from MC ->systematics on norm from PDF, scale, missing VZ*
-> **5-10% uncertainty**

CMS 4.6 fb⁻¹

m_H (GeV/c ²)	ZZ	WZ	Top/WW/ W+jets/Z->ττ	Z+Jets	Total	Signal	Data
400	12 ± 0.13 ± 0.87	4.6 ± 0.14 ± 0.39	0	2.7 ± 2.7	19 ± 0.19 ± 2.9	17 ± 1.6	18

H->ZZ->llvv: results



ATLAS includes $H \rightarrow WW \rightarrow l\nu l\nu$ (13%-77%) and $H \rightarrow ZZ \rightarrow 4l$, $H \rightarrow ZZ \rightarrow 2l2q$
 (exclusive cuts wrt to dedicated analysis to avoid double counting in combination)

H → ZZ → ljj

□ 2l from Z + 2jets from Z

(to improve m_H resolution:
kinematic fit M(Z → jj) in CMS,
m_Z/m_{jj} rescaling in ATLAS)

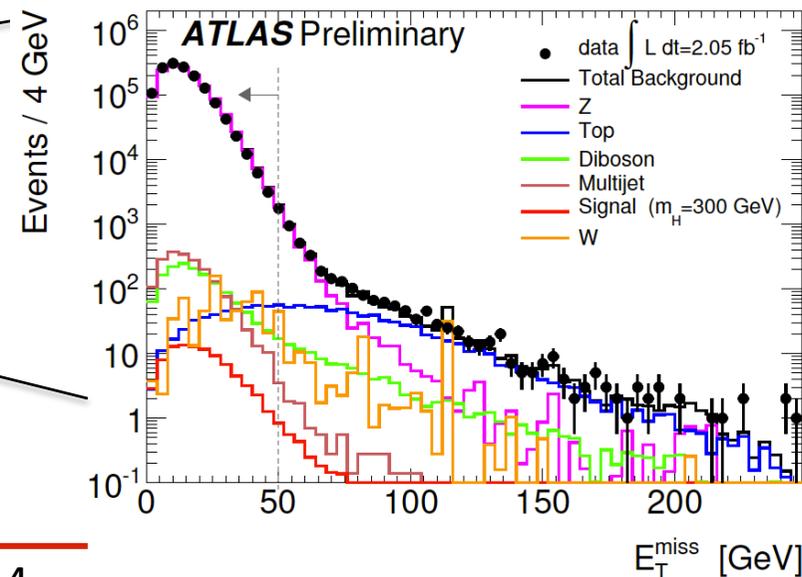
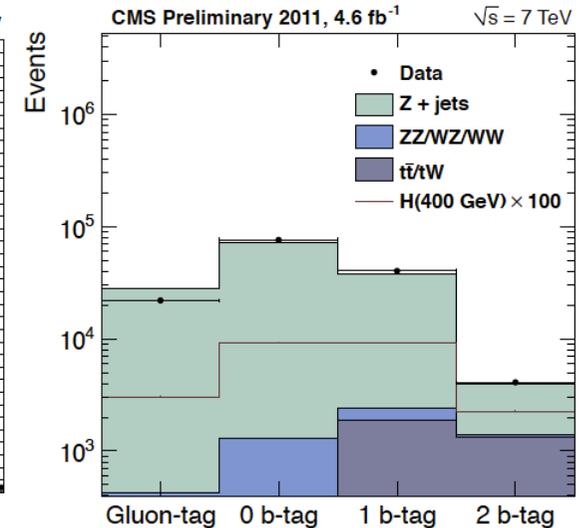
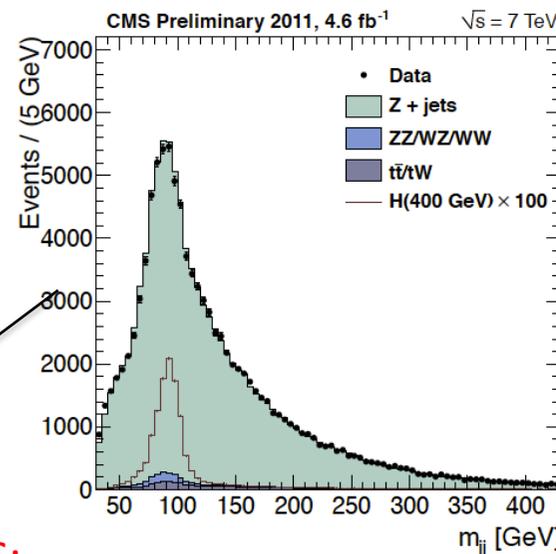
□ Main background Z+jets

• still peaking ~M_Z because of
kinematic p_T cuts

• no b-jets → b categories (CMS:
0,1,2 bjet; ATLAS: 2 bjets or less)

• (CMS remove jet from gluons
with qg-likelihood discriminant)

□ In 2 btag ttbar relevant → MET cut
(MET significance in CMS)



H->ZZ->lljj: final cuts

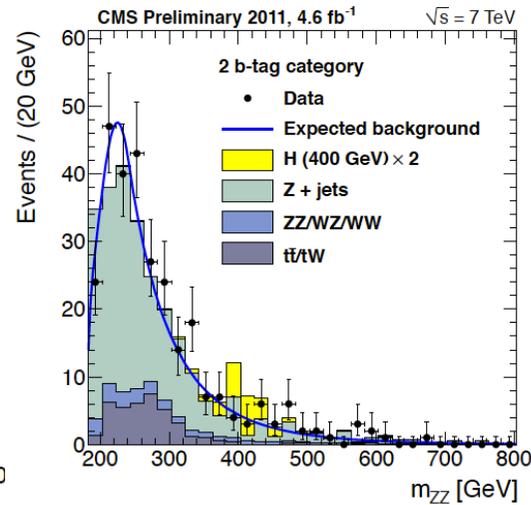
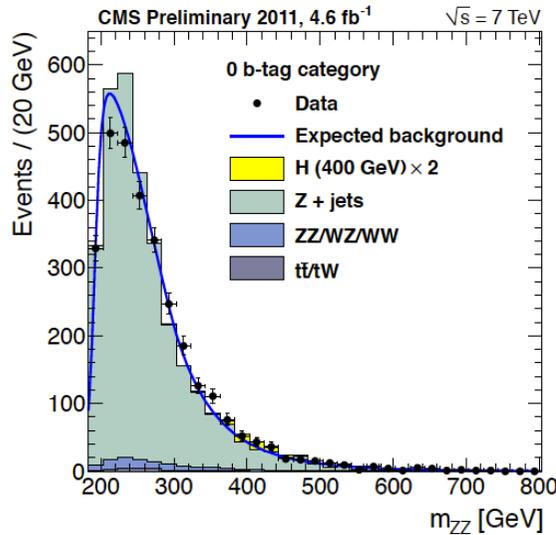
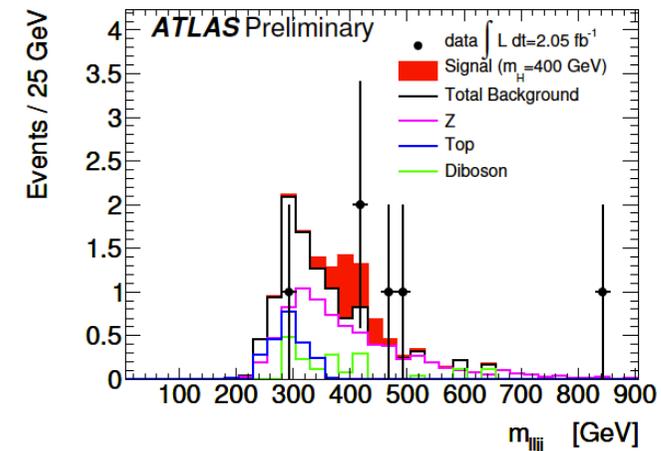
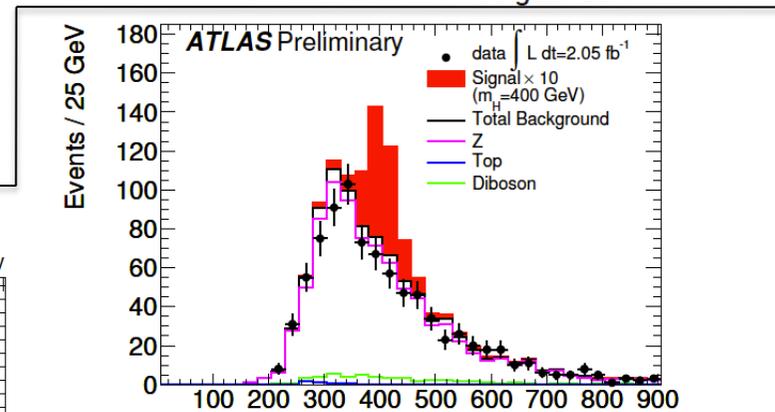
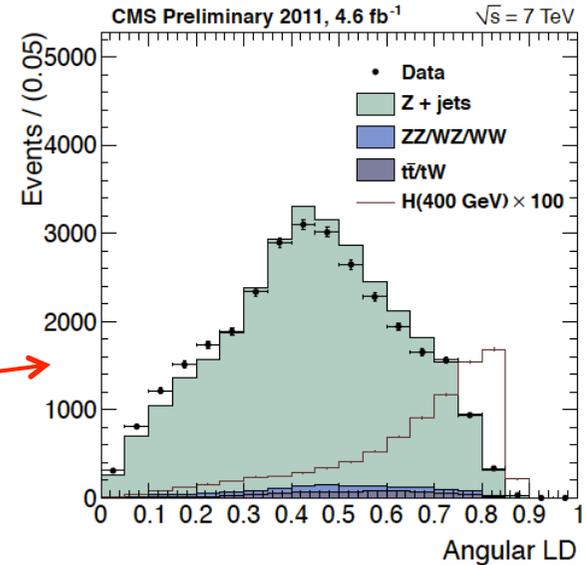
□ **ATLAS**: higher p_T jets + small $\Delta\phi(jj)$, $\Delta\phi(ll)$ -> **high Z p_T** for high mass Higgs

□ **CMS**: exploit **full angular info**
H(spin0) -> ZZ -> 4f :

likelihood discriminant built with angles distributions

->no $m(jj)$ bias -> background from sidebands

□ UL from **fit to $m(lljj)$**



H->ZZ->lljj: background control

□ ATLAS:

▪ Z+jets

- **shape from MC** -> uncertainty from Pythia-AlpGen comparison
- **normalization from data M(jj) sidebands**
-> uncertainty from data statistics: 1-12%
depend. on mH and untagged (16%)

▪ ttbar from MC

- xcheck in MET>50 GeV, M(ll) sidebands
-> uncertainty from data statistics +
theoretical uncertainty (~15-20%)

□ CMS:

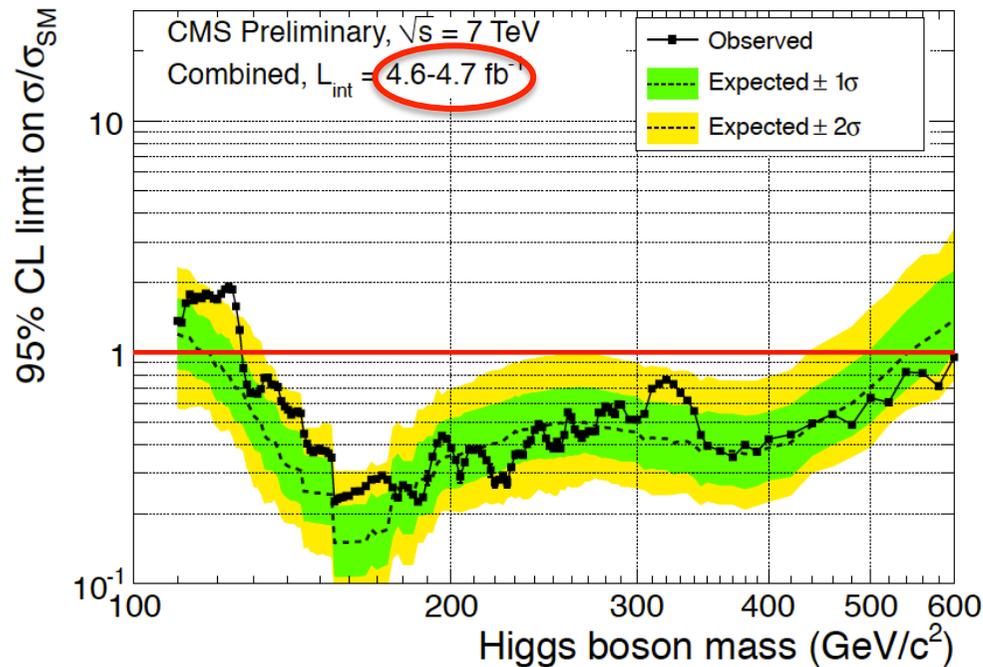
▪ All background at once, data-driven (highly dominated by Z+jets)

- **shape and normalization from data M(jj) sidebands**
- rescaled SB/SR from MC

-> uncertainty from statistics

Combined results: high mass

127-600 GeV 95% CL exclusion



131-237 251-468 GeV 95% CL exclusion

