Azimuthal correlations in photoproduction and deep inelastic *ep* scattering at HERA

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Subject of this analysis

The study of multibody QCD interactions under extreme conditions

Confinement in QCD "hides" the details of such interactions We can collide **heavy nuclei** to open them up in the laboratory



Central Question:

Is a similar kind of multibody environment created in much smaller *ep* systems produced at HERA?

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Outline

- Introduction to heavy-ion $(\mathbf{A} + \mathbf{A})$ collisions
 - A new form of matter expected
 - Pillars of heavy-ion measurements
- Surprising observation in $\boldsymbol{p} + \boldsymbol{p}$ collisions at the LHC
- ZEUS analysis in ep

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A new form of matter expected



Lattice QCD predicts a sharp rise of the energy density of matter around T=170 GeV. The transition is caused by an increase in the number of degrees of freedom. Hadrons \rightarrow deconfined quarks and gluons (**Quark-Gluon Plasma**.)

A new form of matter expected



- The shear magnitude of current accelerator beam energies and the miniscule size of a nucleus creates enormous energy densities and temperatures.
- Both the LHC and RHIC reach temperatures expected in the early Universe about 1 microsecond after the Big Bang.
- Observations in heavy-ion collisions can tell us about the early Universe.

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Pillars of heavy-ion measurements: Jet Quenching



- Most of the time, the hardest scatter will occur off center.
- This means that one parton will traverse a greater path length through the collision zone.

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Pillars of heavy-ion measurements: Jet Quenching



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- A dramatic suppression of back-to-back jets was found at RHIC (Jet Quenching).
- Interpreted as an indication of substantial final-state rescattering.

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- A dramatic suppression of back-to-back jets was found at RHIC (Jet Quenching).
- Interpreted as an indication of substantial final-state rescattering.
- Jet Quenching was confirmed at the LHC.

Pillars of heavy-ion measurements: Elliptic Anisotropy



• The spatial configuration of the initial scattering typically has a large elliptic component.

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Pillars of heavy-ion measurements: Elliptic Anisotropy



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• **Rescattering** between the produced partons converts this spatial eccentricity into an asymmetry in momentum space.

Pillars of heavy-ion measurements: Elliptic Anisotropy



- The spatial configuration of the **initial scattering** typically has a large elliptic component.
- **Rescattering** between the produced partons converts this spatial eccentricity into an asymmetry in momentum space.
- The elliptic asymmetry is clearly evident in two-particle correlations: collectivity.

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Evolution of a heavy-ion collision



- The initial scattering takes place over an extended region.
- A subsequent stage of rescattering produces a thermally equilibrated system.
- This fluid of QCD matter is called the **quark-gluon plasma (QGP)**.
- Relativistic hydrodynamics describes the evolution of the QGP.

QED fluid and QCD fluid

QED fluid



QCD fluid



Girolamo Sferrazaa Papa

QGP in heavy-ion collisions lives only a few yoctoseconds (10^{-24})

Surprising observation in $oldsymbol{p}+oldsymbol{p}$ collisions at the LHC

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Double ridge in two-particle correlations



- Two particle correlations in heavy-ion collisions show a clear **double ridge**, which is interpreted as a sign of fluid-like behaviour in QCD—The QGP.
- C(Δη, Δφ) = S(Δη, Δφ)/B(Δη, Δφ),
 S and B are formed from pairs from the same- and mixed-events, respectively.

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Double ridge in two-particle correlations



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- C(Δη, Δφ) = S(Δη, Δφ)/B(Δη, Δφ),
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- The start of the LHC revealed that high-multiplicity p + p collisions also have a double-ridge!
- Such collisions were thought to be too small to produce a thermally equilibrated QGP.

Motivation for the analysis



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- C(Δη, Δφ) = S(Δη, Δφ)/B(Δη, Δφ),
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- The start of the LHC revealed that high-multiplicity p + p collisions also have a double-ridge!
- Such collisions were thought to be too small to produce a thermally equilibrated QGP.
- What about even more fundamental *ep* scattering at HERA??

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How small can a colliding system be while still exhibiting the collective behaviour typically associated with the quark–gluon plasma observed in heavy-ion collisions?

What kind of environment could collectivity evolve from?

I'll present recently posted measurements of azimuthal correlations in neutral current DIS and photoproduction with the ZEUS detector: arXiv:2106.12377 (Submitted to JHEP)

The HERA collider and experiments



- Location: DESY, Hamburg, Germany
- Data taking: 1992 2007
- 27.5 GeV electrons/positrons 920 GeV protons $\rightarrow \sqrt{s} = 318 \text{ GeV}$
- HERA I+II: 500 pb^{-1} per experiment

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Deep inelastic scattering (DIS)



- DIS is defined by large virtualities: $Q^2 \gg \Lambda_{\rm QCD}^2. \label{eq:Q2}$
- Transverse radius (*R*_t) and longitudinal length (*L*) of the probed region are given by:

$$egin{aligned} R_t &\sim rac{1}{Q} \ L &\sim rac{1}{m_{ ext{proton}\, imes}} \end{aligned}$$

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• Neutral current (NC) DIS involves the exchange of photon or Z boson.

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Photoproduction (PhP)



- Photoproduction (γp) is defined by small virtualities: Q² ≪ Λ²_{OCD}.
- Exchanged photon may fluctuate into quarks and gluons.
- Larger interaction regions are probed.
- Multiparton Interactions are possible.

• Scattering is hadron-like.

Multiparton Interactions (MPI)



- MPI occur when there's more than one 2 → 2 partonic scattering between the beam particles in a given event.
- If the scatterings are sufficiently hard $(p_T \gtrsim 1 \text{ GeV})$, they can be modeled in an event generator like PYTHIA 8.
- Established feature in high-multiplicity hadronic collisions. Not conclusively observed in *ep* scattering so far.

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Illustration of MPI growth



- Rough illustration of how nMPI grows from DIS to heavy-ions
- *N*_{coll}: number of binary nucleon-nucleon collsions
- N_{nn}^{partonic} : number of parton scatterings per binary nucleon-nucleon collision
- Estimates for N_{coll} taken from

 Ann. Rev. Nucl. Part. Sci. 57, 205 (2007)
 PRC 97 024905 (2018).
- Estimates for N_{nn}^{partonic} taken from PYTHIA 8

ep photoproduction: subsequent rescattering phase possible



- The initial scattering is shown here with 3 MPIs (black dots)
- Unlike in DIS, the spatial extent of this "initial state" is finite with an irregular shape, in general.
- Subsequently, a phase of **rescattering may occur**, whereby a local thermal equilibrium might form.

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Observables for this analysis

To search for collective behavoiur and MPI in ep scattering, we measured the following:

- Charged hadron multiplicity distribution: $dN/dN_{\rm ch}$
- Transverse momentum distribution: dN/dp_{T}
- Pseudorapidity distribution: $dN/d\eta$
- Two-particle azimuthal correlation functions: $c_n\{2\}$ and $\mathcal{C}(\Delta\eta,\Delta\varphi)$
- Four-particle cumulant azimuthal correlations: c_n {4}

Additional detailed measurements in DIS alone can be found in our previous paper JHEP 04 (2020) 070.

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Two- and four-particle correlation functions

Two-particle azimuthal correlations are measured:

 $c_n\{2\} = \langle \langle \cos n(\phi_i - \phi_j) \rangle \rangle$

 φ_i is the azimuthal angle of particle i

n is the harmonic (n=1, 2 studied here) Four-particle cumulant correlations are also measured:

$$C_n\{4\} = \langle (\cos n(\phi_i + \phi_j - \phi_k - \phi_l)) \rangle c_n\{4\}(p_{T,1}) = C_n\{4\}(p_{T,1}) - 2 c_n\{2\}(p_{T,1}) c_n\{2\}$$

where $p_{T,1}$ is the transverse momentum of particle i

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ZEUS detector



Charged particles are tracked in the central tracking detector (CTD) and micro vertex detector (MVD) in a 1.43 T magnetic field.

Depleted uranium calorimeters.

The barrel and rear ones are used to help identify the scattered electron.

A fully contained event is characterized by $\sum_{i} (E_i - P_{z,i}) = 55 \text{ GeV}$ due to energy and momentum conservation.

Track selection



Track selection for correlation analysis

- Reject scattered electron (if detected)
- $-1.5 < \eta < 2.0$
- $0.1 < p_T < 5.0 \text{ GeV}$
- ${\small \bullet}~\geq 1$ MVD hit
- DCA_{XY,Z} < 2 cm
- $\Delta R > 0.4$ (cone around electron)

$$N_{
m ch} = \sum_{i}^{N_{
m rec}} w_i^{(1)}$$

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DIS event selection



DIS Event selection (0.2 M)

- $N_{
 m ch} \geq 20$
- DIS triggers
- electron probability > 90%

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$$Q^2 = -(k - k')^2 > 5 \text{ GeV}^2$$

- $k'_0 > 10 \text{ GeV}$
- *r* > 15 cm
- $\theta_e > 1 \text{ rad}$
- $47 < \sum (E_i P_{z,i}) < 69 \text{ GeV}$
- $|V_z| < 30 \text{ cm}$

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Photoproduction event selection



Photoproduction event selection (5 M)

- $N_{
 m ch} \geq 20$
- PhP oriented triggers
- electron probability < 90%
- $k_0' < 15 \text{ GeV}$
- $\sum (E_i P_{z,i}) < 55 \text{ GeV}$
- $|V_z| < 30 \text{ cm}$

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Results: Ridge plots



A near-side peak and away-side ridge are clearly visible. Photoproduction correlations are dimished wrt those in DIS. **No visible double-ridge.**

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Results: Q^2 evolution of $c_1{2}$



Photoproduction correlation strengths ($Q^2 = 0$) are clearly diminished wrt those in DIS.

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Results: Q^2 evolution of $c_2\{2\}$



Photoproduction correlation strengths ($Q^2 = 0$) are clearly diminished wrt those in DIS.

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Multiparton Interactions

Study of MPI in *ep* photoproduction

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Results: $dN/dN_{\rm ch}$



The level of MPI and IR divergencies are controled by the p_{T0} parameter in PYTHIA 8.

It is used to regularize the interaction cross section in PYTHIA 8.

$$rac{d\sigma}{dp_{
m T}^2} \propto rac{lpha_s^2(p_{
m T0}^2+p_{
m T}^2)}{(p_{
m T0}^2+p_{
m T}^2)^2}$$

The energy dependence of this parameter is given by $p_{\rm T0} = p_{\rm T0}^{\rm ref} (W/7 \,{\rm TeV})^{0.215}$, where W is the $\gamma p \sqrt{s}$.

More MPI \rightarrow lower p_{T0}^{ref}

Colour Reconnection (CR) is PYTHIA's modeling of rescattering between partons from different MPIs

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Results: $dN/dp_{ m T}$ and $dN/d\eta$



• The scenarios of no MPI and very many MPI are disfavored.

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Results: $c_1\{2\}$ and $c_2\{2\}$ versus $\Delta\eta$



- Correlation strengths are diluted by MPI.
- The scenarios of no MPI and very many MPI are disfavored.

Results: c_1 {2} and c_2 {2} versus $\langle p_{\mathrm{T}} angle$



• $c_1\{2\}$ versus $\langle p_T \rangle$ not sensitive to MPI and not described well by PYTHIA 8.

• More extreme levels of MPI are favored by $c_2\{2\}$ versus $\langle p_T \rangle$.

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Results: c_1 {4} and c_2 {4} versus $p_{T,1}$



- Four-particle cumulant is positive, which is in contrast to the negative values seen in non-central heavy-ion collisions.
- The scenarios of no MPI and very many MPI are disfavored.

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Summary of measurements

- Measurements of charged-particle azimuthal correlations at high multiplicity have been presented using ZEUS data in *ep* photoproduction (γp) and NC DIS at $\sqrt{s} = 318$ GeV.
- $dN/dN_{\rm ch}$, $dN/dp_{\rm T}$, and $dN/d\eta$ were measured in γp .
- There is no clear indication of a double ridge from $C(\Delta \eta, \Delta \varphi)$ in either γp or DIS.
- Two-particle correlation strengths are markedly diminished in γp wrt DIS.
- While there is no consistent value of p_{T0}^{ref} in PYTHIA 8 that describes all the γp data, the "no MPI" scenario is strongly disfavored.
- Other parameters in PYTHIA 8 such as those related parton showering should be investigated to improve the description of the data.

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Responses to the motivating questions

DIS probes very small length scales in the proton ($\ll 1$ fm). Photoproduction probes much larger scales, which can be as large as the proton itself.

The observations in both regimes **do not** reveal significant collective behaviour like that seen in heavy-ions or high-multiplicity hadronic collisions.

The concept of multiparton interactions provides a useful tool to help understand the emergence of collective behaviour.

It sets the stage for a potential rescattering phase.

	nMPI	Collectivity
ep photoproduction	~ 3	No
pp high-multiplicity	~ 20	Yes

The initial states in both systems may be similar in spatial extent but completely different in the number of MPI.

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Tracking efficiency corrections

The efficiency correction weights for 1-, 2-, and 4-particle distributions are defined as:

$$w^{(n)} = \frac{N_{gen}^n(\vec{x})}{N_{rec}^n(\vec{x})}$$

They are computed differentially in Monte Carlo simulations of the ZEUS detector:

dimension of \vec{x}	One-particle $(n=1)$	Two-particle $(n=2)$	Four-particle (n=4)
	φ	$\varphi_1 - \varphi_2$	$\varphi_1+\varphi_2-\varphi_3-\varphi_4$
x_2	η	$\langle \eta_i - \langle \eta angle angle$	$\langle \eta_i - \langle \eta \rangle angle$
x_3	p_{T}	$\langle p_{T,i} - \langle p_T \rangle \rangle$	$\langle p_{T,i} - \langle p_T \rangle \rangle$
x_4 (charge)	q	$ q_1 + q_2 $	$ q_1 + q_2 + q_3 + q_4 /2$
x_5	-	$N_{ m rec}$	$N_{ m rec}$

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Condensed view of PYTHIA 8 comparisons



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W distribution



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Direct and Resolved event distributions



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nMPI in high-multiplicity p + p PYTHIA 8 at LHC energies



PYTHIA 8 p + p events at $\sqrt{s} = 13$ TeV were generated.

 $\it N_{ch}$ was counted according to the ATLAS acceptance used in PRL 116 172301. $-2.5 < \eta < 2.5, \, 0.4 < \it p_T < 50~GeV$

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Results: Q^2 evolution of $c_1\{2\}$



 $\gamma \pmb{p}$ direct: photon couples directly to a quark in the proton.

 γp resolved: photon splits into quarks and gluons which then scatter with the proton (MPI possible).

Photoproduction correlation strengths ($Q^2 = 0$) are clearly diminished wrt those in DIS.

The LEPTO model of DIS gives a rough qualitative description of the data.

PYTHIA 8 with only the direct component of γp predicts much stronger correlations than the full calculation (direct + resolved).

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Results: Q^2 evolution of $c_2\{2\}$



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