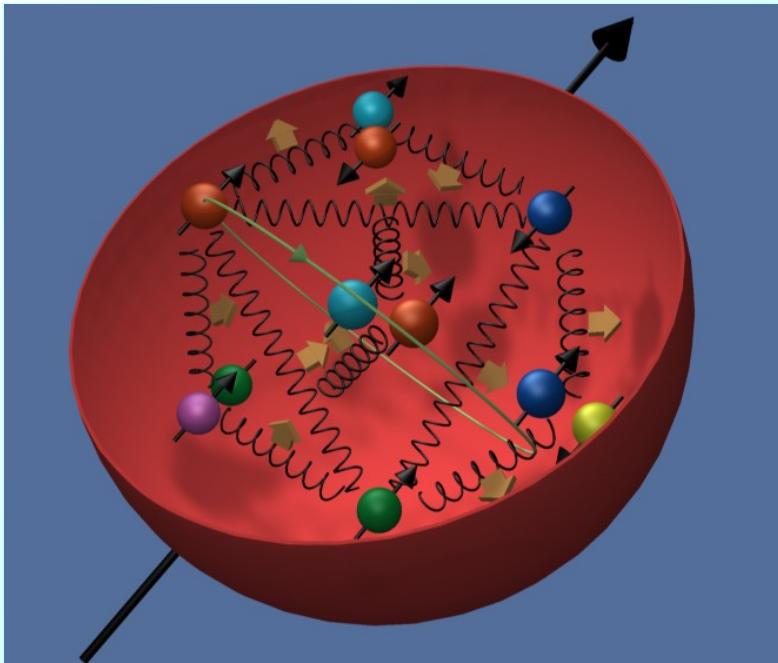


New Results

Klaus Rith

University of Erlangen-Nürnberg



Main HERMES research topics:

- Origin of nucleon spin
- Details of nucleon structure

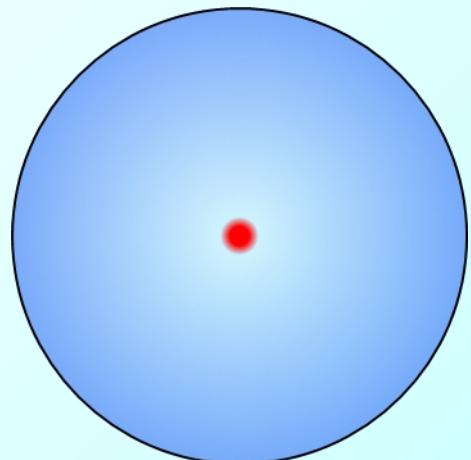
Quark Properties:

- fractional charge
- spin-1/2
- longitudinal momentum xP
- intrinsic transverse momentum p_T
- spatial position r_\perp
- orbital angular momentum L

Atom

(non-relativistic electrons in Coulomb potential)

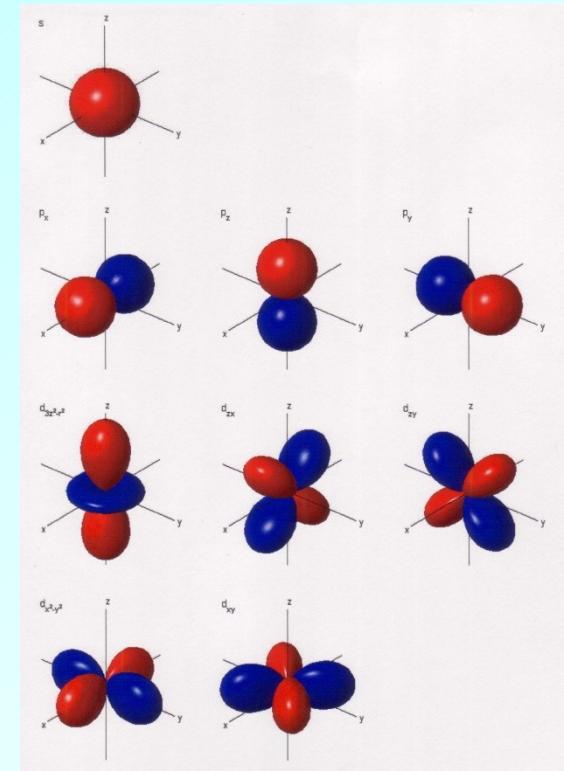
Rutherford



Add angular momentum



Bohr, Schrödinger, ..

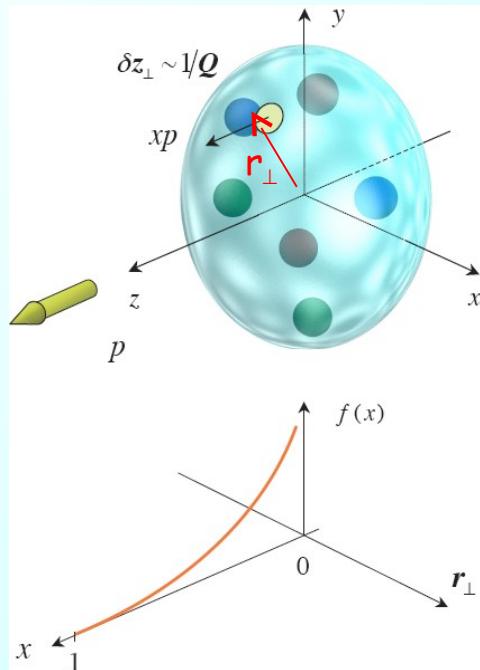


$$\Psi_{n,l,m_l}(r,\theta,\phi)$$

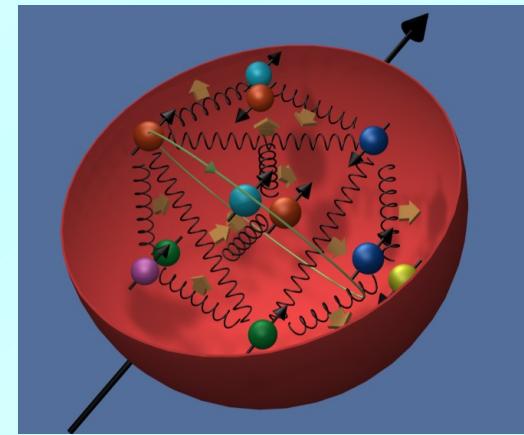
Nucleon

(Relativistic quarks in colour field)

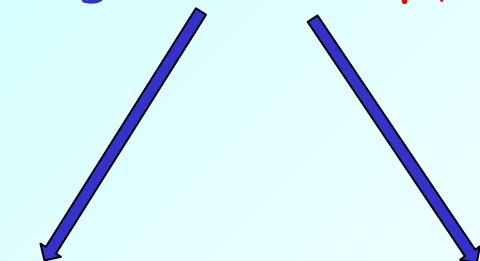
Inclusive DIS



Add angular and
transverse momentum



Wigner DF $W(p_T, r_\perp)$



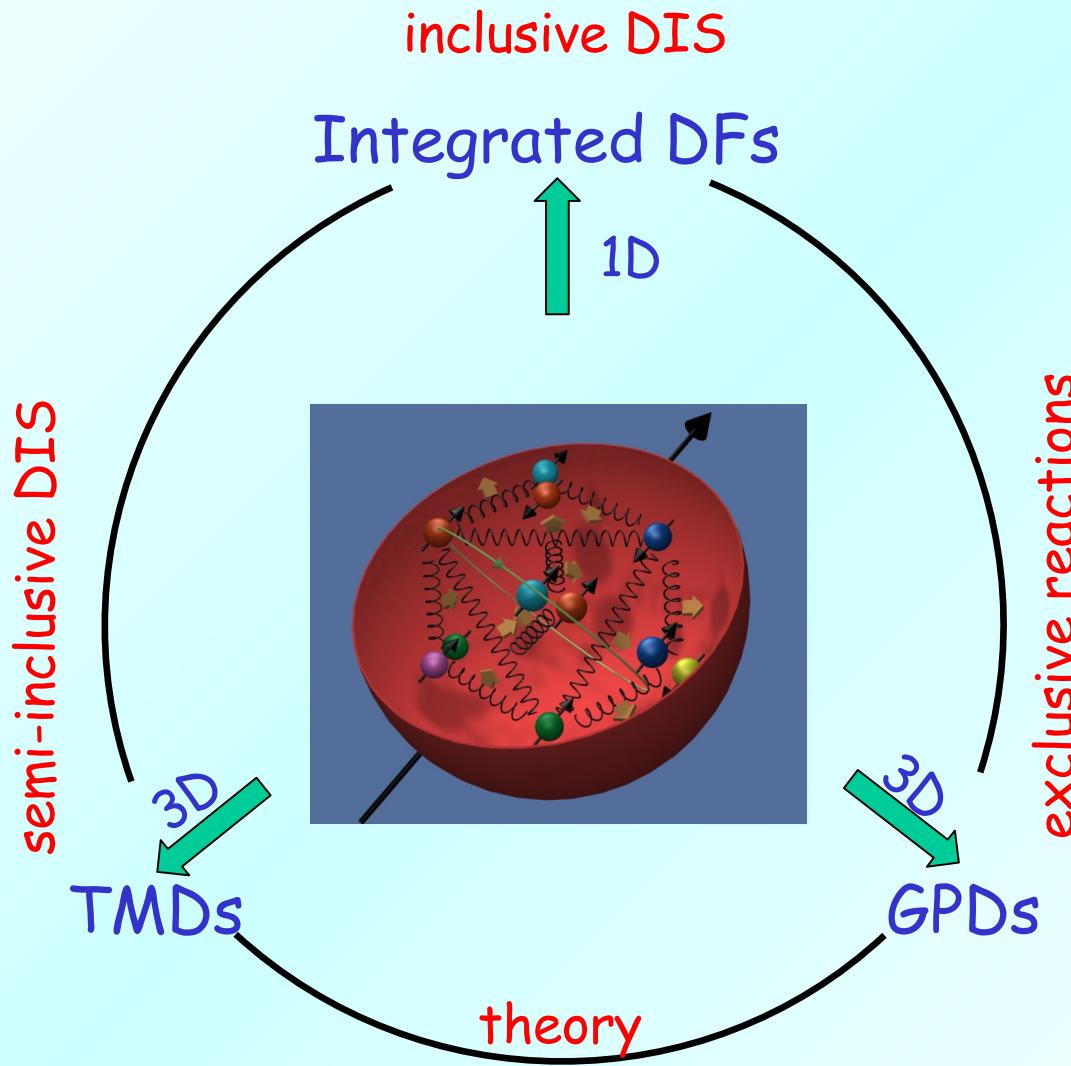
TMDs

GPDs

(p_T -dependence) (r_\perp -dependence)

Number density of quarks
with longitudinal
momentum fraction x

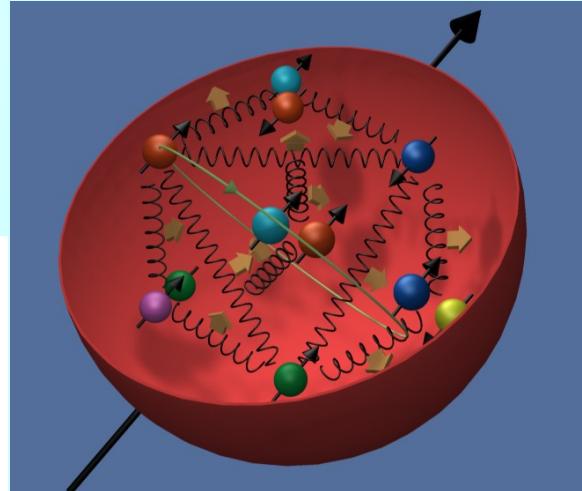
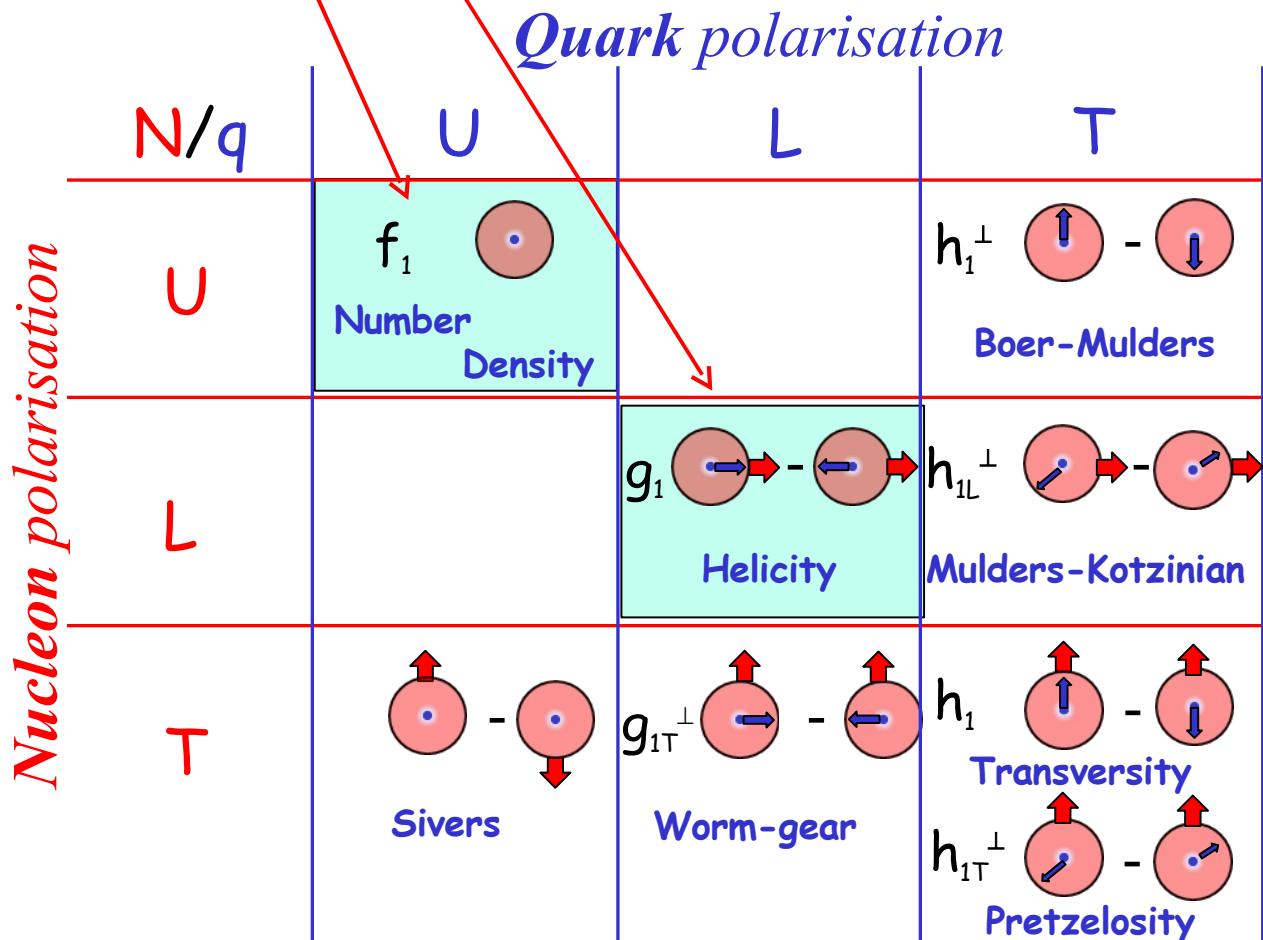
Accessing the nucleon's structure



after G. Schnell

TMDs

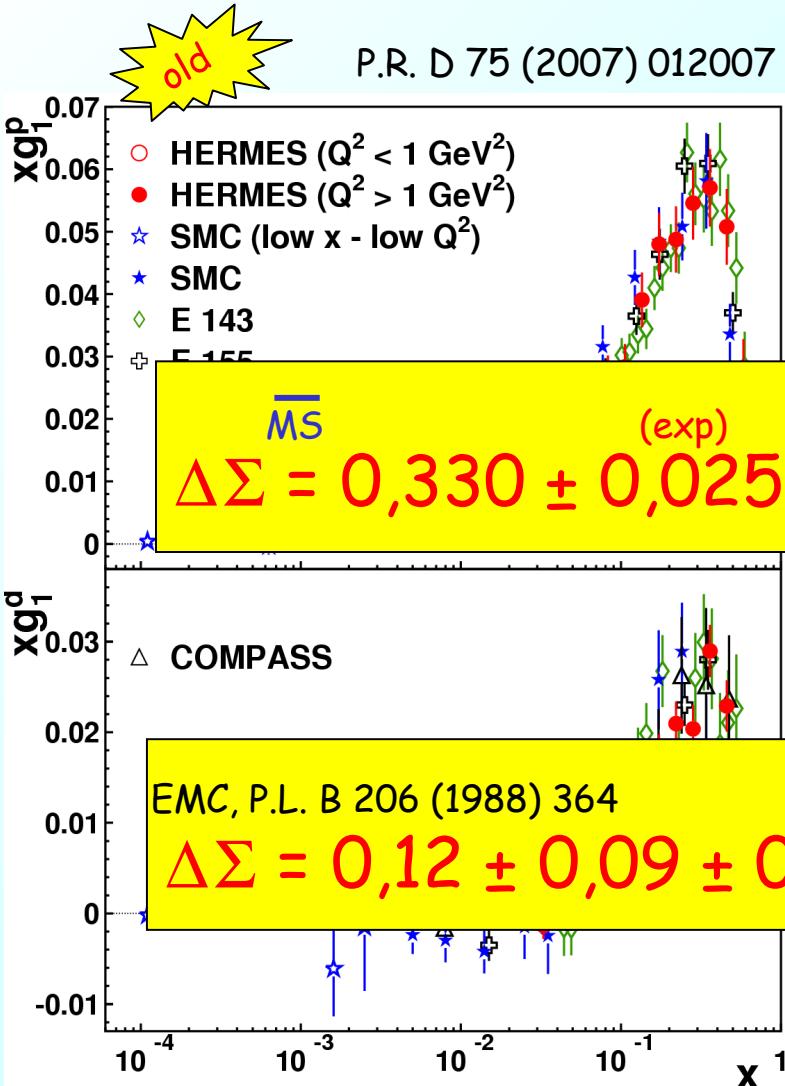
Accessible in inclusive DIS



Nucleon structure described by 8 leading-twist (+ many subleading) quark distributions containing information about quark orbital motion and spin-orbit effects

Inclusive Measurements

Longitudinal double-spin asymmetry: g_1 , $\Delta\Sigma$



$$g_1(x) = \frac{1}{2} \sum_q e_q^2 \Delta q(x)$$

$$\Delta q = \int_0^1 \Delta q(x) dx$$

$$\Delta\Sigma = \sum_q$$

HERMES, JHEP 08 (2010) 130

new

Furthermore: $\Delta g/g = 0,045 \pm 0,034 \pm 0,126$ (high- p_T hadrons)

Unpolarised DIS cross section: F_2

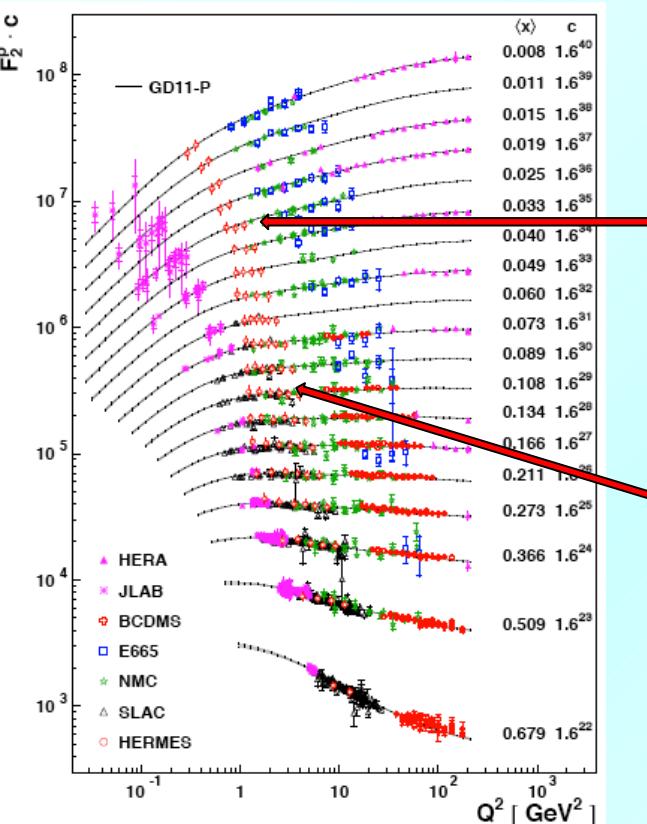
$$\frac{d^2\sigma}{dx dQ^2} = \frac{4\pi\alpha_{em}^2}{Q^4} \frac{F_2(x, Q^2)}{x} \left[1 - y - \frac{Q^2}{4E^2} + \frac{y^2 + Q^2/E^2}{2(1 + R(x, Q^2))} \right]$$

new
JHEP 05 (2011) 126

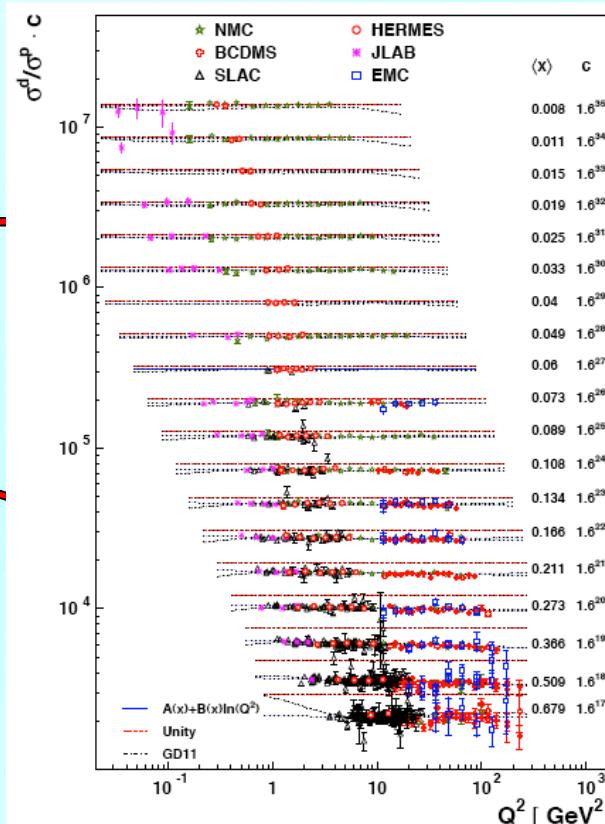
Exploring perturbative to non-perturbative
regime in an unmeasured x - Q^2 region
 $0.006 < x < 0.9$ $0.1 \text{ GeV}^2 < Q^2 < 20 \text{ GeV}^2$

From global fit: HERMES
relative normalisation ~2%
for p and d and ~0.5% for
the ratio

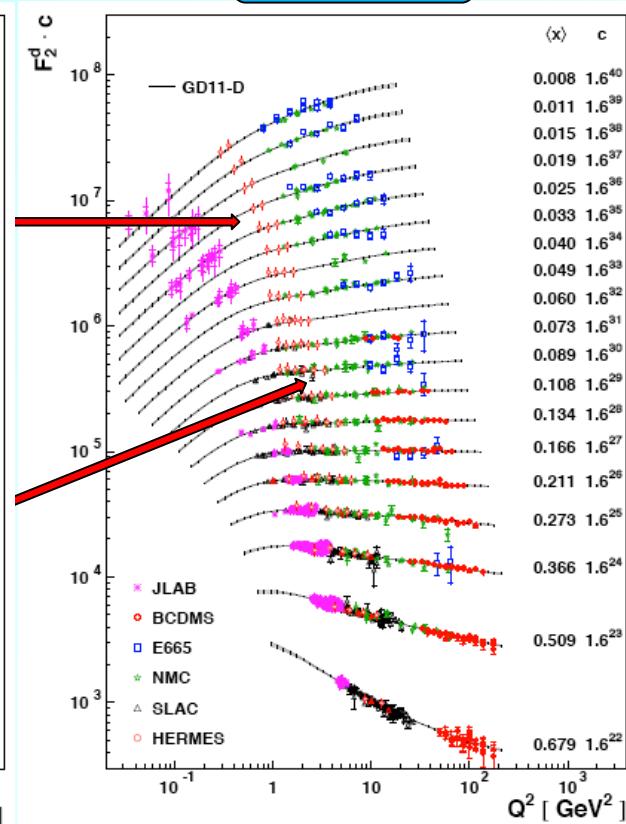
Proton



Ratio $\sigma^d/\sigma^p (F_2^d/F_2^p)$

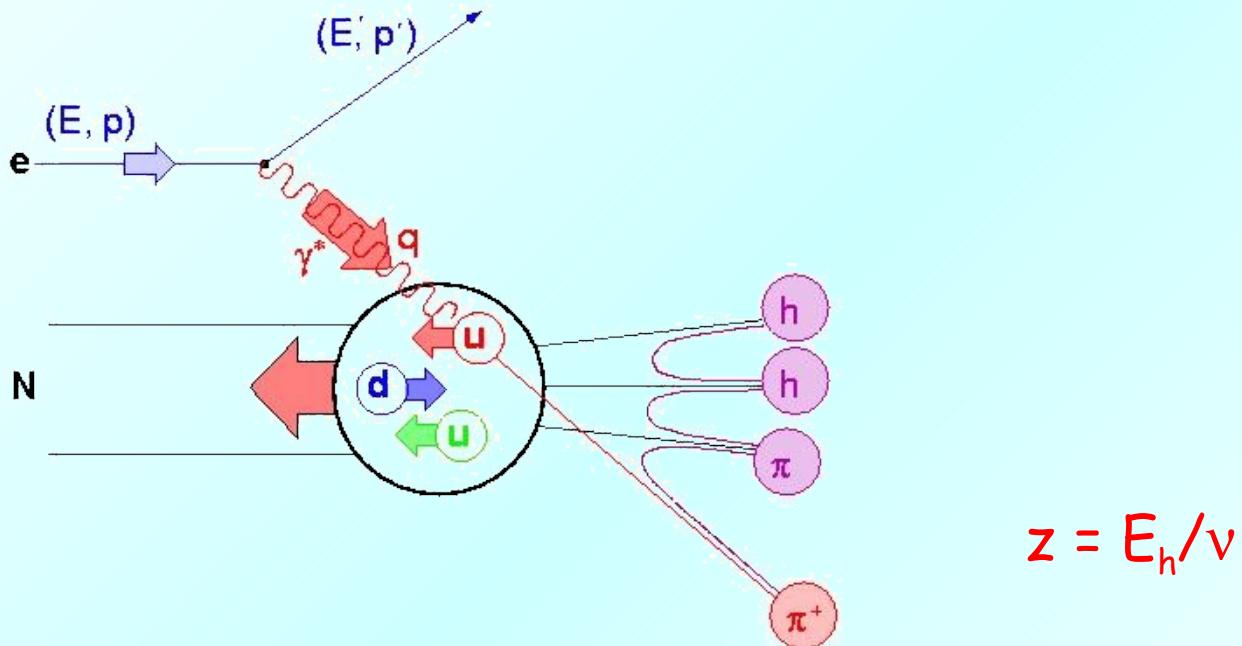


Deuteron



Semi-inclusive Measurements

Semi-inclusive Deep-Inelastic Scattering



$$\text{Factorisation} \rightarrow \sigma^{eN \rightarrow ehX} = \sum_q \sigma^{eq \rightarrow eq} \otimes DF^{N \rightarrow q} \otimes FF^{q \rightarrow h}$$

$DF(x, Q^2)$: Parton Distribution Function - $q(x, Q^2) \equiv f_1^q(x, Q^2)$,
 $\Delta q(x, Q^2) \equiv g_1^q(x, Q^2)$, $\delta q(x, Q^2) \equiv h_1^q(x, Q^2)$, ...

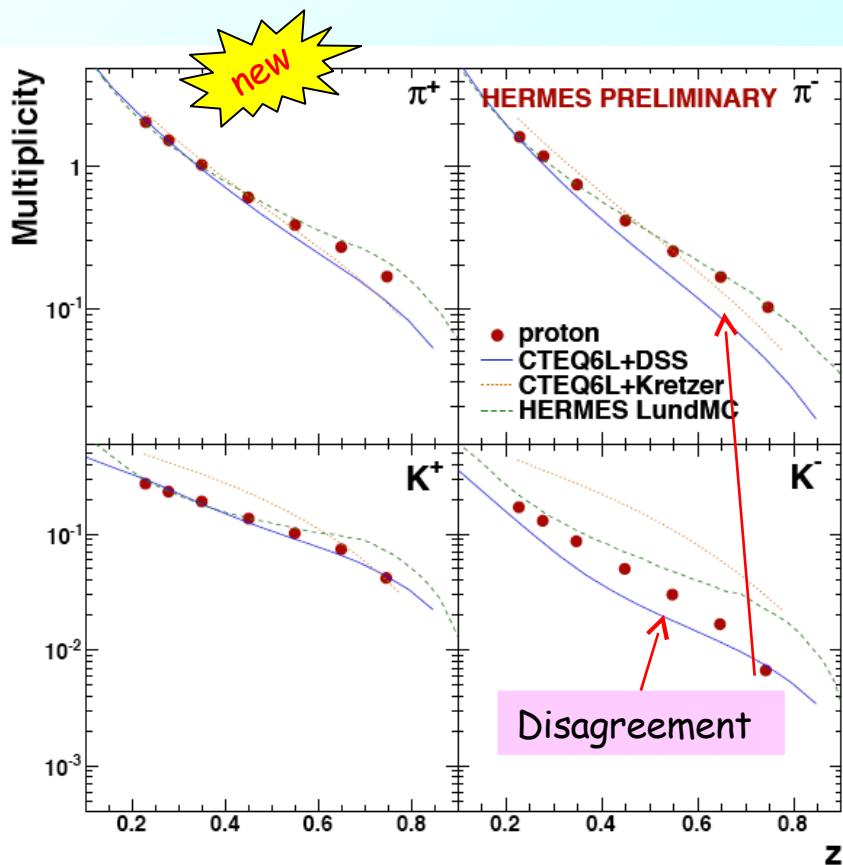
$FF(z, Q^2)$: Fragmentation Function - $D_{1^q \rightarrow h}(z, Q^2)$, $H_{1^{\perp q} \rightarrow h}(z, Q^2)$, ...

Charged-hadron multiplicities I

LO interpretation:

$$\sigma_{UU} \propto f_1^q \otimes D_1^{q \rightarrow h}$$

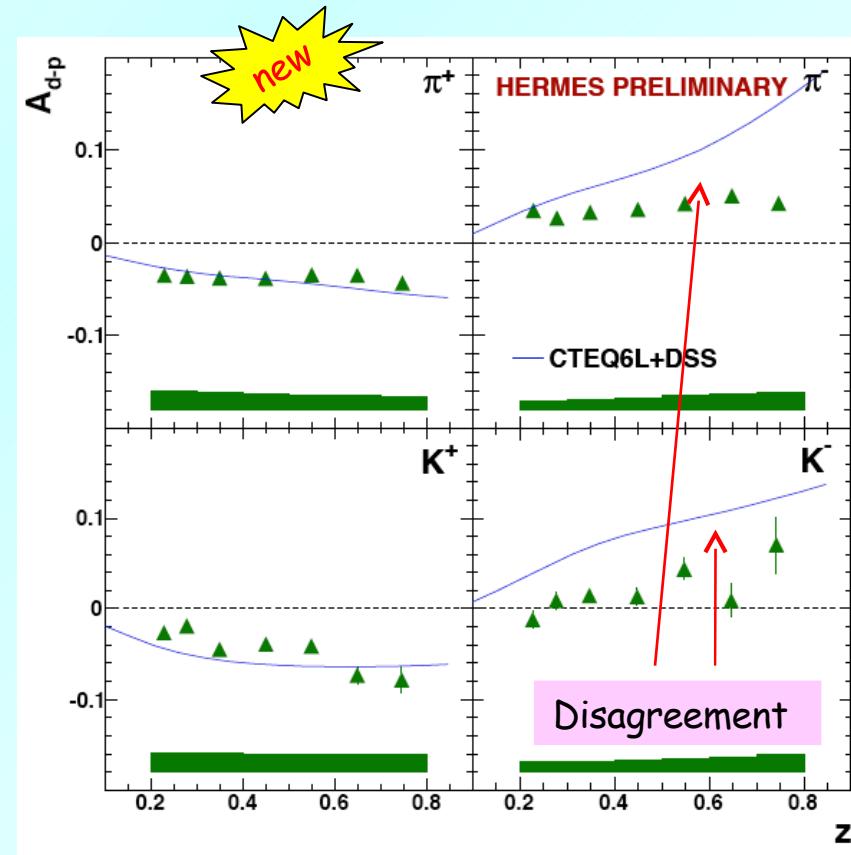
$$M_N^h = \frac{1}{N_N^{DIS}(Q^2)} \frac{dN_N^h(z, Q^2)}{dz} = \frac{\sum_q e_q^2 \int dx f_{1q}(x, Q^2) D_{1q}^h(z, Q^2)}{\sum_q e_q^2 \int dx f_{1q}(x, Q^2)}$$



Proton-deuteron asymmetry

$$A_{d-p}^h \equiv \frac{M_d^h - M_p^h}{M_d^h + M_p^h}$$

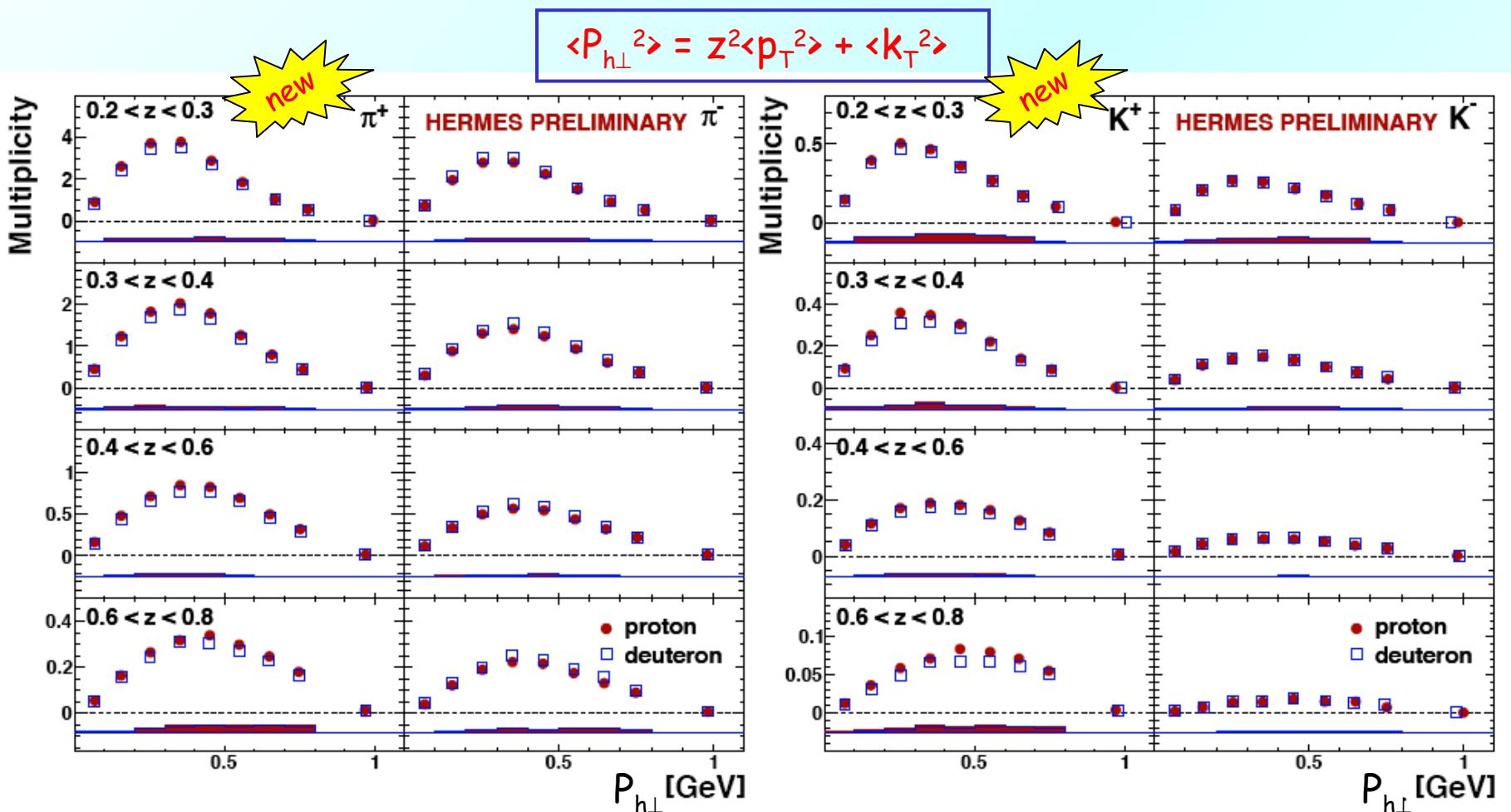
Reflects different flavor content
Correlated systematics cancel



Charged-hadron multiplicities II

- Disentanglement of z and $P_{h\perp}$ dependences
- Access to intrinsic quark p_T and fragmentation k_T

$$\sigma_{UU} \propto f_1^q \otimes D_1^{q \rightarrow h}$$



Double-spin asymmetry A_1^h

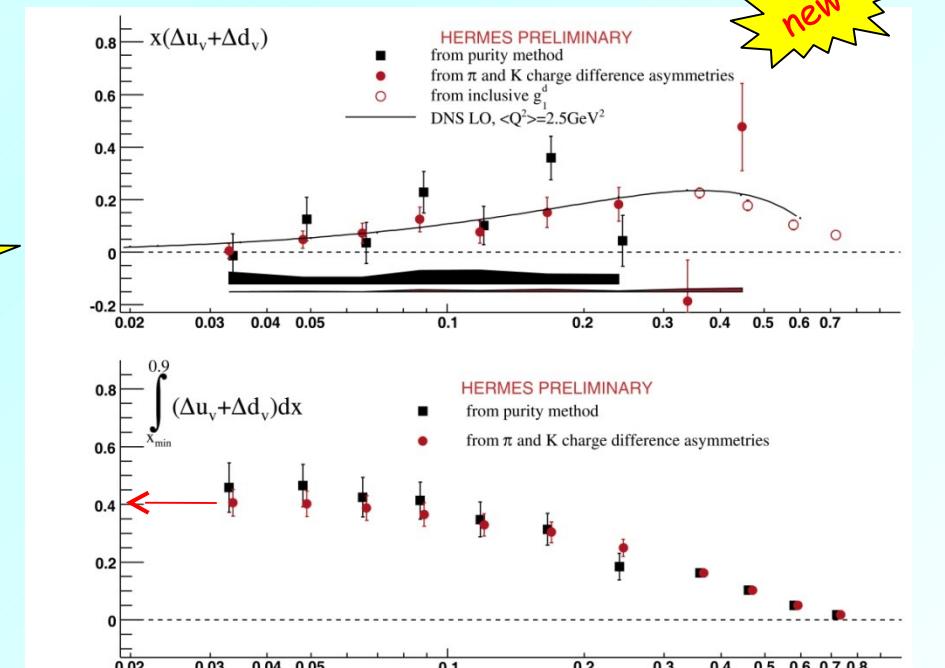
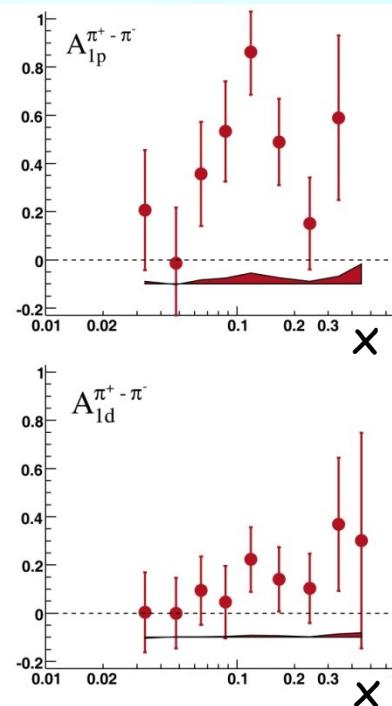
Refined studies extending the work
in Phys. Rev. D 71 (2005) 012003

$$\sigma_{LL} \propto g_1^q \otimes D_1^{q \rightarrow h}$$

With charge conjugation
symmetry in fragmentation

$$D_{1,q}^{h+} = D_{1,q}^{h-}$$

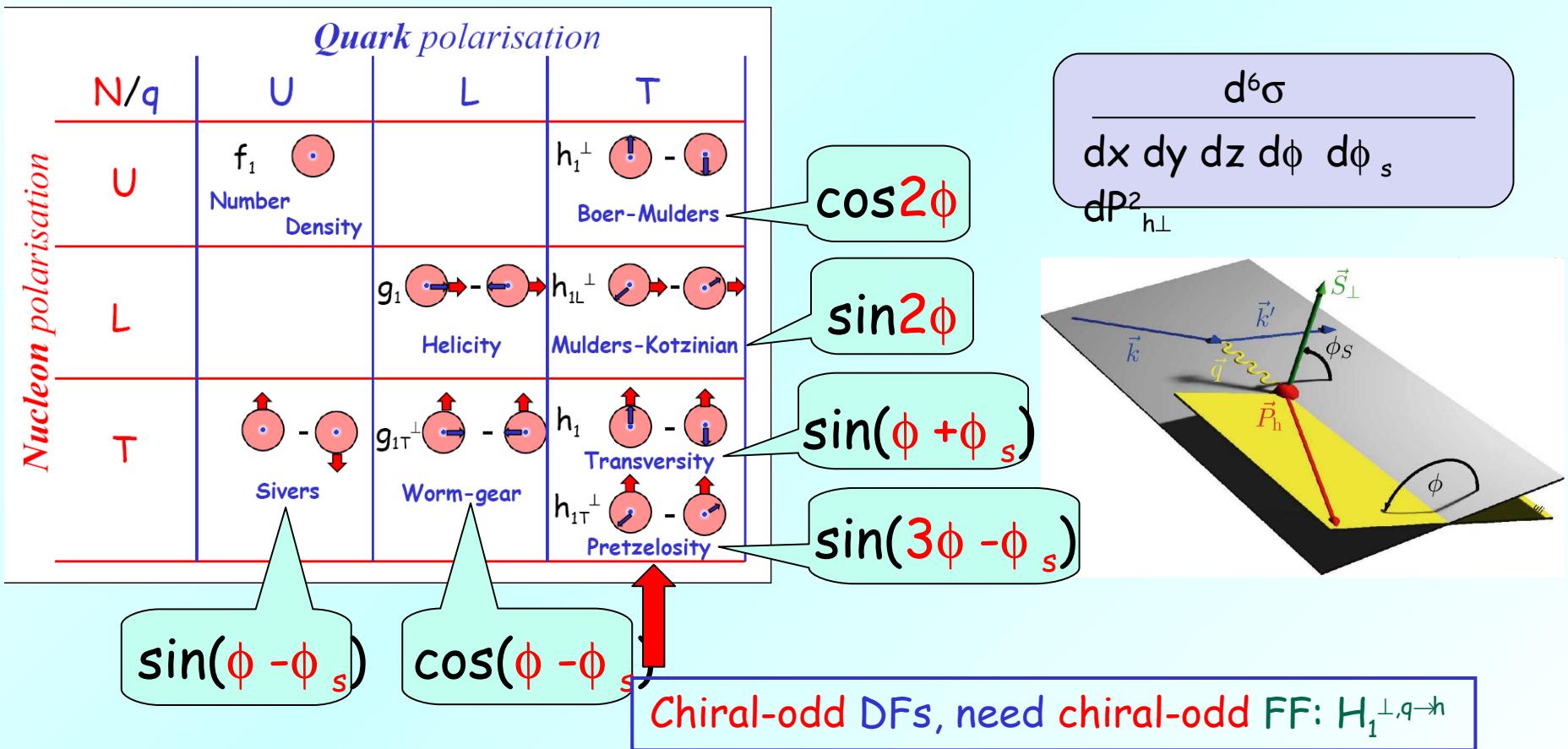
$$A_{1d}^{h^+-h^-} = \frac{\Delta u_v +}{\Delta d_w + d_v}(x)$$



Integral over sum of valence distributions
compatible with $\Delta\Sigma$
Sea contribution to nucleon spin small

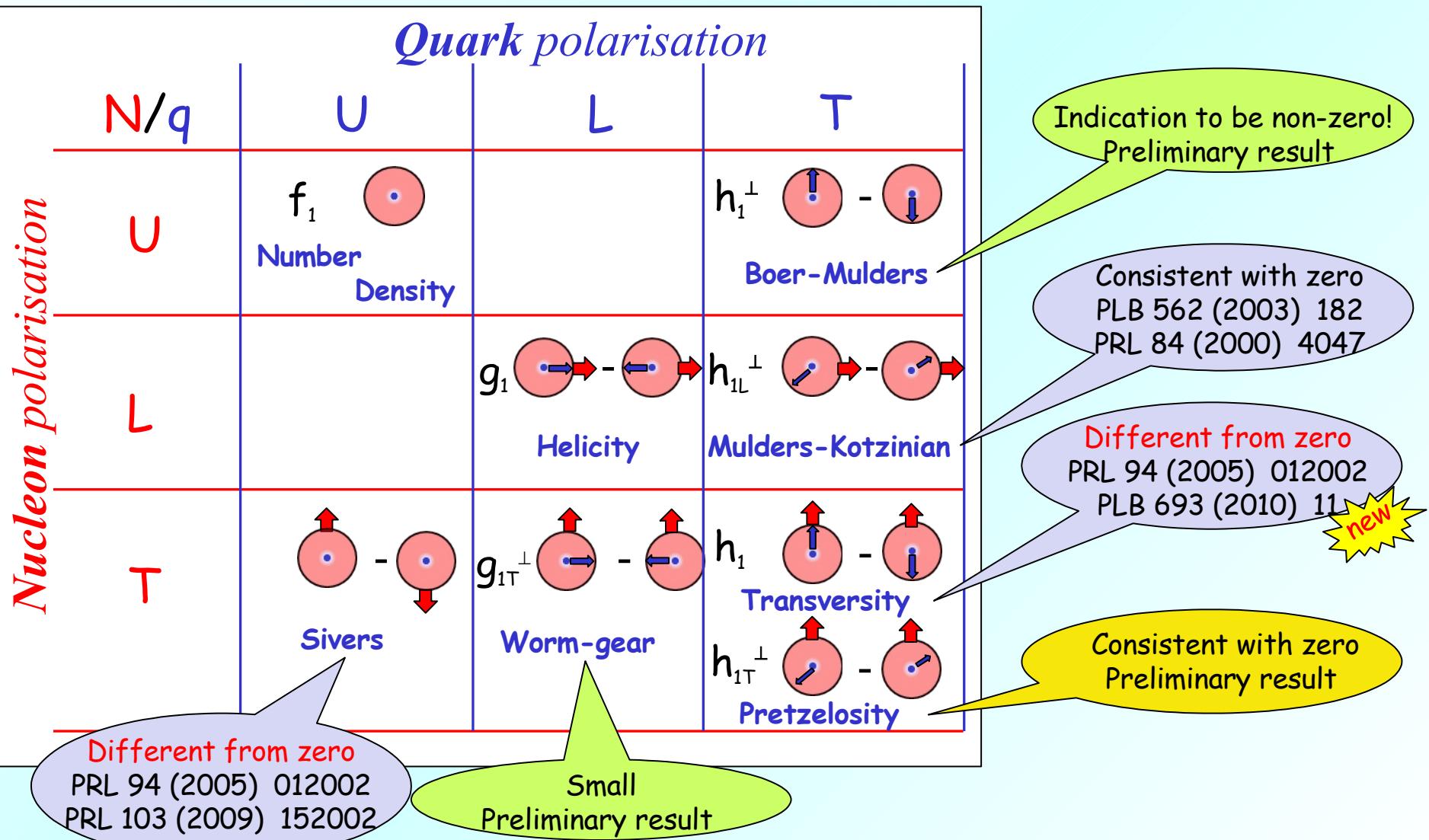
Leading-twist TMDs

- Nucleon structure described in leading-twist by 8 transverse-momentum dependent quark distributions (TMDs)
- HERMES has access to all of them through specific azimuthal modulations (ϕ , ϕ_s) of the cross section thanks to the polarised beam and target

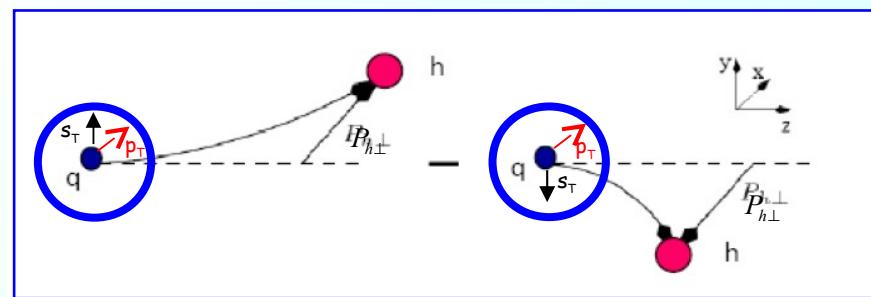


Leading-twist TMDs

Pioneering measurements by HERMES



Boer-Mulders DF $h_1^{\perp,q}$



\bar{q}/q	U	L	T
U	f_{1q}		h_{1q}^\perp
L		g_{1q}^\perp	h_{1L}^\perp
T	f_{1T}	g_{1T}	h_1 h_{1T}

transversely polarised quarks
with p_T in unpolarised nucleon

$$\sigma_{UU} \cos 2\phi \propto h_1^{\perp,q} \otimes H_1^{\perp,q} \rightarrow$$

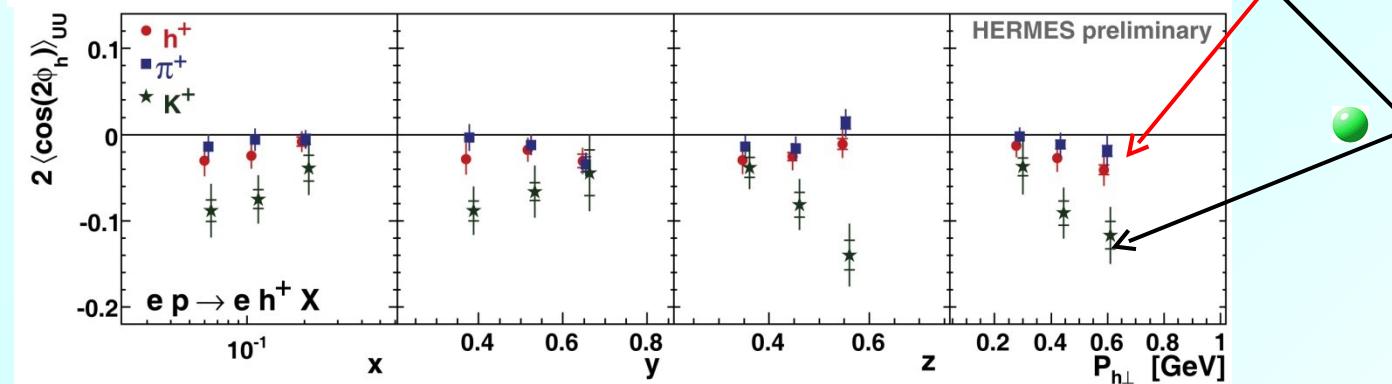
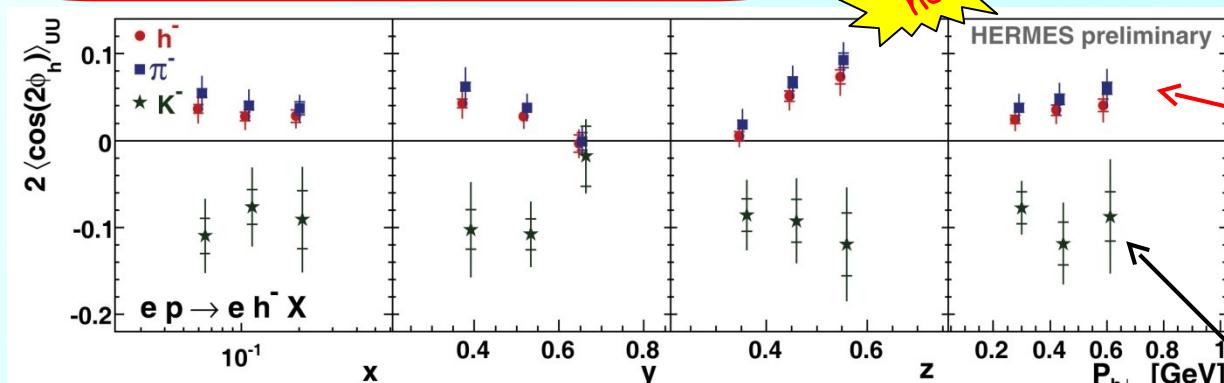
h_1^\perp is chiral-odd and
naive T-odd (like
 f_{1T}^\perp) requires
FSI/ISI

Opposite sign for
 π^+ and π^- , larger
magnitude for π^-

$h_{1,u}^\perp$ and $h_{1,d}^\perp$
have same sign

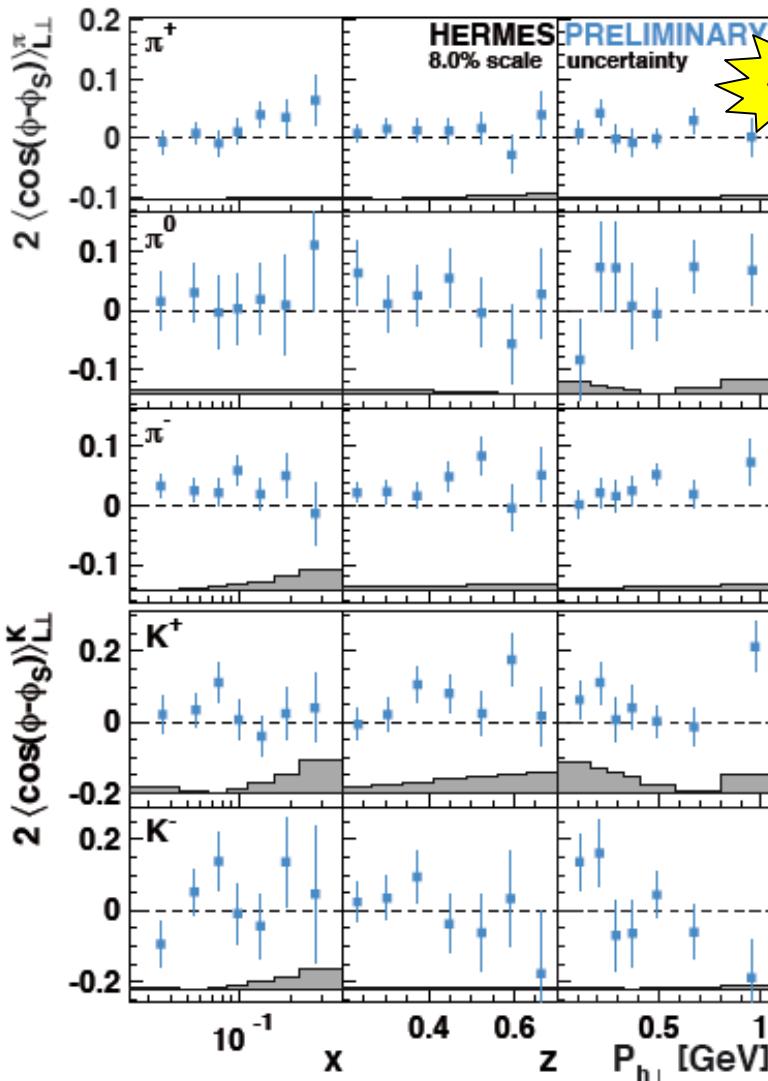
Large signal with
same sign for K^\pm

sea fragmentation
important



Worm-gear DF $g_{1T}^{\perp,q}$

longitudinally polarised quarks in transversely polarised nucleon



$$\sigma_{LT} \cos(\phi - \phi_s) \propto g_{1T}^{\perp,q} \otimes D_1^{q-h}$$

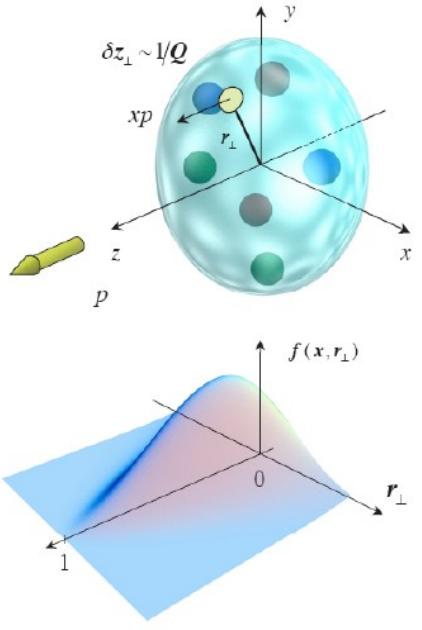
- Related to parton orbital motion: requires interference between wave functions with OAM difference by 1 unit

Slightly non-zero

- $g_{1T}^{\perp,q} = -h_{1L}^{\perp,q}$ (supported by many models)
 - $$g_{1T}^{\perp,q} \approx \frac{1}{x} \int_x^\infty \frac{dy}{y} g_1^q(y)$$
- (Wandzura-Wilczek type approximation)

Exclusive Measurements

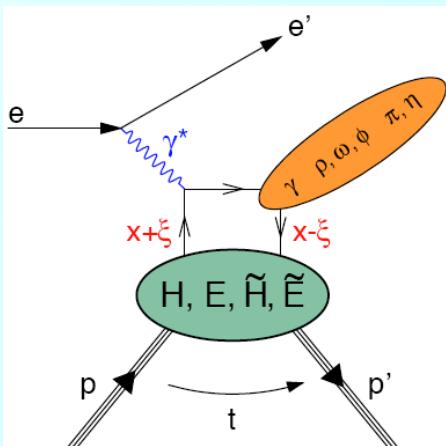
Generalised parton distributions



- Generalisation of Form Factors (moments of GPDs) and PDFs (forward limit)
- Correlated information about **longitudinal momentum xp** and **transverse spatial position r_{\perp}**
- Ji relation:

$$J_q = 1/2 \Delta \Sigma + L_q = \lim_{t \rightarrow 0} \int_{-1}^1 dx \times [H(x, \xi, t) + E(x, \xi, t)]$$

Access: exclusive processes



- Final state sensitive to different **GPDs**
- Spin- $\frac{1}{2}$ target: 4 chiral-even leading-twist quark **GPDs**
- H, \tilde{H} (E, \tilde{E}) conserve (flip) nucleon helicity
- Vector mesons (p, ω, ϕ) H, E
- Pseudoscalar mesons (π, η) \tilde{H}, \tilde{E}
- DVCS (γ) $H, E, \tilde{H}, \tilde{E}$

Hard exclusive ρ^0 -meson production I

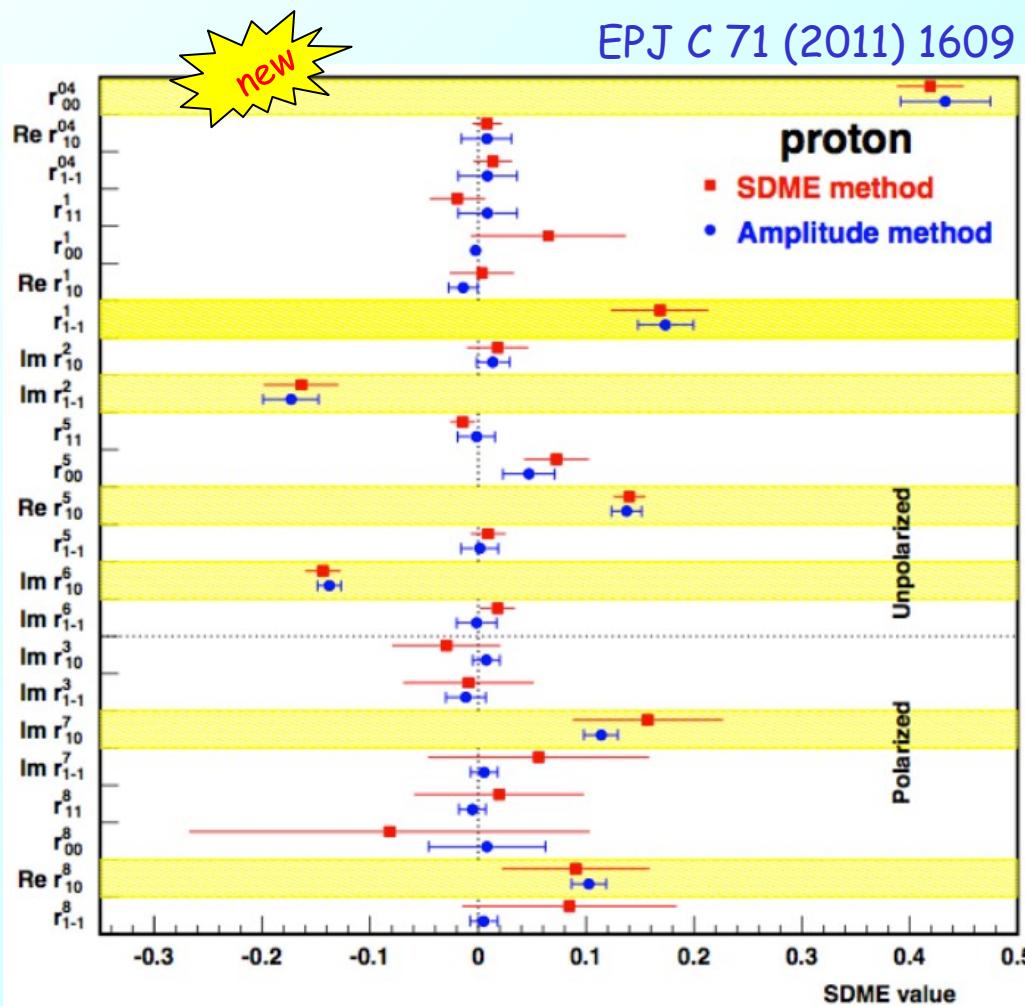
Meson SDMEs

EPJC 62 (2009) 659

Photon SDMEs

$$r_{\lambda_V \mu_V}^{\eta} = \frac{1}{2N} \sum_{\lambda_\gamma \mu_\gamma \lambda'_N \lambda_N} F_{\lambda_V \lambda'_N \lambda_\gamma \lambda_N} \Sigma_{\lambda_\gamma \mu_\gamma}^{\eta} F_{\mu_V \lambda'_N \mu_\gamma \lambda_N}^*$$

EPJ C 71 (2011) 1609



Helicity amplitudes

$$F_{\lambda_V \lambda_\gamma} = T_{\lambda_V \lambda_\gamma} + U_{\lambda_V \lambda_\gamma}$$

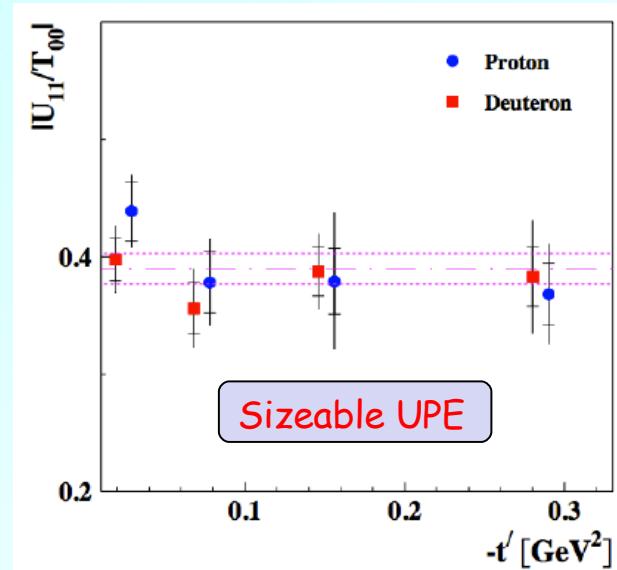
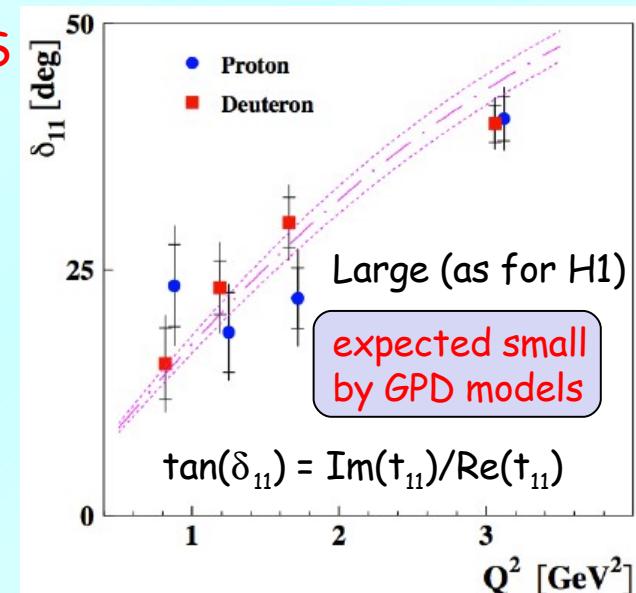
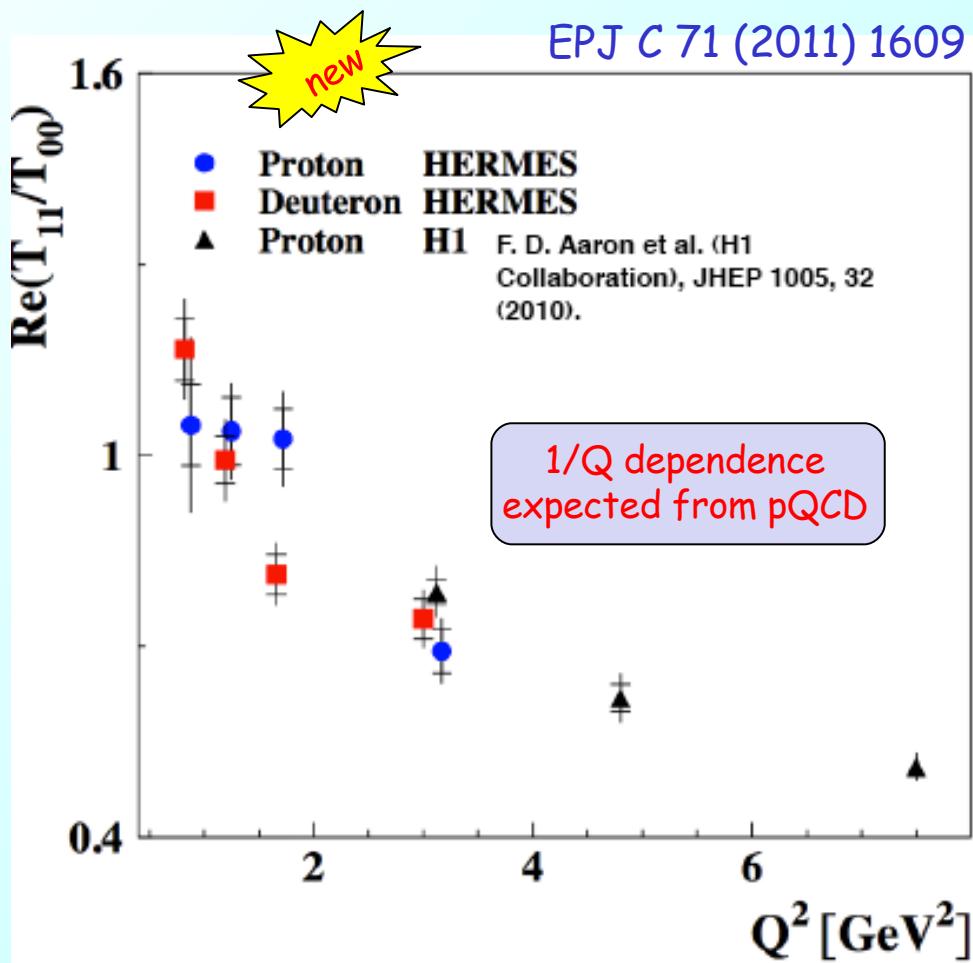
- Helicity amplitudes are the fundamental quantities to be compared with theory
- They form a basis for the SDMEs
- Re-derived SDMEs consistent with published ones
- Enhanced sensitivity for polarised SDMEs

Hard exclusive ρ^0 -meson production II

Hierarchy predicted by theory, confirmed by HERMES

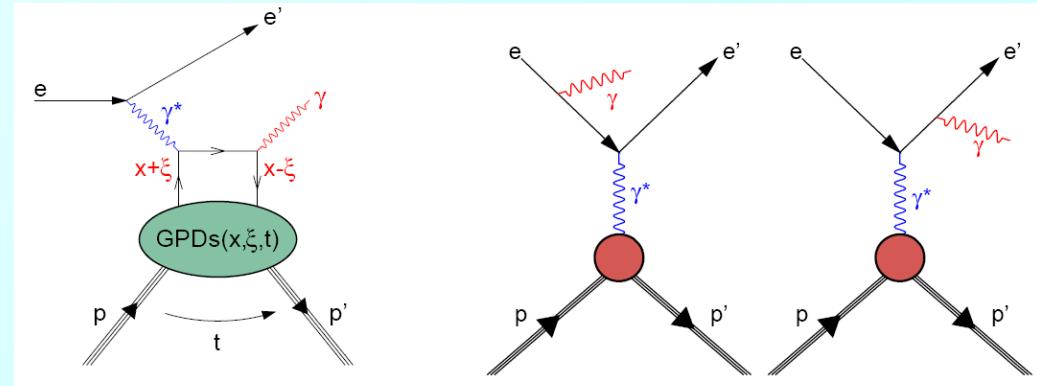
$$|T_{00}|^2 \approx |T_{11}|^2 \gg |U_{11}|^2 > |T_{01}|^2 \gg |T_{10}|^2 \dots$$

$$\gamma^* L \rightarrow p_L \quad \gamma^* T \rightarrow p_T \quad \gamma^* T \rightarrow p_L \quad \gamma^* L \rightarrow p_T$$



Deeply Virtual Compton Scattering & GPDs

- Theoretically cleanest way to access GPDs
- Interference between DVCS and Bethe-Heitler amplitude
- $|T_{\text{DVCS}}| \ll |T_{\text{BHL}}| @ \text{HERMES}$

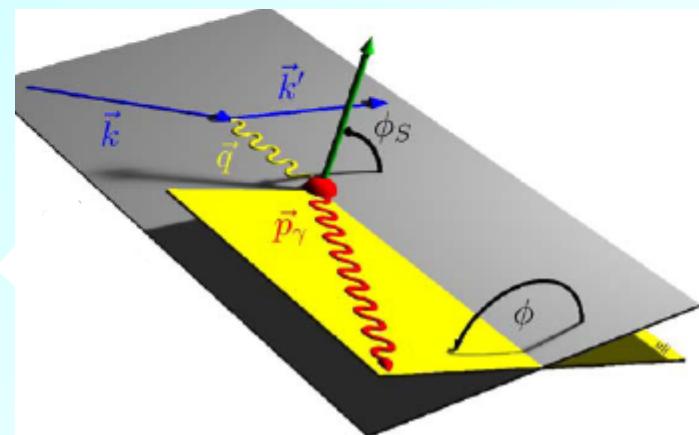


- Access to GPD combinations through azimuthal asymmetries

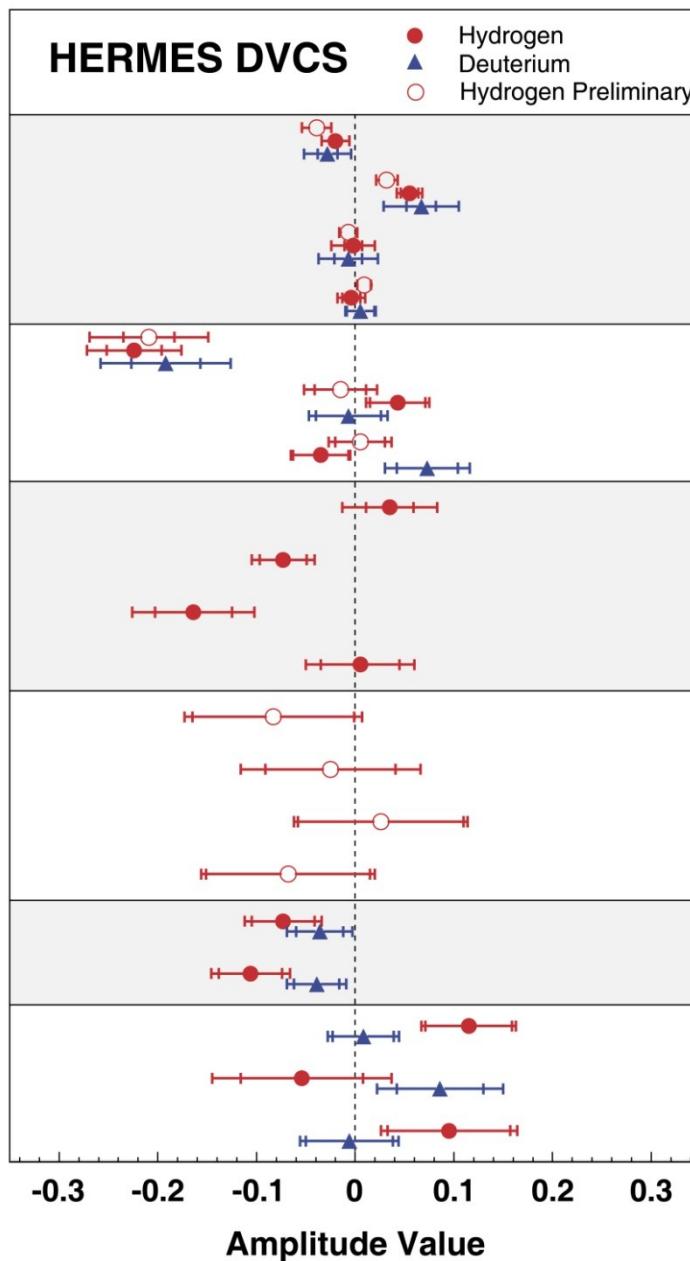
A_{XY}
beam target
polarisation

HERMES: Complete set of asymmetries

- Both beam charges
- Both beam helicities
- Unpolarised H, D and nuclear targets
- Longitudinally polarised H and D targets
- Transversely polarised H target



DVCS asymmetries measured @ HERMES

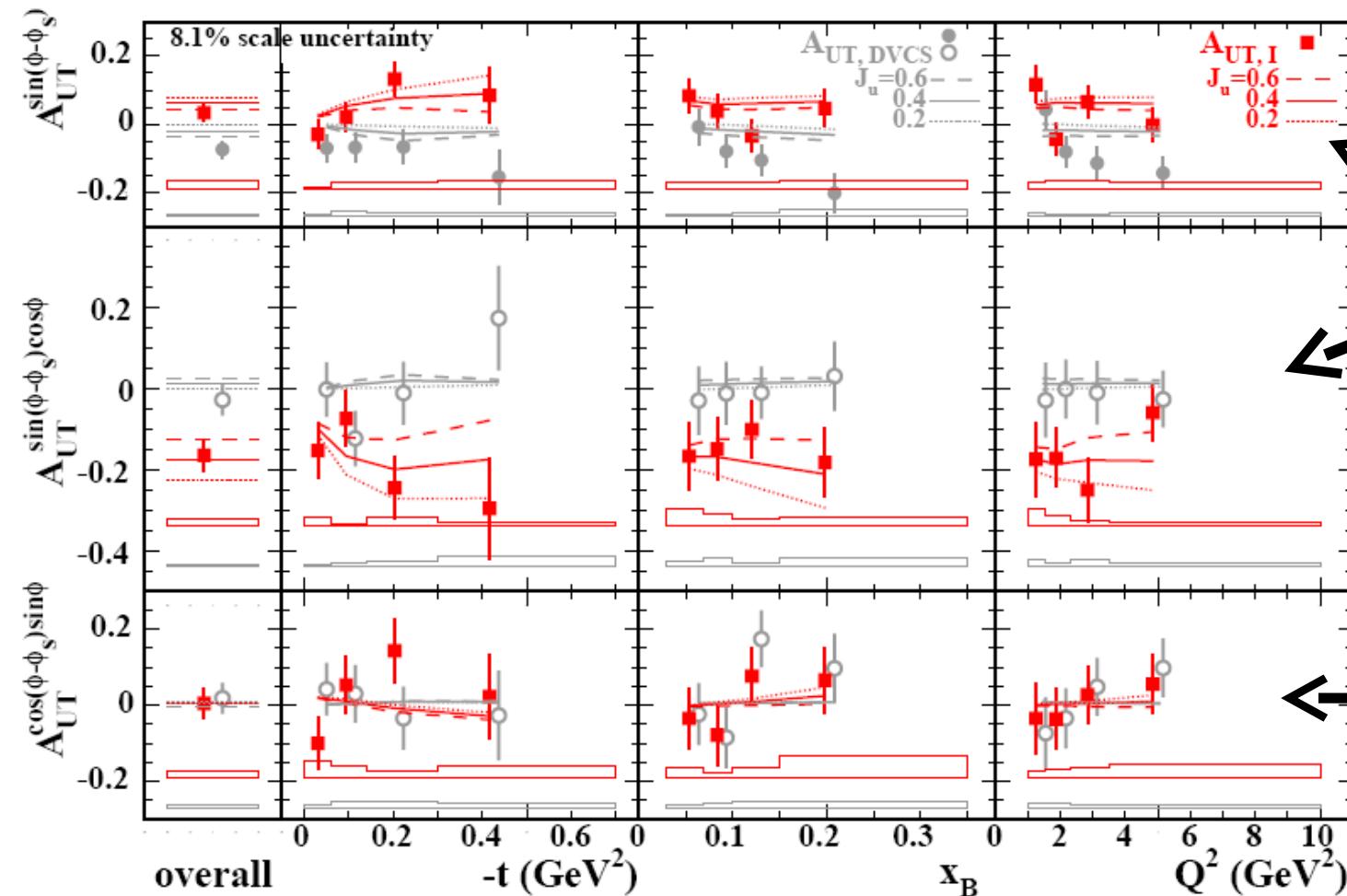


- Beam charge asymmetry
GPD H
H: PRL 87 (2001) 182001
PR D 75 (2007) 011103
JHEP 11 (2009) 083
D: Nucl. Phys. B 829 (2010) 1
- Beam helicity asymmetry
GPD H
- Transverse target-spin asymmetry
GPD E
H: JHEP 06 (2008) 066
- Transverse double-spin asymmetry
GPD E
H: arXiv:1106.2990 **new!**
- Longitudinal target spin asymmetry
GPD \tilde{H}
H: JHEP 06 (2010) 019
D: Nucl. Phys. B 842 (2011) 265 **new!**
- Longitudinal double spin asymmetry
GPD \tilde{H}

DVCS: transverse target asymmetry A_{UT}

Sensitive to GPD E

JHEP 06 (2008) 066



sensitive to J_u :

$$\text{Im}(F_2 \mathcal{H} - F_1 \mathcal{E}) \cdot \sin(\phi - \phi_s) \cos(n\phi)$$

NOT sensitive to J_u :

$$\text{Im}(F_2 \tilde{\mathcal{H}} - (F_1 + \xi F_2) \tilde{\mathcal{E}}) \cdot \cos(\phi - \phi_s) \sin \phi$$

Model: VGG with variation of J_u , while $J_d=0$

DVCS transverse double-spin asymmetry A_{LT}

Beam charge

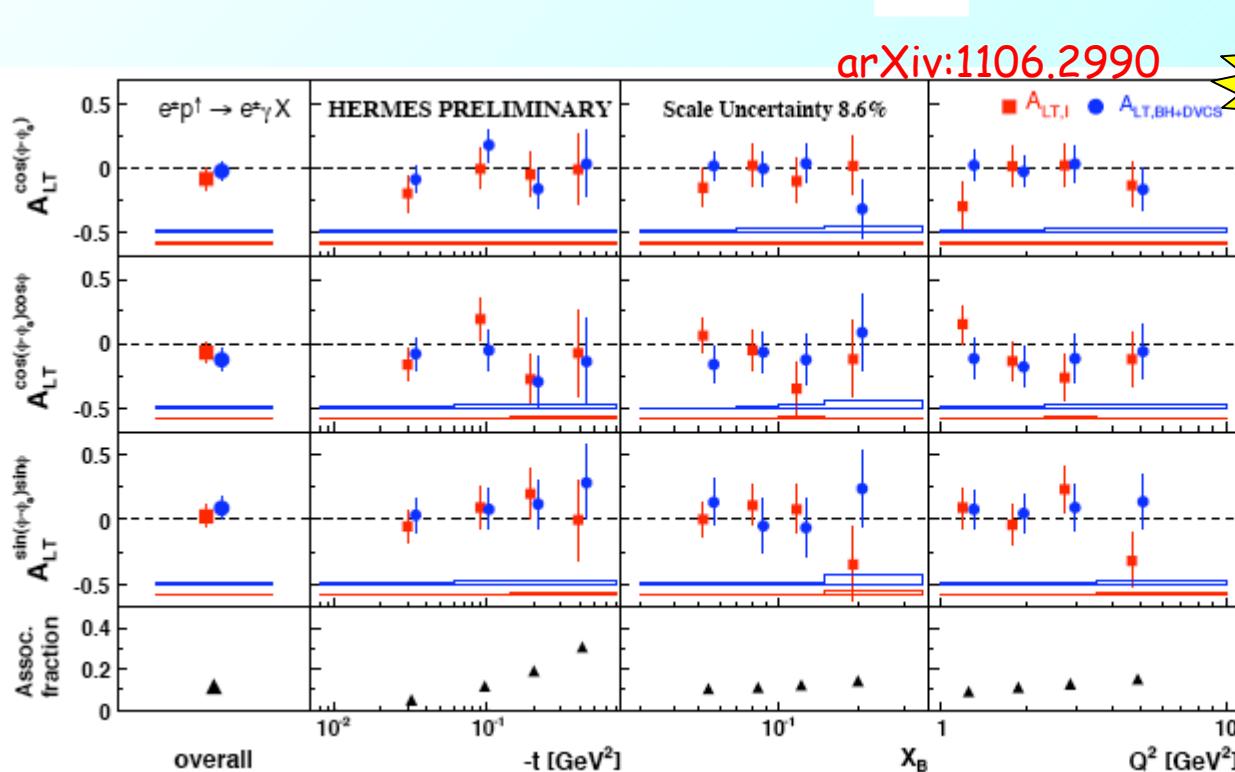
Beam polarisation

Target polarisation

$$\langle \mathcal{N}(e_\ell, P_l, S_t, \phi, \phi_S) \rangle \propto \sigma_{UU}(\phi) [1 + \dots + P_l S_t \mathcal{A}_{LT}^{BH+DVCS} + e_\ell P_l S_t \mathcal{A}_{LT}^I]$$

arXiv:1106.2990

new



$$\propto A_{LT}^{\cos(\phi - \phi_S) \cos(\phi)}$$

$$\propto \frac{\text{Re}[F_2 \tilde{\mathcal{H}} - (F_1 + \xi F_2) \tilde{\mathcal{E}}]}{\text{Re}[\mathcal{H}\mathcal{E}^* - \mathcal{E}\mathcal{H}^* - \xi(\tilde{\mathcal{H}}\tilde{\mathcal{E}}^* - \tilde{\mathcal{E}}\tilde{\mathcal{H}}^*)]}$$

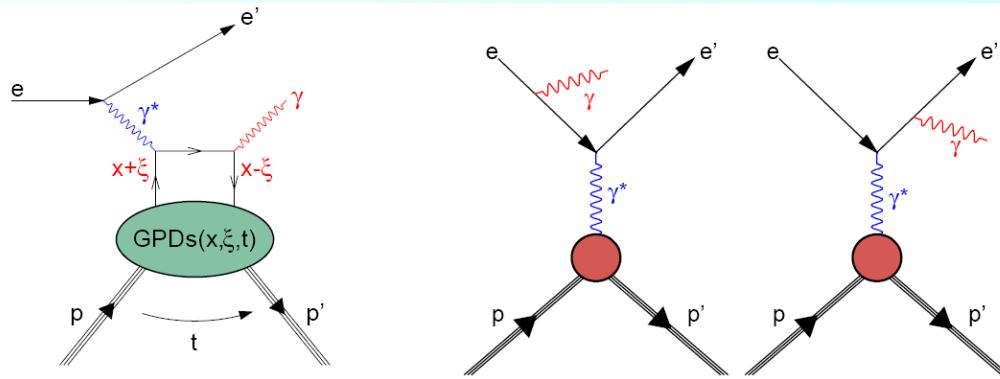
$$\propto \frac{\text{Re}[F_2 \mathcal{H} - F_1 \mathcal{E}]}{\text{Re}[-\tilde{\mathcal{H}}\mathcal{E}^* - \tilde{\mathcal{H}}^*\mathcal{E} + \xi(\mathcal{H}\tilde{\mathcal{E}}^* + \tilde{\mathcal{E}}\mathcal{H}^*)]}$$

Sensitive to both GPDs
entering the Ji sum rule

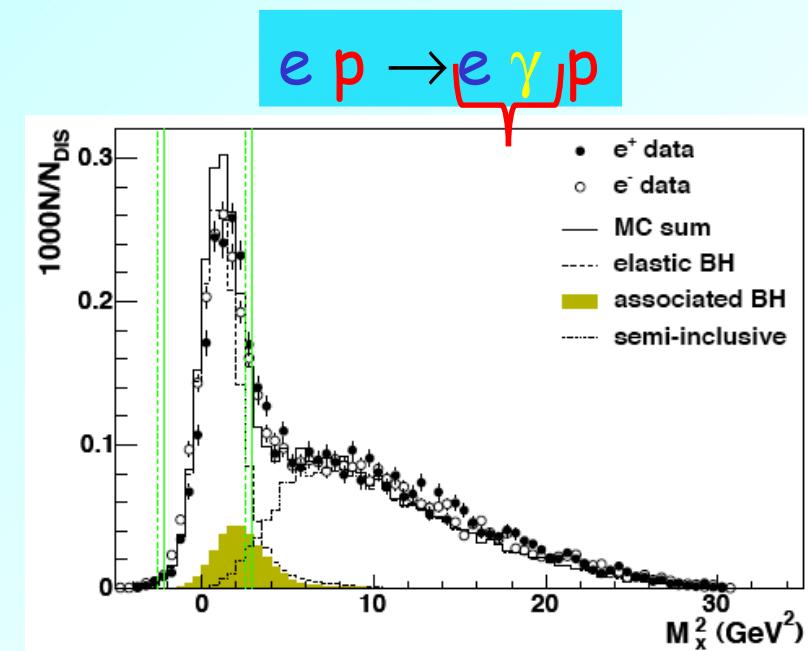
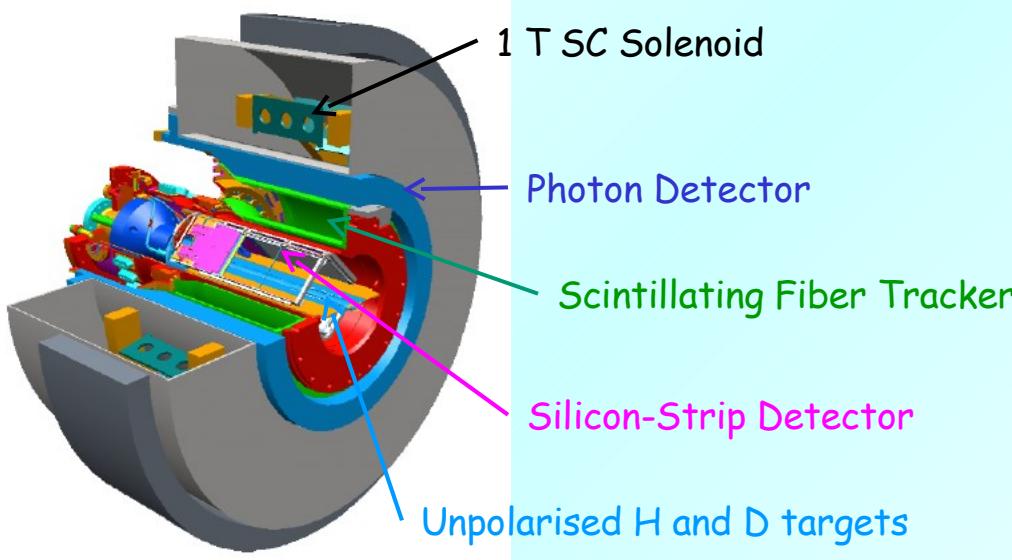


Consistent with zero, cancellations between E and H
Sensitivity to J_u suppressed by kinematic factors

DVCS with Recoil Detector



Recoil Detector to tag exclusivity

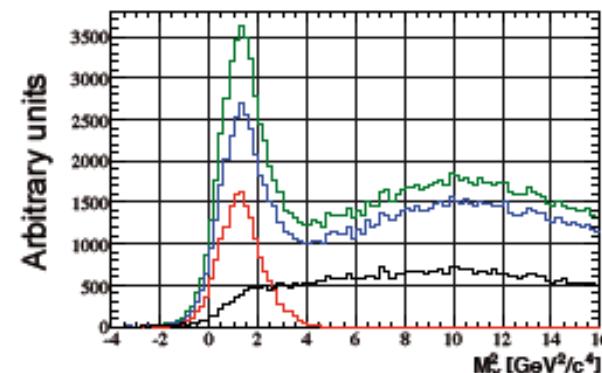
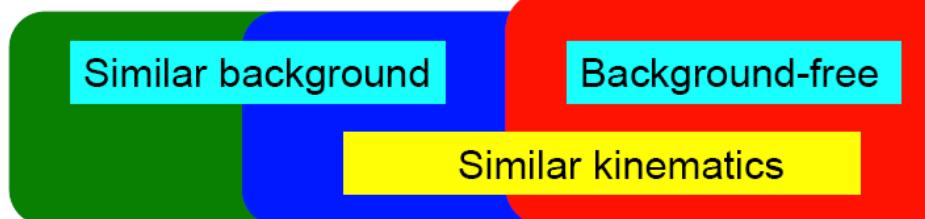


DVCS with Recoil Detector

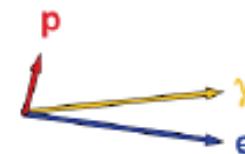
Without Recoil Detector

In Recoil Detector acceptance

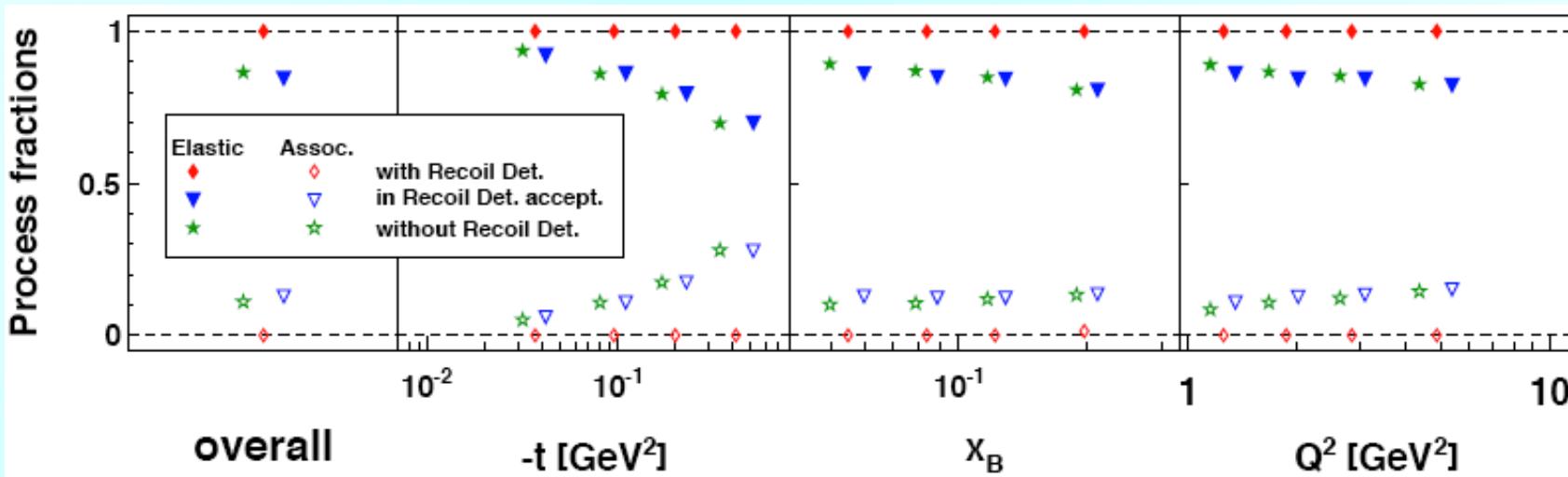
With Recoil Detector



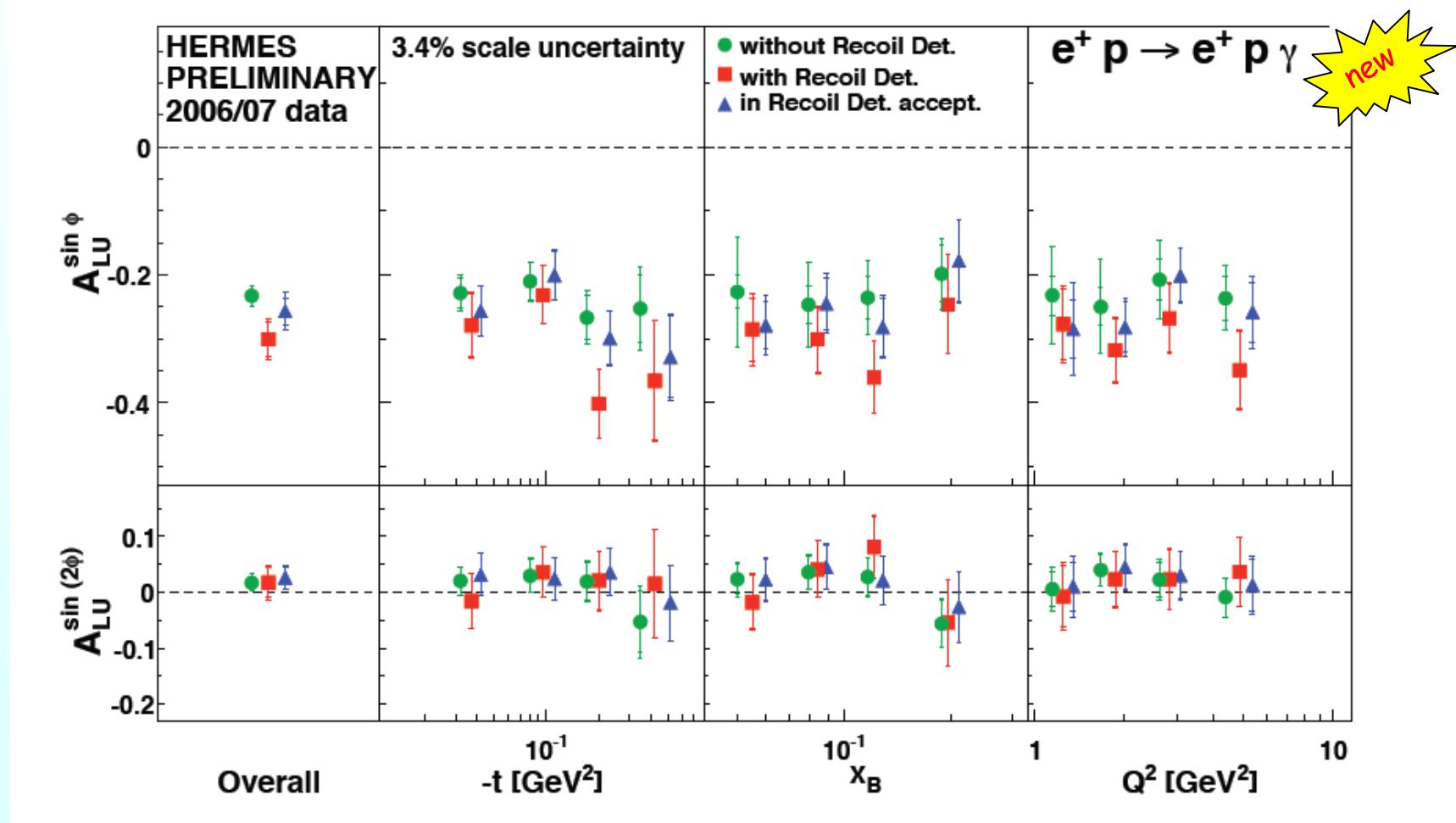
Kinematic event fitting technique: all 3 particles in the final state detected should satisfy
4-constraints on energy-momentum conservation



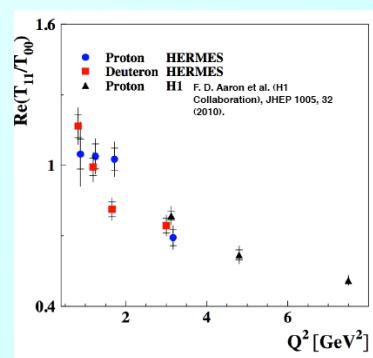
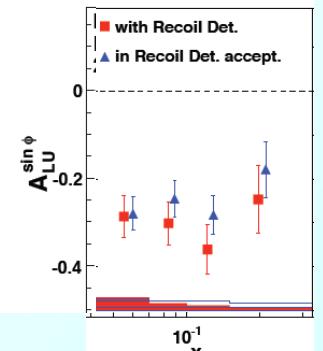
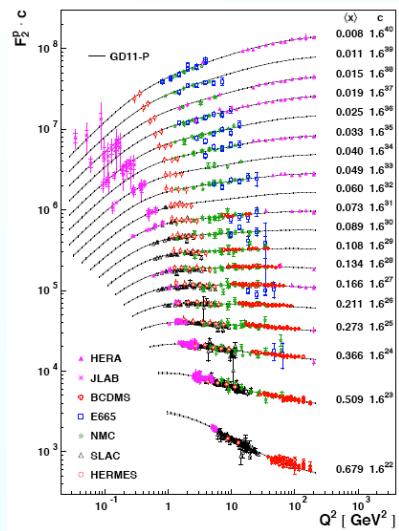
- No requirement for Recoil
- Charged recoil track in acceptance
- Kinematic fit probability > 1 %
- Kinematic fit probability < 1 %



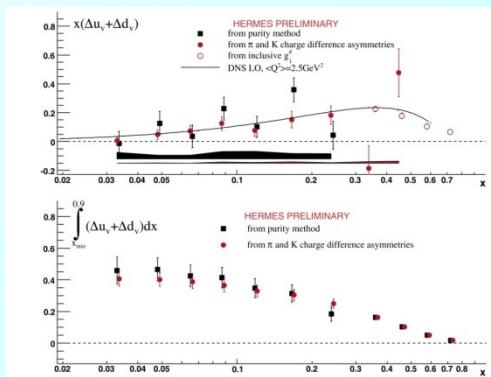
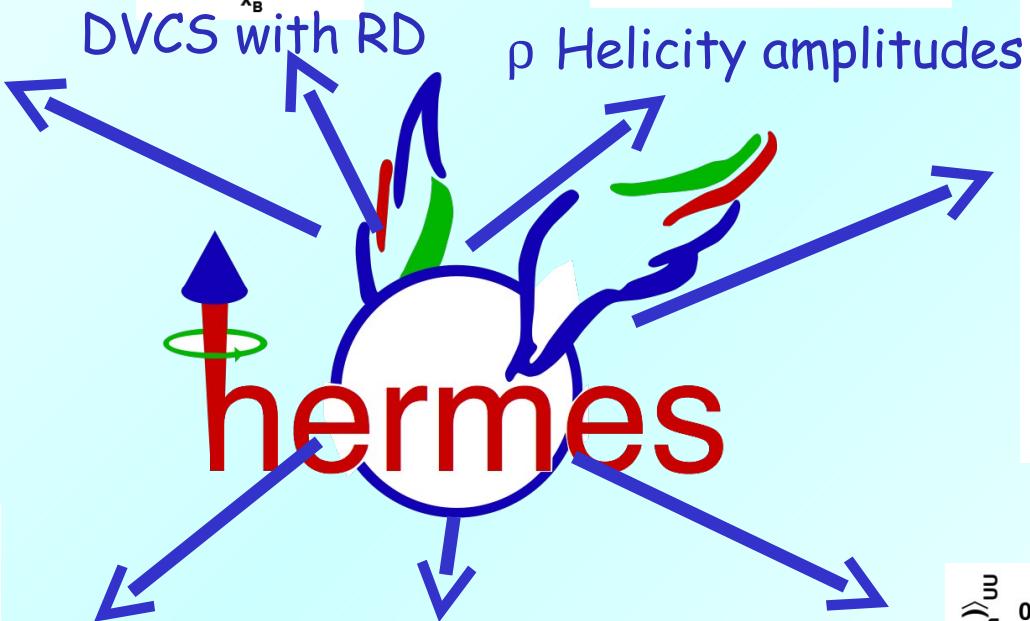
Pure elastic DVCS



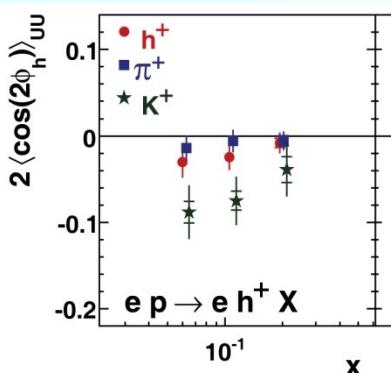
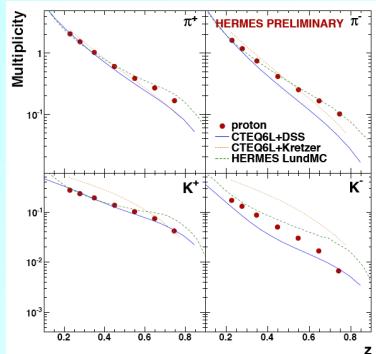
Indication that leading amplitude for pure elastic process is slightly larger than for unresolved signal (elastic + associated)



F_2



$x(\Delta u_v + \Delta d_v)$

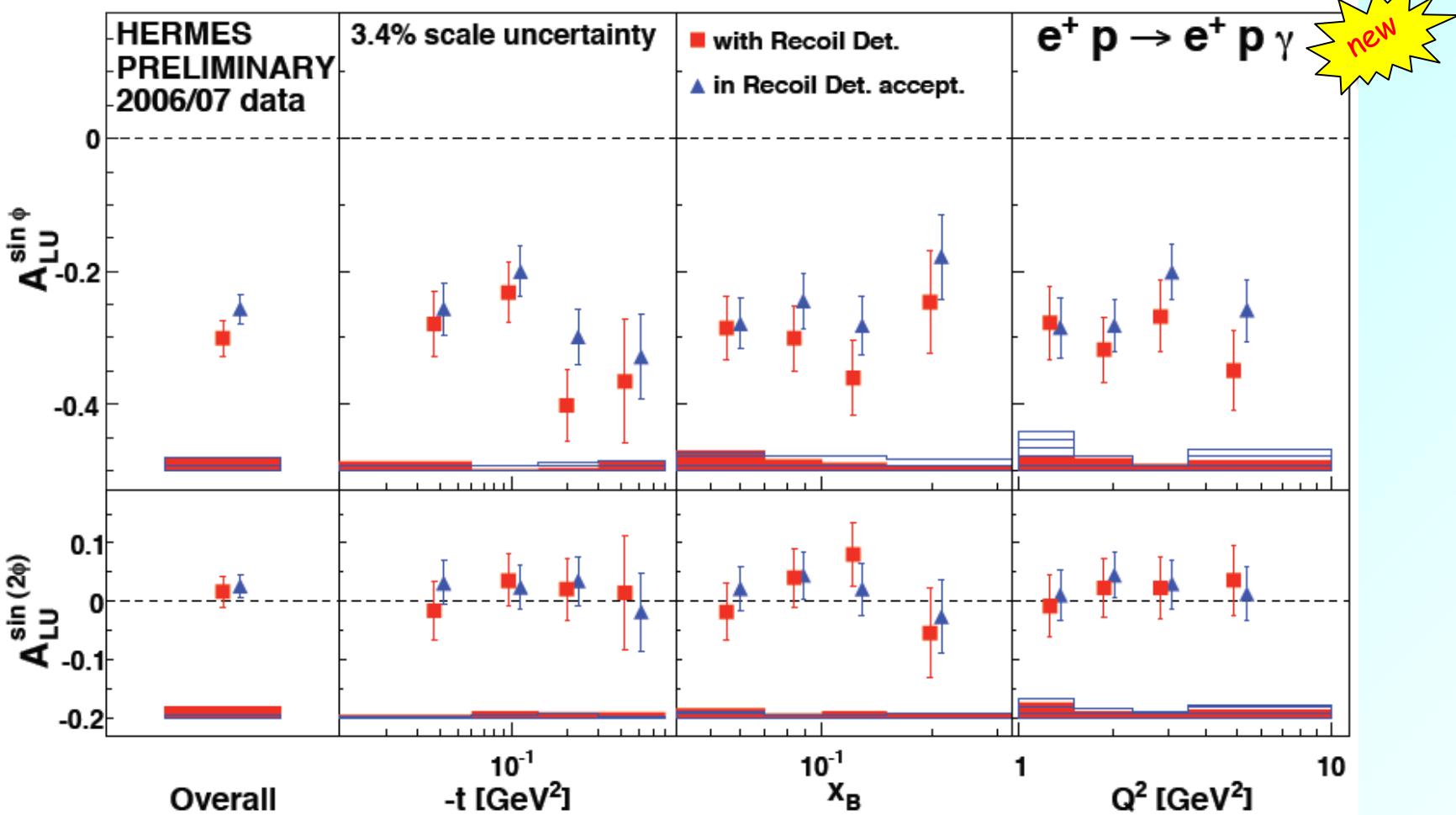


Boer-Mulders DF

30

Backups

Pure elastic DVCS



● Indication that leading amplitude for pure elastic process is slightly larger than for unresolved signal (elastic + associated)

Double-spin asymmetry A_1^h

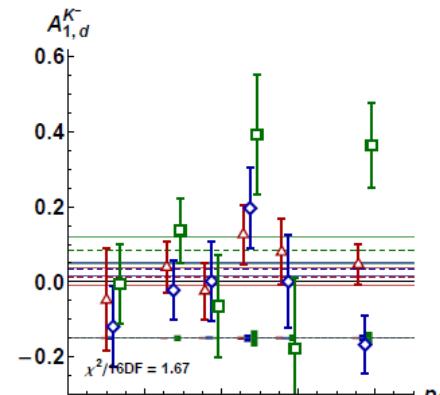
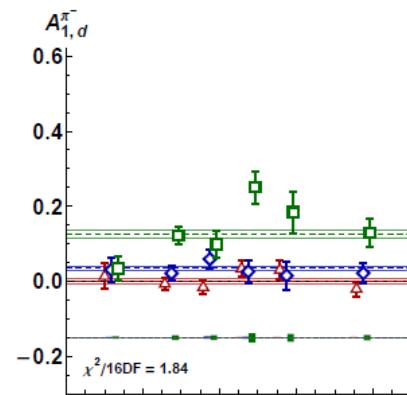
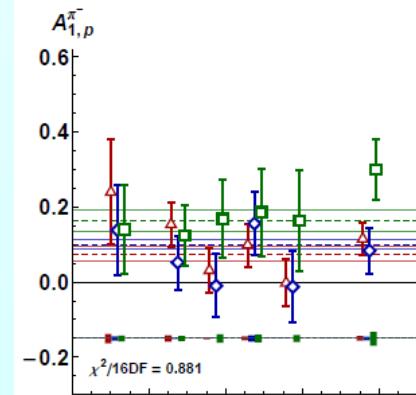
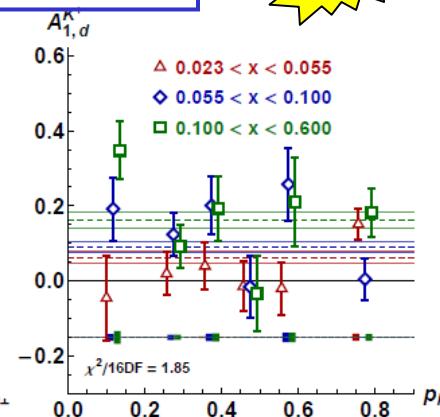
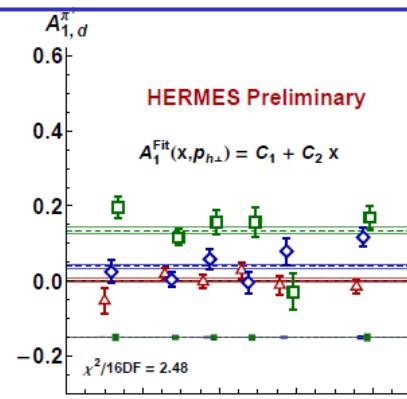
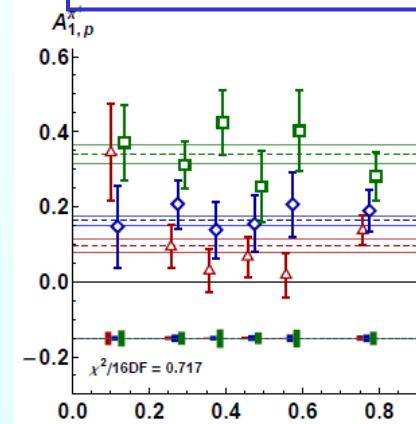
Refined studies extending the work
in Phys. Rev. D 71 (2005) 012003

$$\sigma_{LL} \propto g_1^q \otimes D_1^{q \rightarrow h}$$

$A_1^h(x, P_{h\perp})$

2D - dependencies

new



$P_{h\perp}$ [GeV]

- Sensitive to differences in transverse momentum dependence of g_1 and f_1

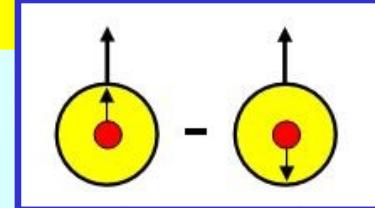
- No significant $P_{h\perp}$ dependence observed

Transversity, Collins Amplitudes

Transversity DF

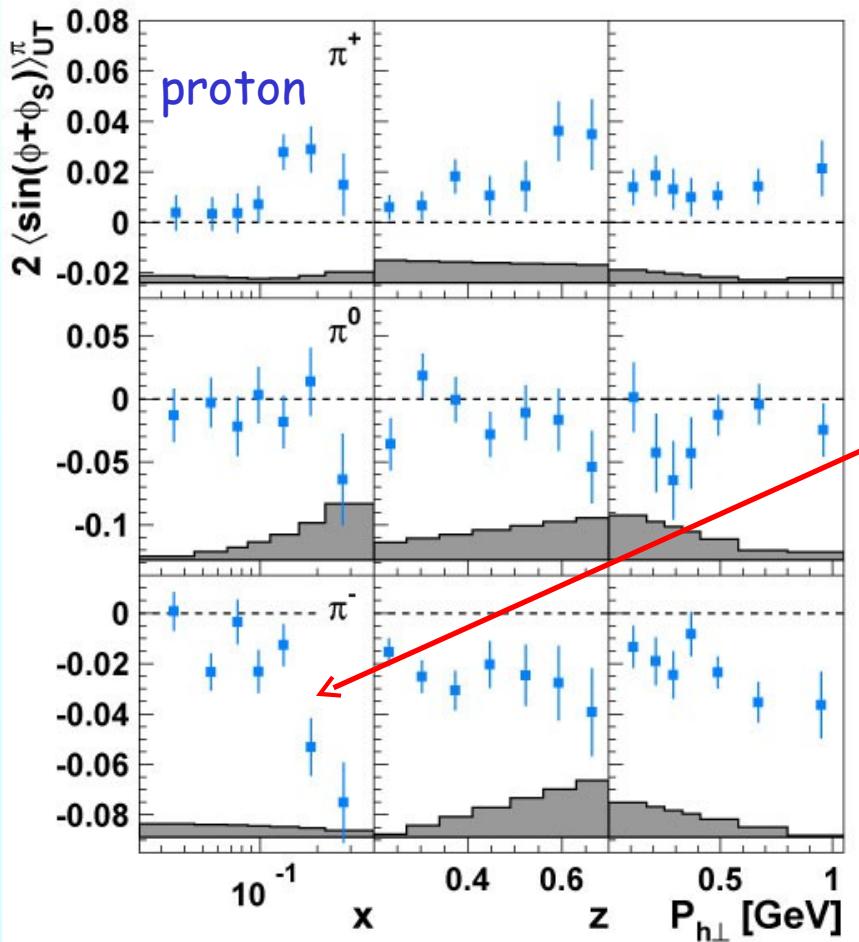
$$2\langle \sin(\phi + \phi_s) \rangle \quad h_{UT}^q \propto h_1^q(x) \otimes H_1^{\perp q}(z)$$

Collins FF



N/q	U	L	T
U	f_1		h_{1T}
L		g_1	h_{1L}
T	f_{1T}	g_{1T}	h_1

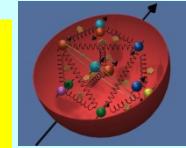
arXiv:1006.4221



- Both Collins fragmentation function and transversity distribution function are sizeable
- Surprisingly large π^- asymmetry
- Possible source: large contribution (with opposite sign) from unfavored fragmentation,

$$H_{1\perp, \text{disf}} \approx - H_{1\perp, \text{fav}}$$

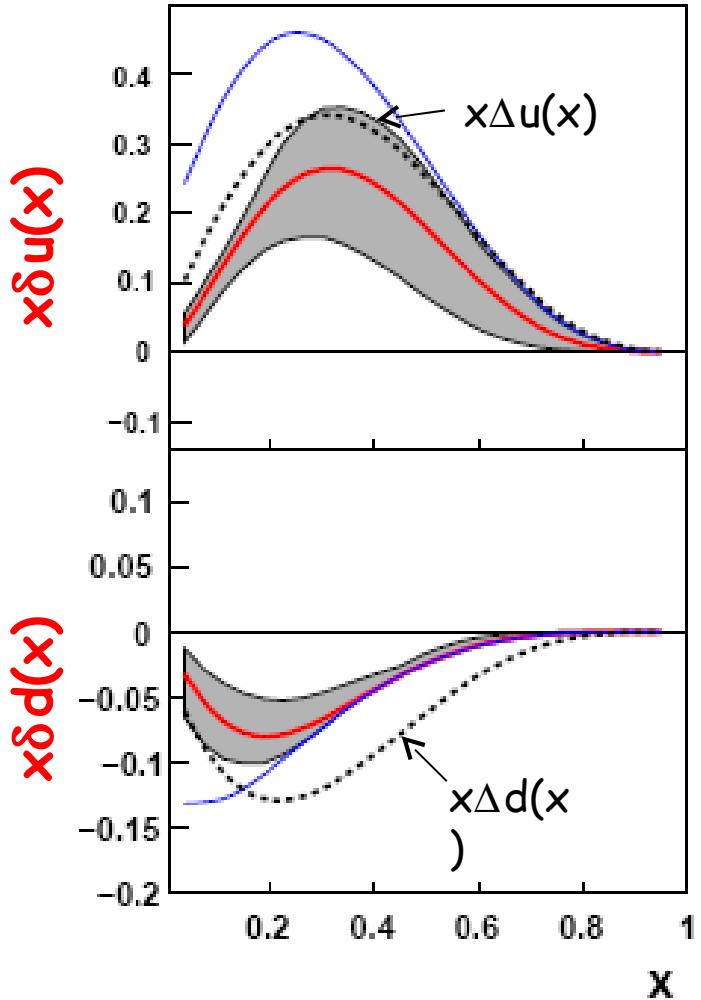
Extraction of Transversity



Fit to HERMES ($e p \xrightarrow{\uparrow} e h X$), COMPASS ($\mu d \xrightarrow{\uparrow} \mu h X$),

BELLE ($e^+ e^- \rightarrow h^+ h^- X$) data

M. Anselmino et al., Nucl. Phys. Proc. Suppl. 191 (2009) 98

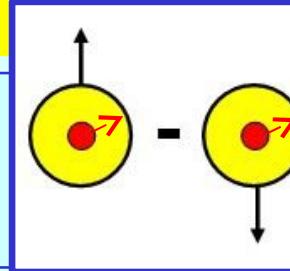


Sivers Amplitudes for Pions

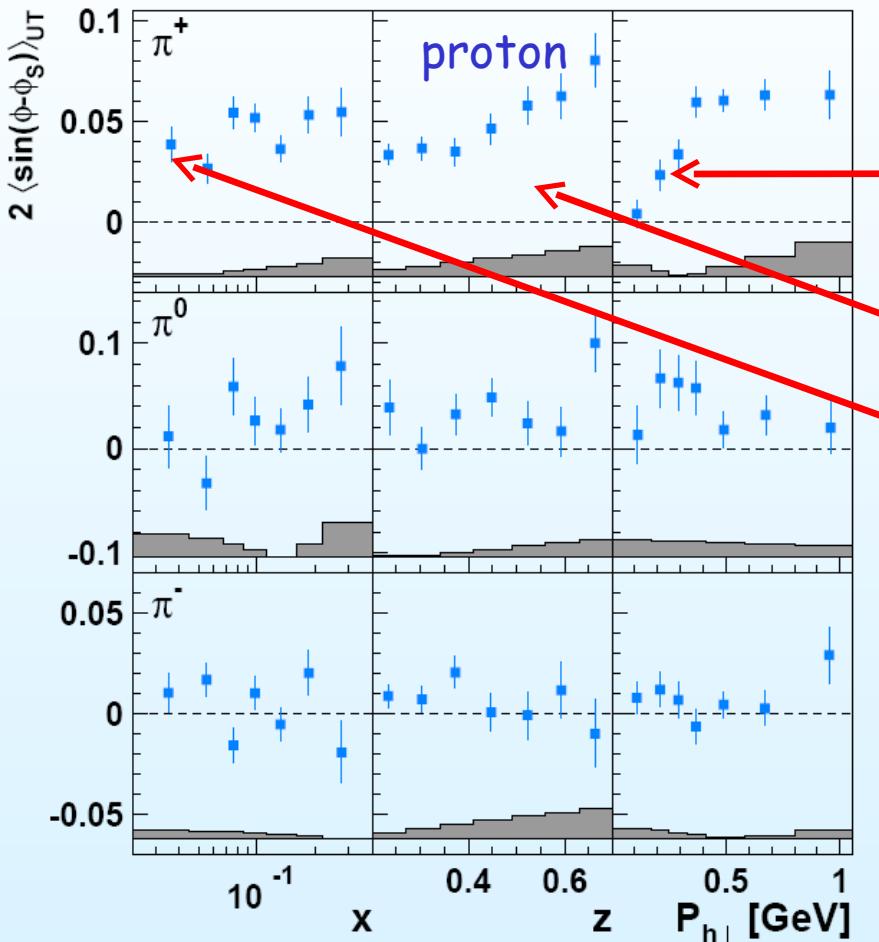
Sivers DF

$$2\langle \sin(\phi - \phi_s) \rangle \quad h_{UT} \propto f_{1T}^{\perp, q}(x) \otimes D_1^q(z)$$

PRL 103 (2009) 152002



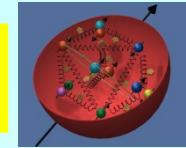
N/q	U	L	T
U	f_1		h_1^\perp
L		g_1^\perp	h_{1L}
T	f_{1T}	g_{1T}	h_1 h_{1T}



- First observation of non-zero Sivers DF in DIS
- Rise at low $P_{h\perp}$, plateau at high $P_{h\perp}$
- Clear rise with z
- Non-zero at low x

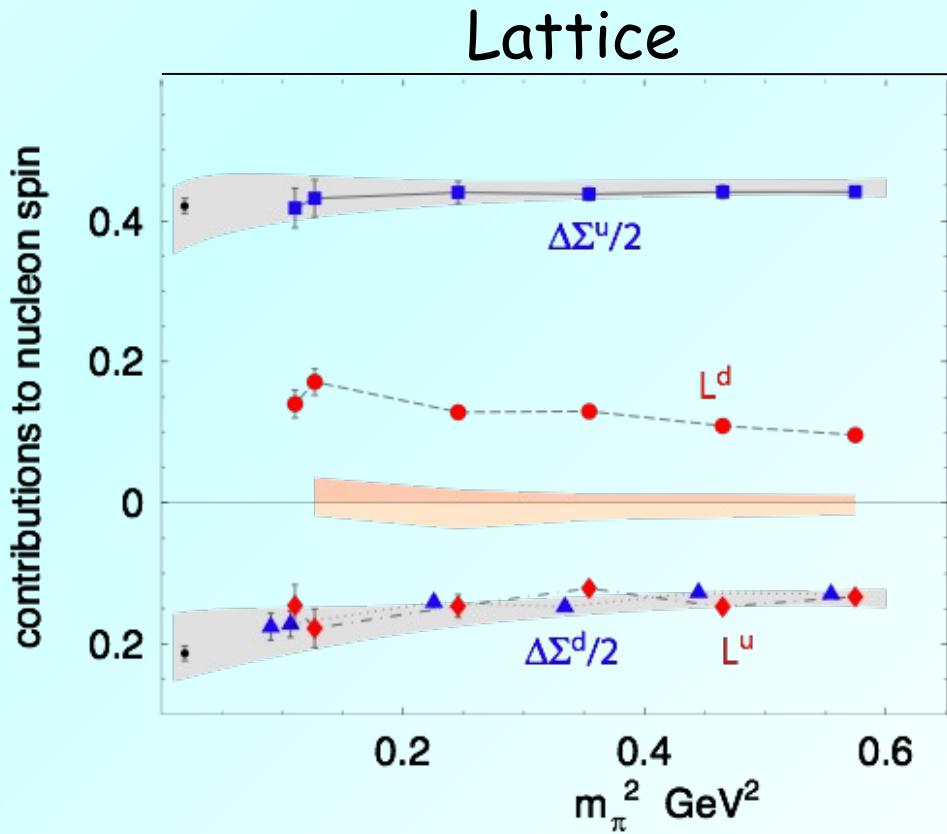
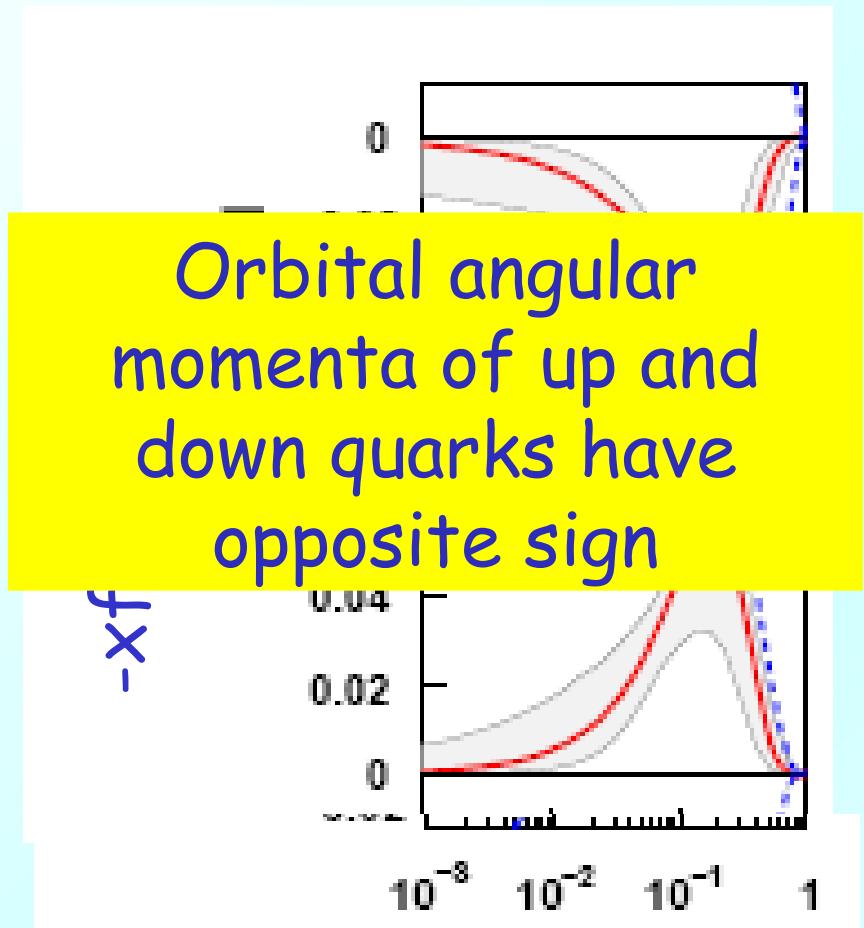
Experimental evidence for orbital angular momentum L_q of quarks
 But: Quantitative contribution of L_q to nucleon spin still unclear

Sivers distribution



Fit to HERMES ($e p \rightarrow e h X$) and COMPASS ($\mu d \rightarrow \mu h X$) data

M. Anselmino et al., Phys. Rev. D79 (2009) 054010



$$L_d \approx -L_u \approx 0.2$$

$$L_d + \Delta d/2 \approx 0 !??$$