

Testing dark energy models with atom interferometry

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University of Nottingham

Outline:

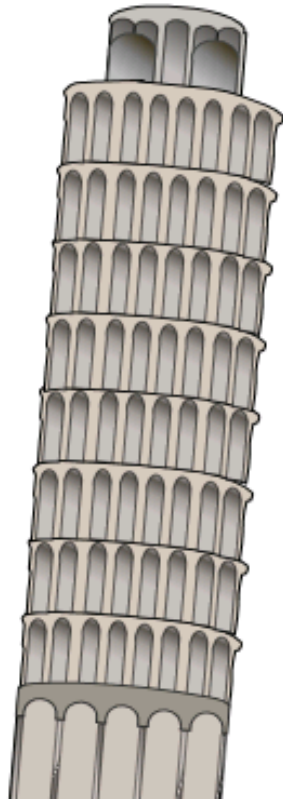
Dark energy and screened fifth forces

How to search for screening

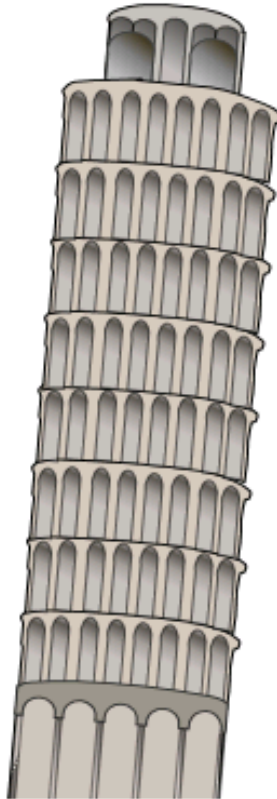
Atom interferometry constraints

Testing the Equivalence Principle

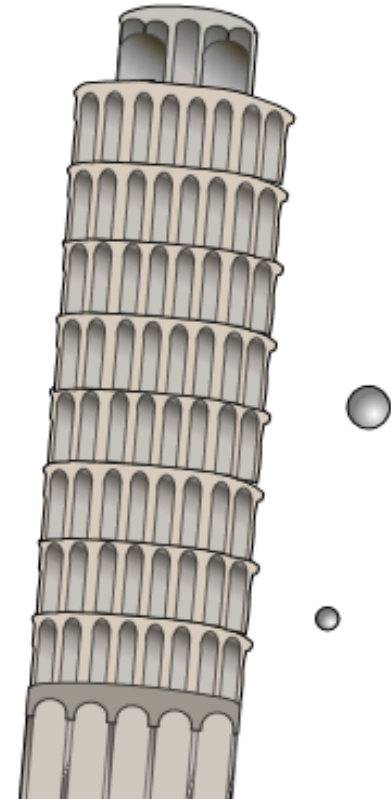
Do large objects and small objects fall at the same rate?



Old idea

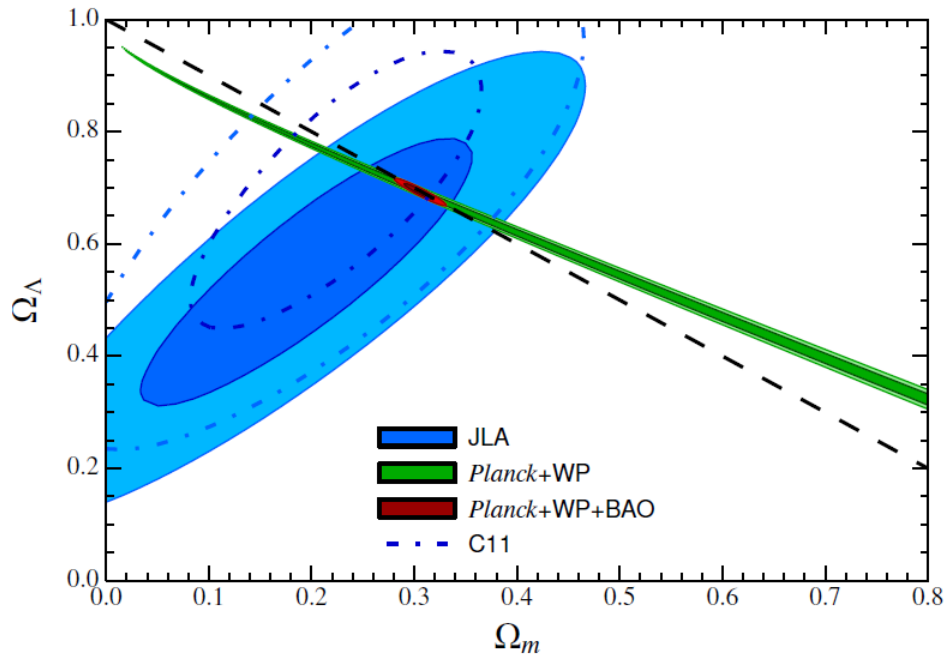


Galileo



Dark Energy?

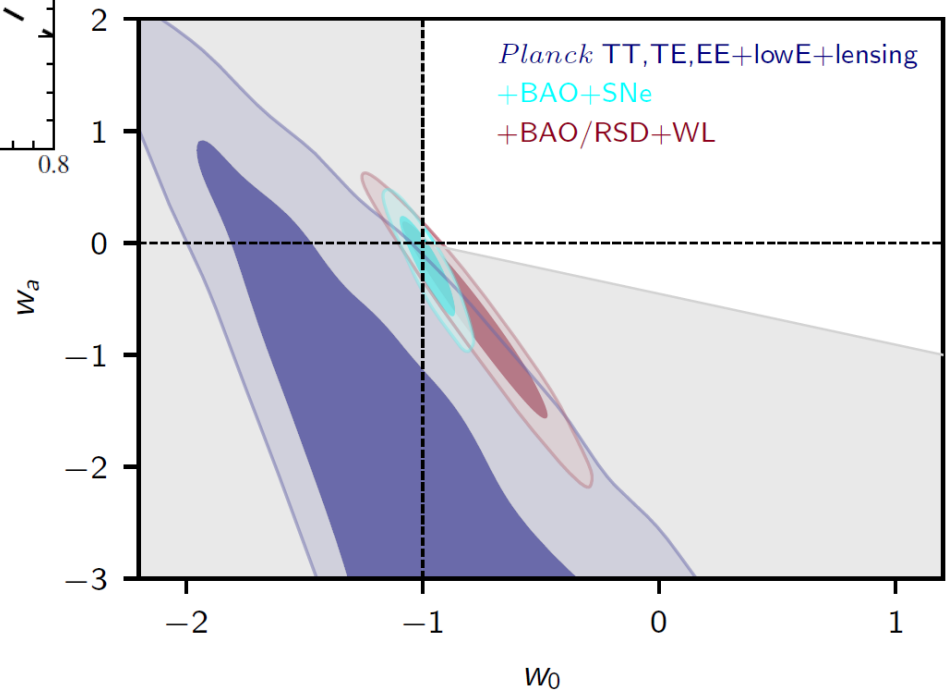
Dark Energy Today



$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4}T_{\mu\nu}$$

Equation of state:

$$w = w_0 + (1 - a)w_a$$



Why Introduce Light Scalar Fields?

A new type of matter eg dark energy

- Quintessence directly introduces new fields
- New, light (fundamental or emergent) scalars

A modification of gravity

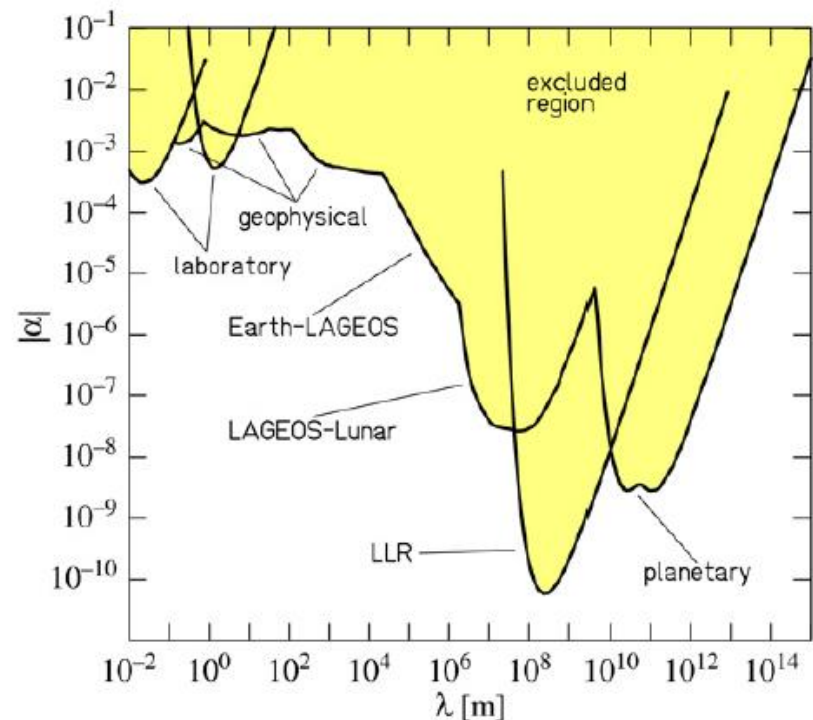
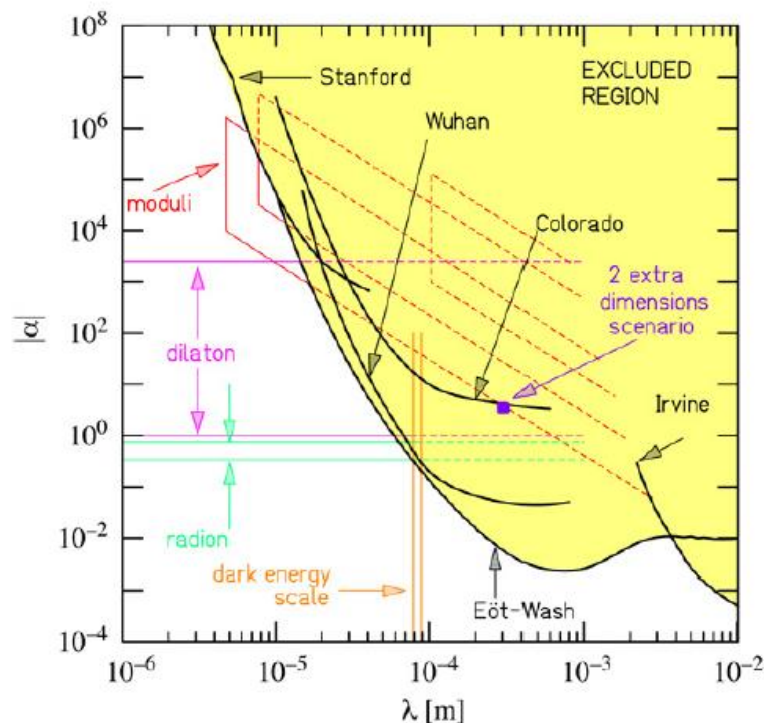
- General Relativity is the unique interacting theory of a Lorentz invariant, massless, helicity-2 particle
Papapetrou (1948). Weinberg (1965).
- New physics in the gravitational sector will introduce new degrees of freedom, typically Lorentz scalars

Dark matter could also be a light scalar

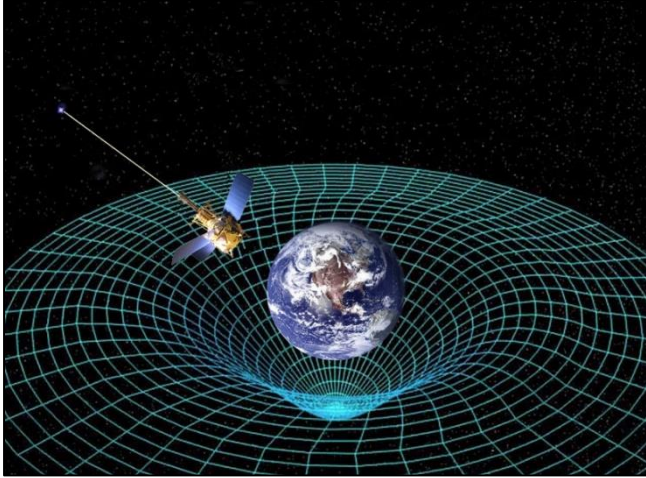
Yukawa Fifth Forces

A long range Yukawa fifth force is excluded to a high degree of precision in the solar system

$$V(r) = -\frac{G\alpha m_1 m_2}{r} e^{-m_\phi r}$$

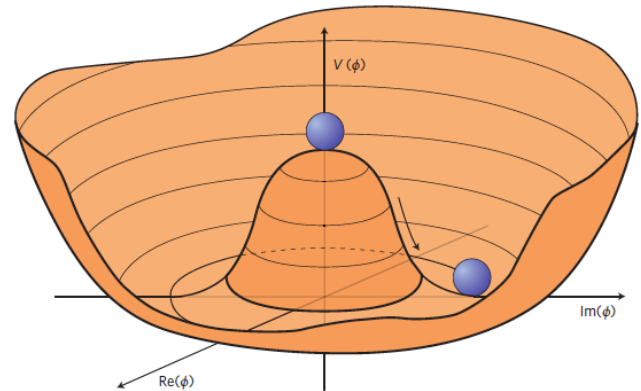


Is the New Physics Linear?



General relativity is
a non-linear theory

Higgs scalar has a
non-linear potential



New Physics is Non-linear: Screening Mechanisms

- **Locally weak coupling**

Symmetron and varying dilaton models

Pietroni (2005). Olive, Pospelov (2008). Hinterbichler, Khoury (2010). Brax et al. (2011).

- **Locally large mass**

Chameleon models

Khoury, Weltman (2004).

- **Locally large kinetic coefficient**

Vainshtein mechanism, Galileon and k-mouflage
models

Vainshtein (1972). Nicolis, Rattazzi, Trincherini (2008).

Babichev, Deffayet, Ziour (2009).

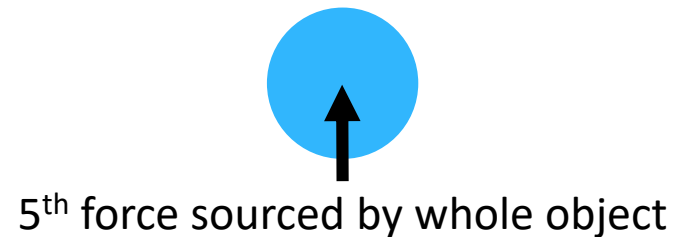
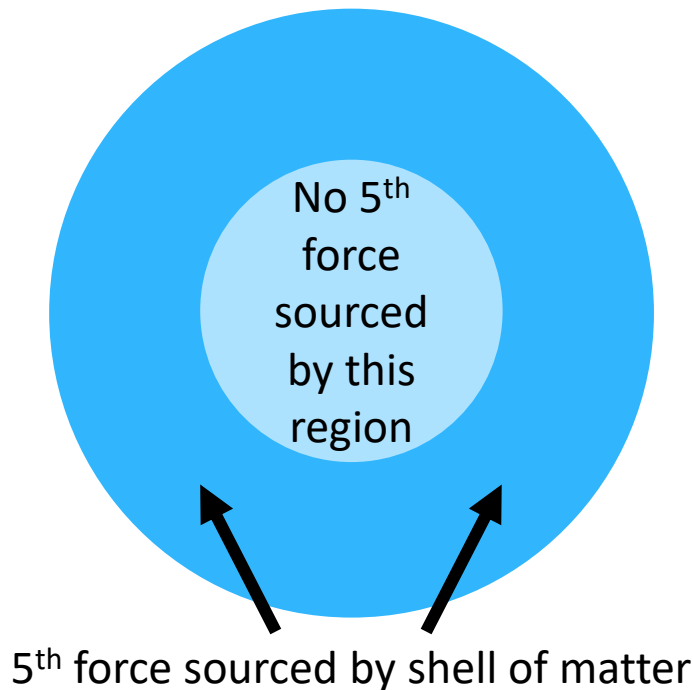
Screening Phenomenology

Change the dependence on distance

- Vainshtein screening

Change the way in which matter sources the scalar field

- thin-shell effect



Large objects are screened
from the 5th force

The Chameleon



A scalar field with canonical kinetic terms, non-linear potential, and direct coupling to matter

$$S_\phi = \int d^4x \sqrt{-g} \left(-\frac{1}{2}(\partial\phi)^2 - V(\phi) - A(\phi)\rho_m \right)$$

$$V(\phi) = \frac{\Lambda^5}{\phi}, \quad A(\phi) = \frac{\phi}{M} ,$$

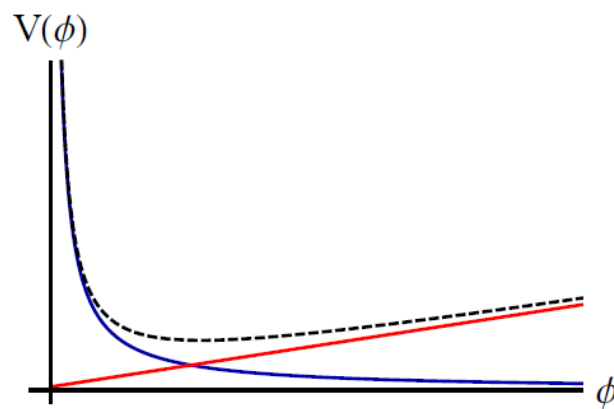
Khoury, Weltman. (2004). Image credit: Nanosanchez
Equivalent description as Higgs portal model:
CB, Copeland, Millington, Spannowsky. (2018)

Varying Mass

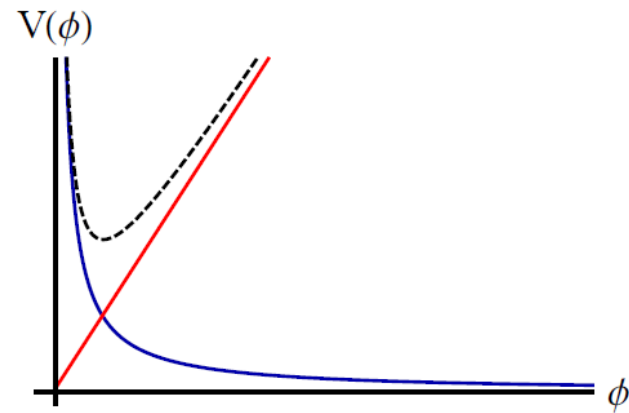
Dynamics governed by an effective potential

$$V_{\text{eff}} = \frac{\Lambda^5}{\phi} + \frac{\phi}{M} \rho$$

Non-linearities in the potential mean that the mass of the field depends on the local energy density



Low density

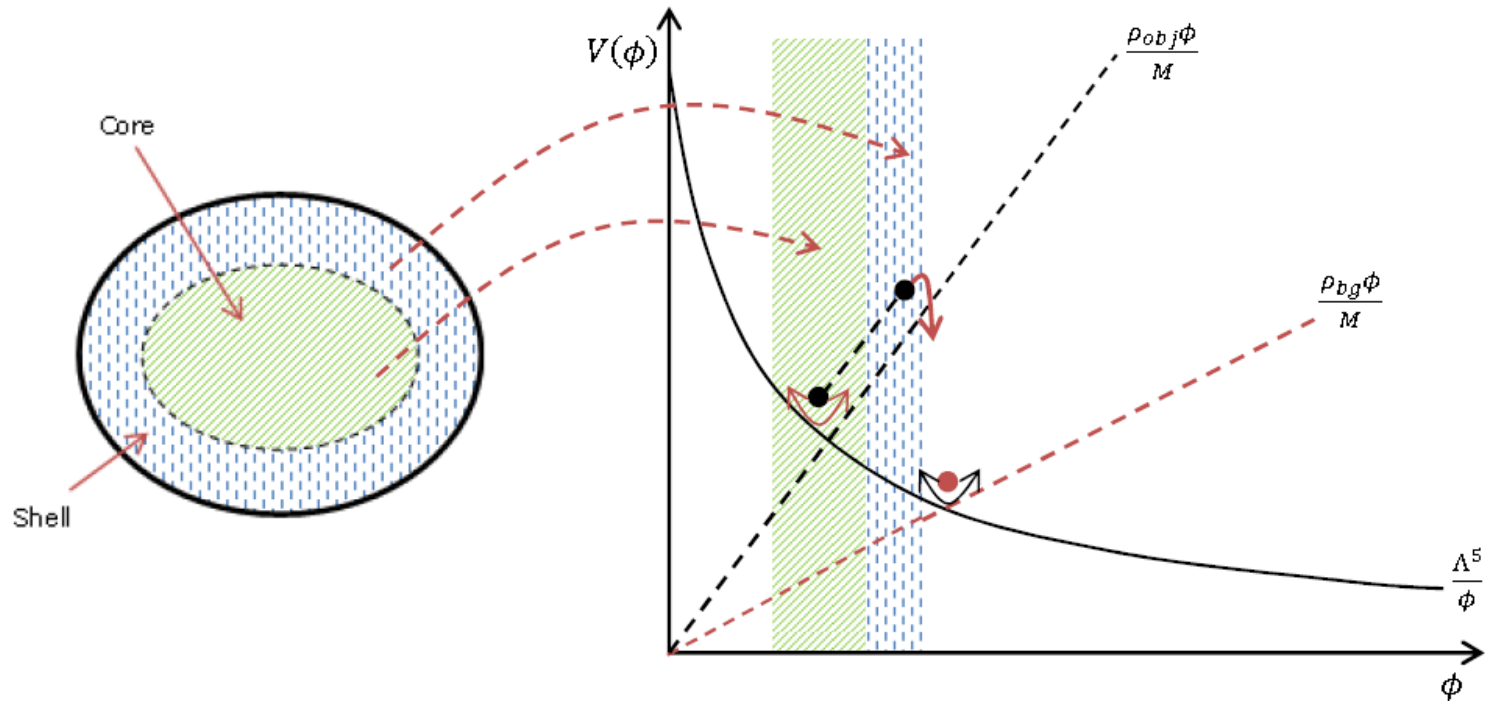


High density

Equivalent description as Higgs portal model:
CB, Copeland, Millington, Spannowsky. (2018)

Chameleon Screening

The increased mass makes it hard for the chameleon field to adjust its value



The chameleon potential well around 'large' objects is shallower than for canonical light scalar fields

The Scalar Potential

Around a static, spherically symmetric source of constant density

$$\phi = \phi_{\text{bg}} - \lambda_A \frac{1}{4\pi R_A} \frac{M_A}{M} \frac{R_A}{r} e^{-m_{\text{bg}} r}$$

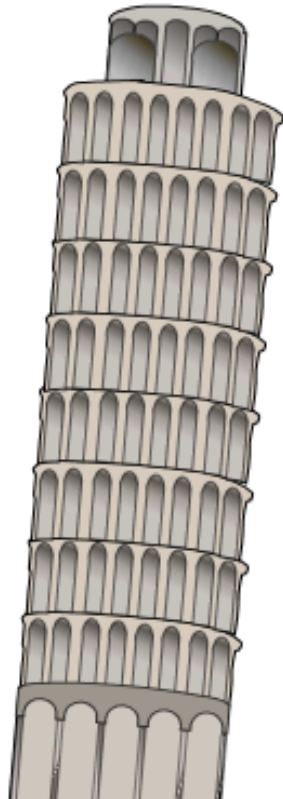
$$\lambda_A = \begin{cases} 1, & \rho_A R_A^2 < 3M\phi_{\text{bg}} \\ 1 - \frac{S^3}{R_A^3} \approx 4\pi R_A \frac{M}{M_A} \phi_{\text{bg}}, & \rho_A R_A^2 > 3M\phi_{\text{bg}} \end{cases}$$

This determines how ‘screened’ an object is from the chameleon field

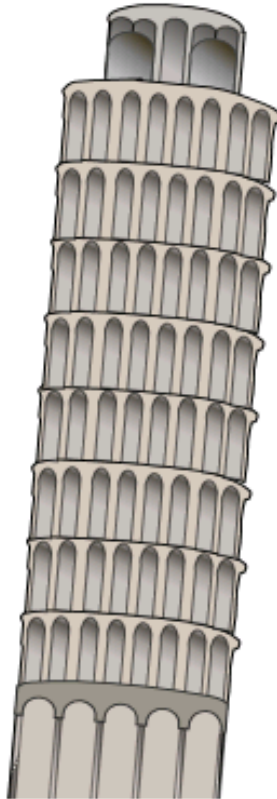
Ideal experiments use unscreened test masses e.g. atomic nuclei, neutrons, microspheres

Testing the Equivalence Principle

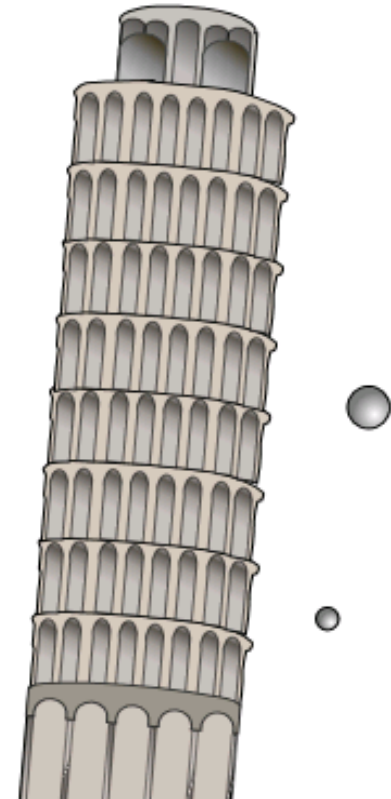
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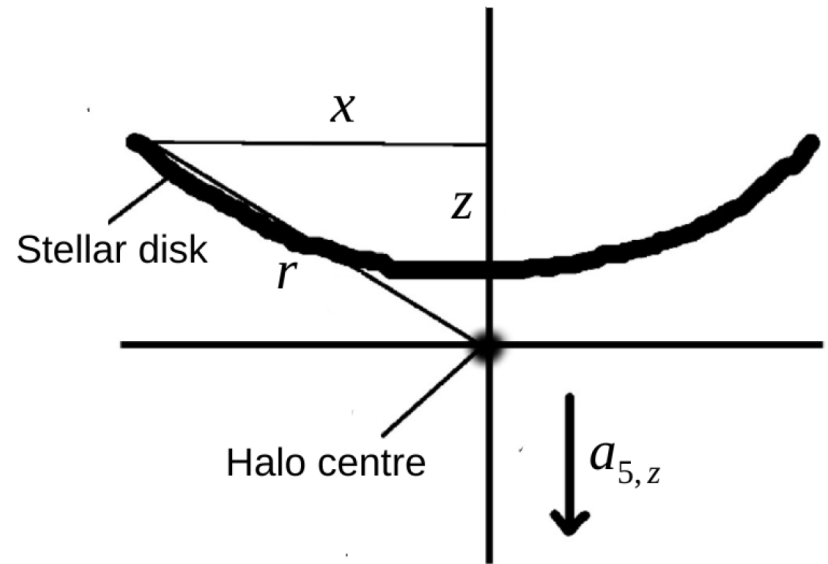
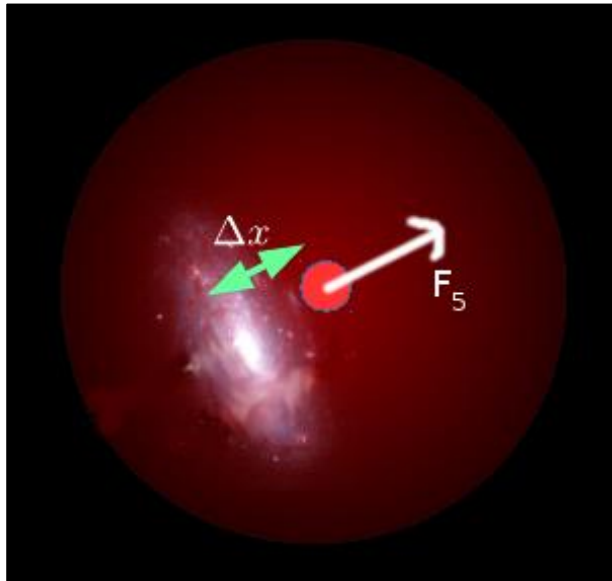


Dark Energy?

Tests on Galactic Scales

Different components of a dwarf galaxy may fall in a gravitational field at different rates

- Stars are screened, gas and dark matter are not
- Look for gas-star offsets & warping of galactic discs



Hui, Nicolis, Stubbs. (2009). Jain, VanderPlas. (2011)

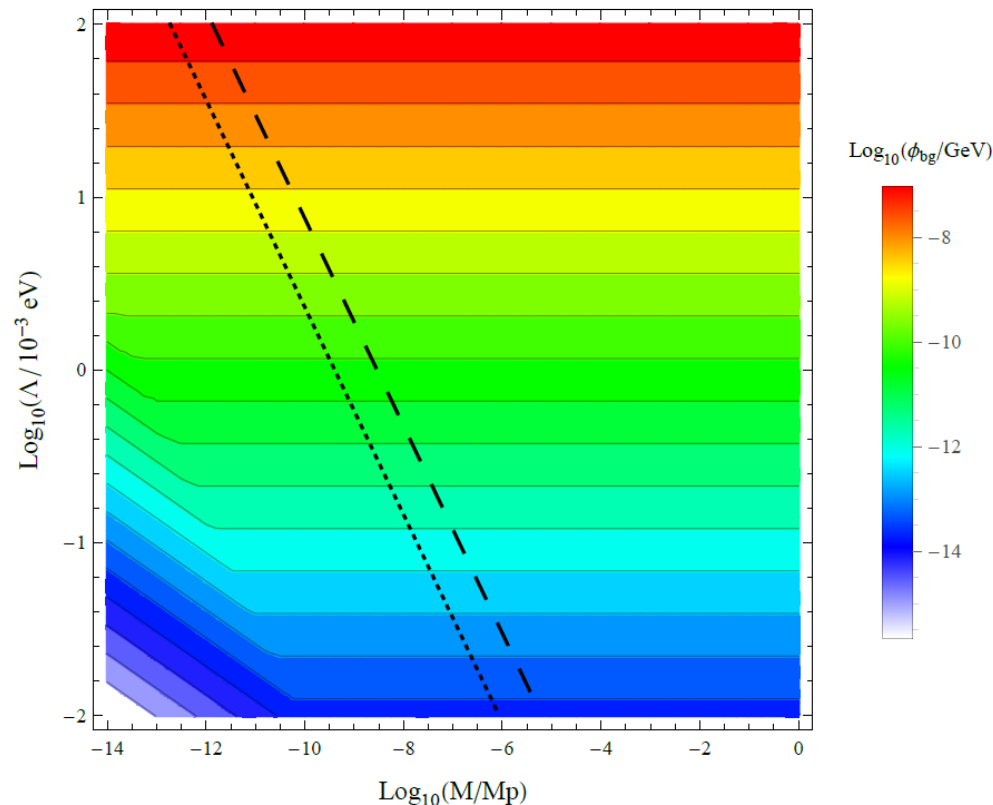
Desmond, Ferreira, Lavaux, Jasche. (2018)

Desmond, Ferreira. (2020)

Why Atom Interferometry?

In a spherical vacuum chamber, radius 10 cm, pressure 10^{-10} Torr

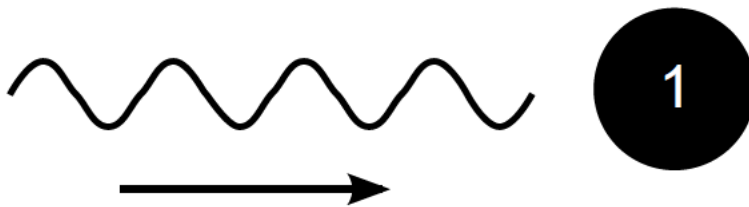
Atoms are unscreened above black lines
(dashed = caesium, dotted = lithium)



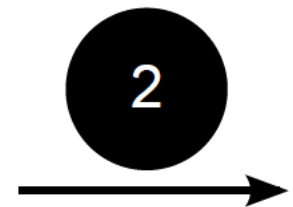
Atom Interferometry

An interferometer where the wave is made of atoms

Atoms can be moved around by absorption of laser photons

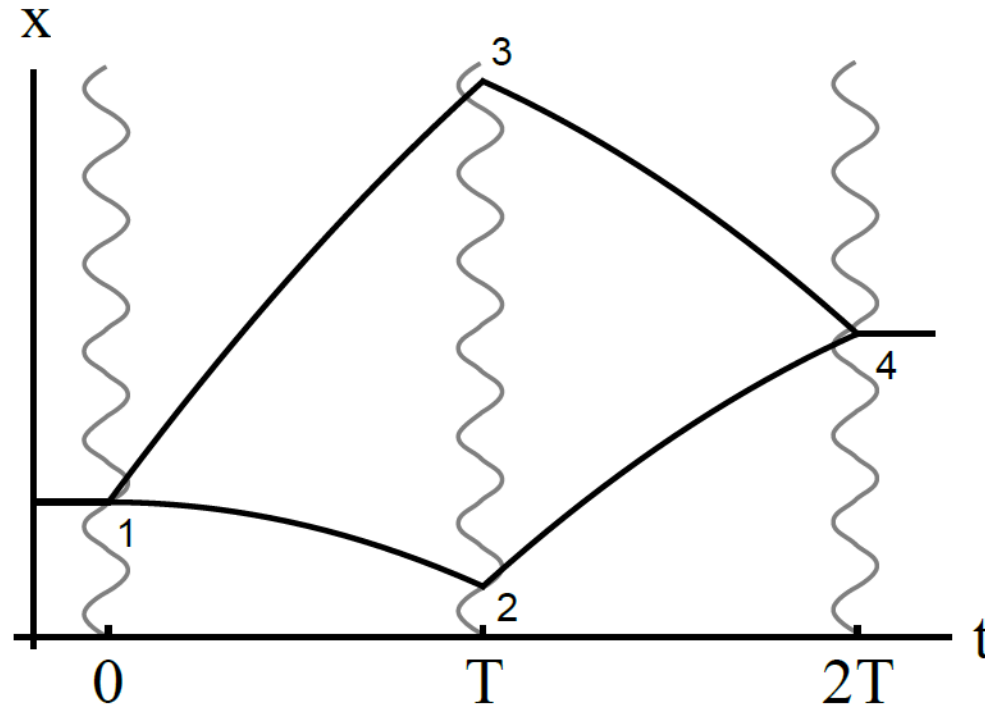


Photon Momentum = k
Atom in ground state



Atom in excited state
with velocity = V

Atom Interferometry



Probability measured in excited state at output

$$P = \cos^2 \left(\frac{kaT^2}{2} \right)$$

The Atomic Wavefunction

The probability of measuring atoms in the unexcited state at the output of the interferometer is a function of the wave function phase difference along the two paths

$$P \propto \cos^2 \left(\frac{\varphi_1 - \varphi_2}{2} \right)$$

For freely falling atoms the contribution of each path has a phase proportional to the classical action

$$\theta[x(t)] = C e^{(i/\hbar)S[x(t)]}$$

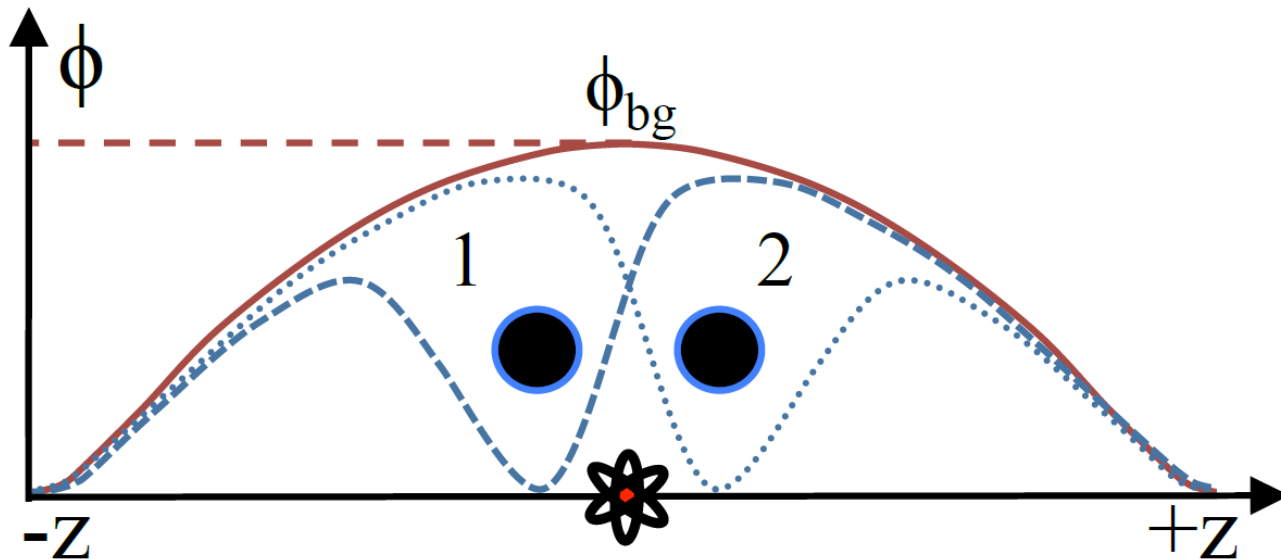
Additional contributions from interactions with photons, proportional to

$$(i/\hbar)(\omega t - \vec{k} \cdot \vec{x})$$

Atom Interferometry for Chameleons

The walls of the vacuum chamber screen out any external chameleon forces

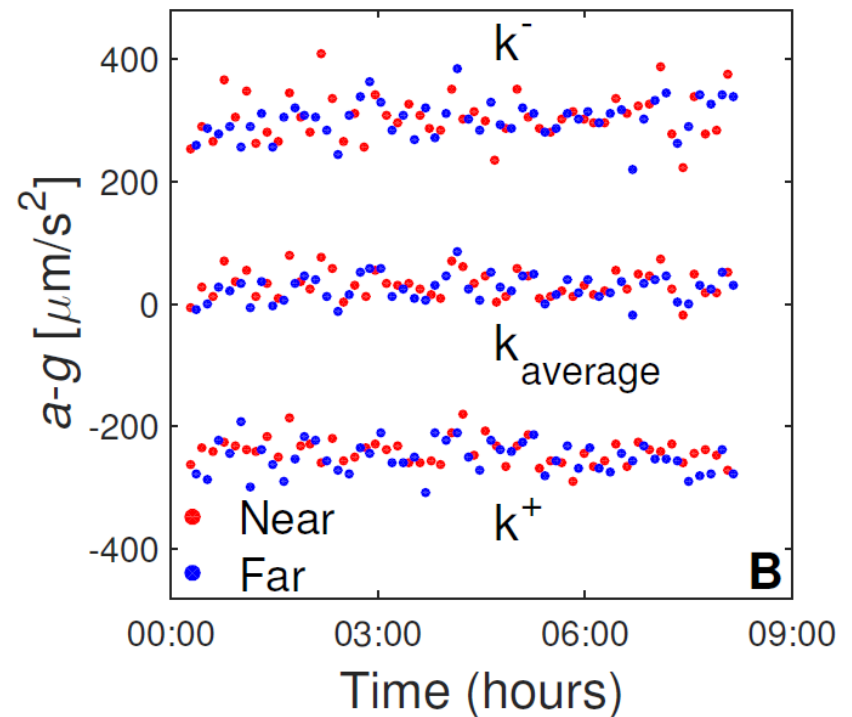
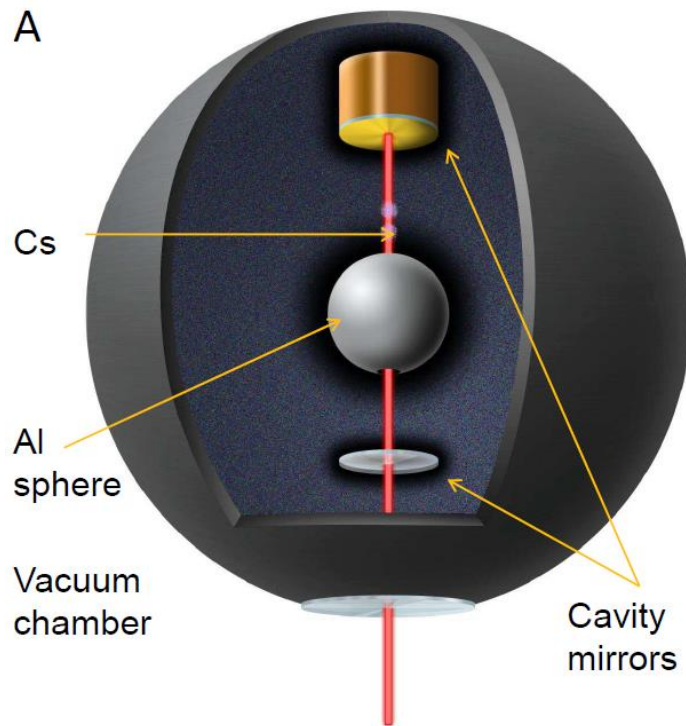
Macroscopic spherical mass, produces chameleon potential felt by cloud of atoms



Berkley Experiment

Using an existing set up with an optical cavity, looking for a signal on top of the Earth's magnetic field

Anomalous acceleration = $11 \pm 24 \text{ nm s}^{-2}$



Jaffe, Haslinger, Xu, Hamilton, Upadhye, Elder, Khoury, Müller. (2017)
Elder, Khoury, Haslinger, Jaffe, Müller, Hamilton. (2016)

Imperial Experiment

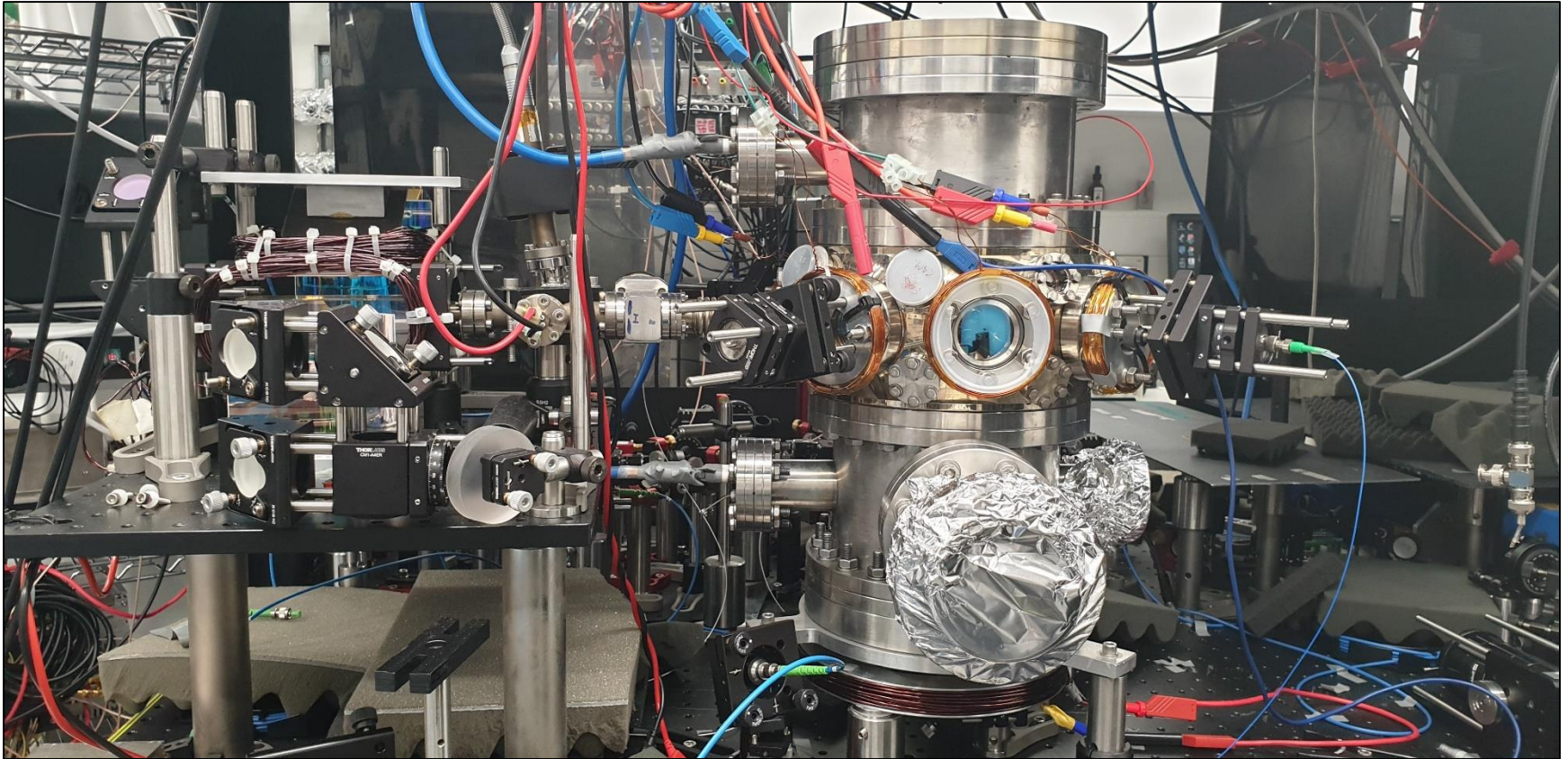
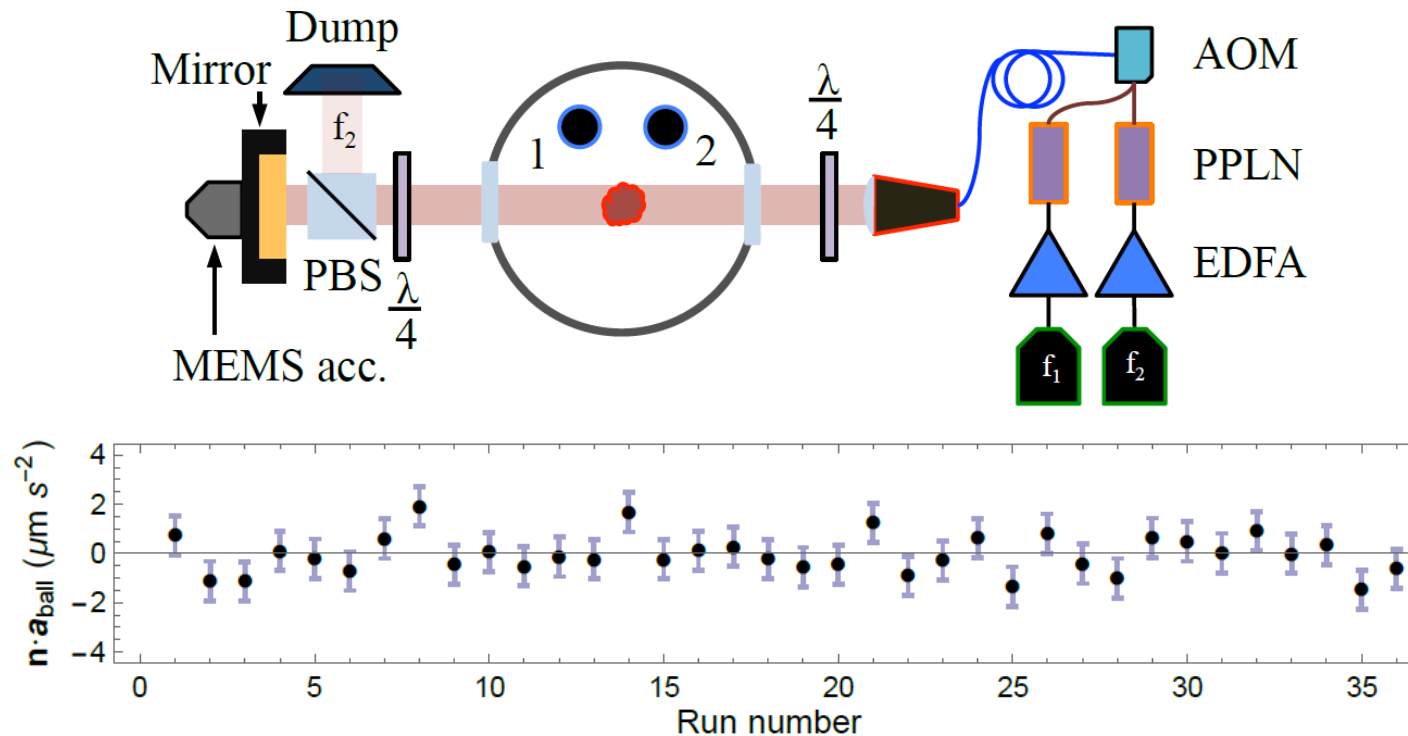


Image credit: E. Hinds

Imperial Experiment

Dedicated chameleon experiment, insensitive to the Earth's gravitational field

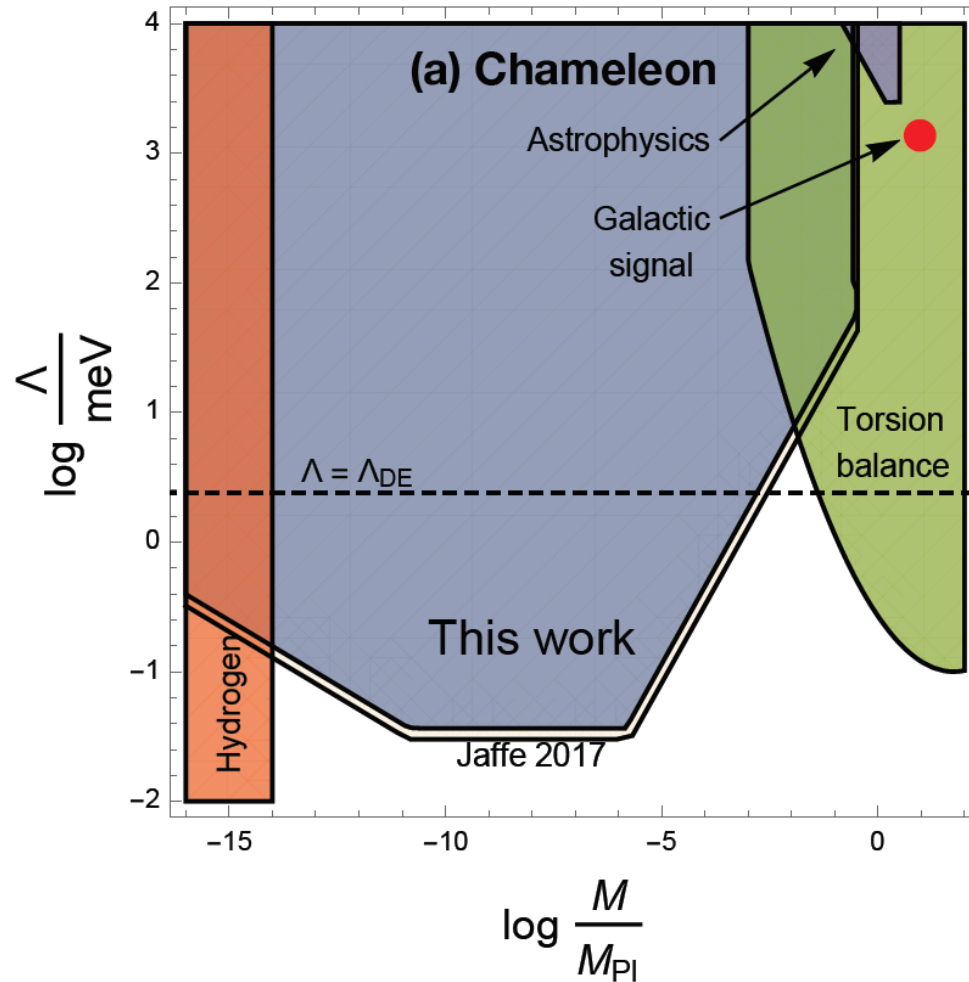
Anomalous acceleration = $-77 \pm 201 \text{ nm s}^{-2}$



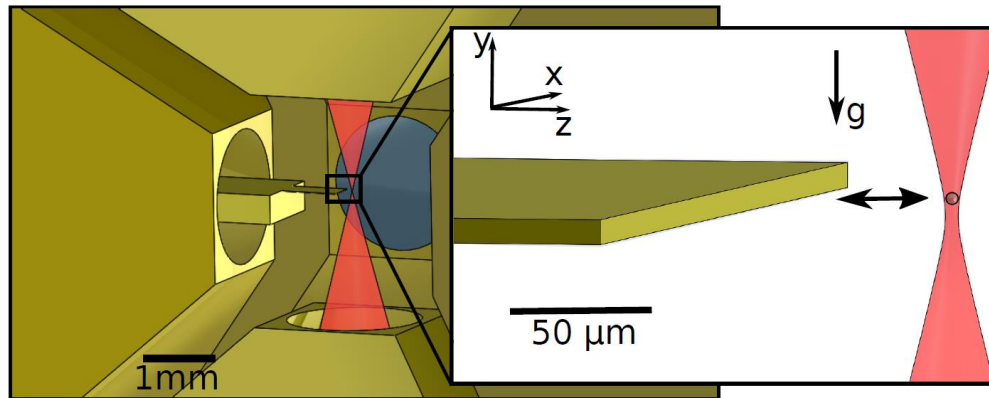
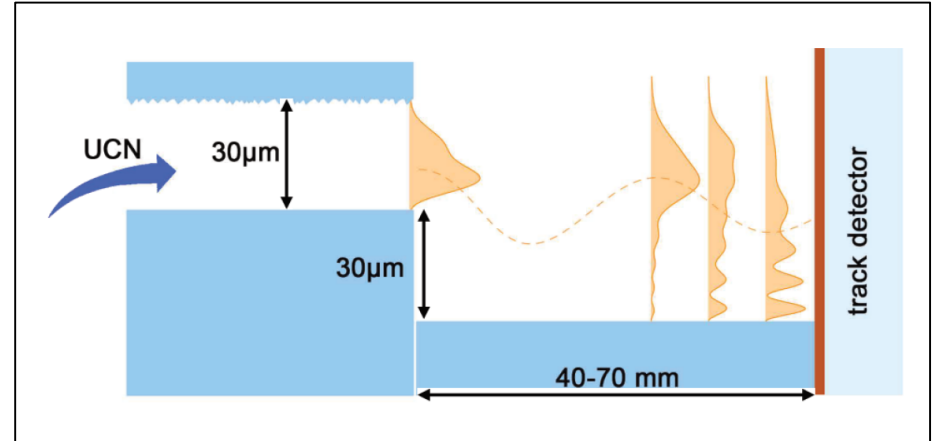
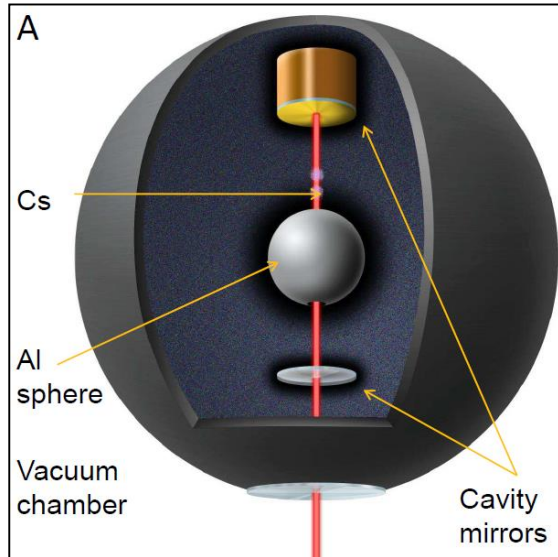
Sabulsky, Dutta, Hinds, Elder, CB, Copeland. arXiv:1812.08244

See also: Jaffe et al. (2017)

Imperial Experiment



Laboratory Searches – EP violation



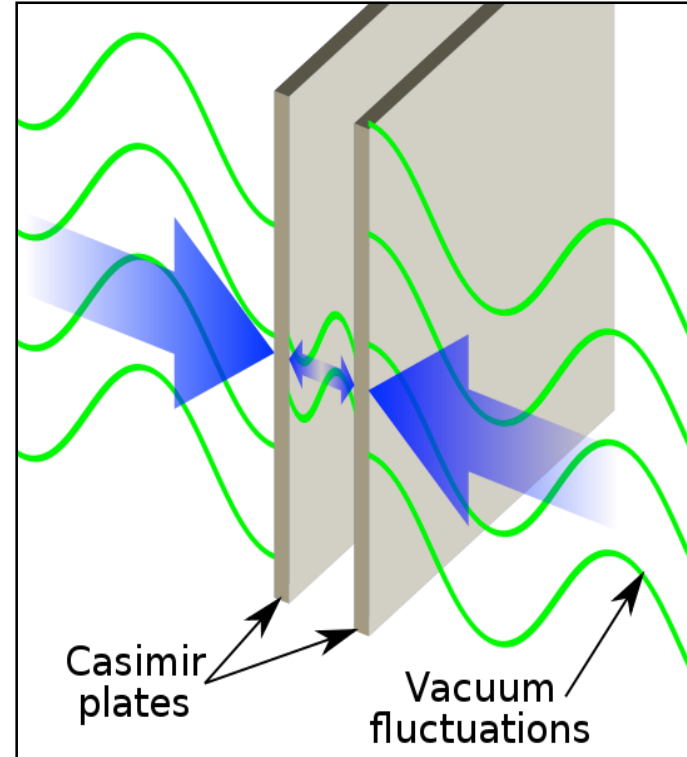
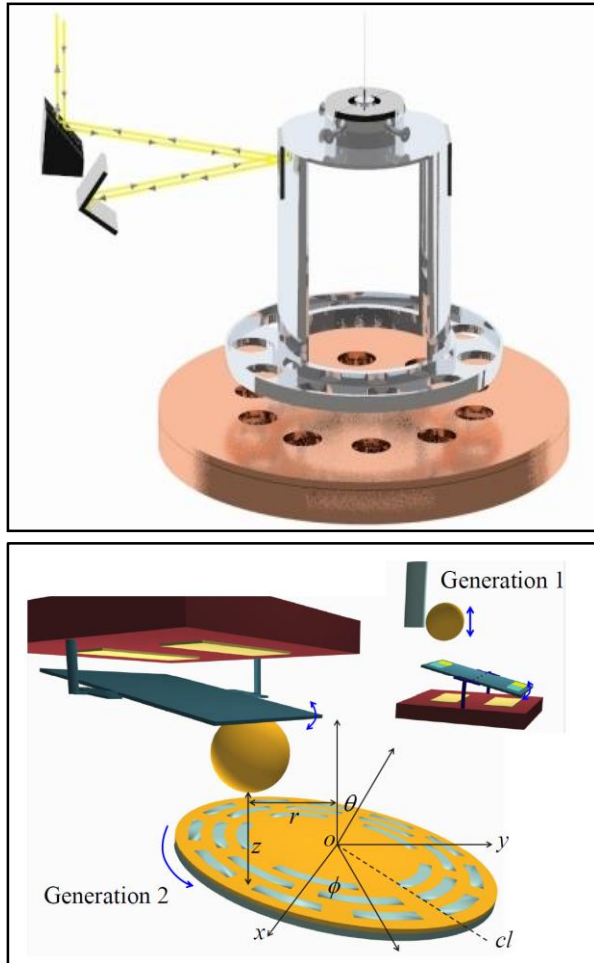
Jaffe, Haslinger, Xu, Hamilton, Upadhye, Elder, Khoury, Müller. (2017)

Rider, Moore, Blakemore, Louis, Lu, Gratta. (2016)

Ivanov, Hollwieser, Jenke, Wellenzohn, Abele (2013)

For a review see CB & Sakstein (2017)

Laboratory Searches – Short Range Forces

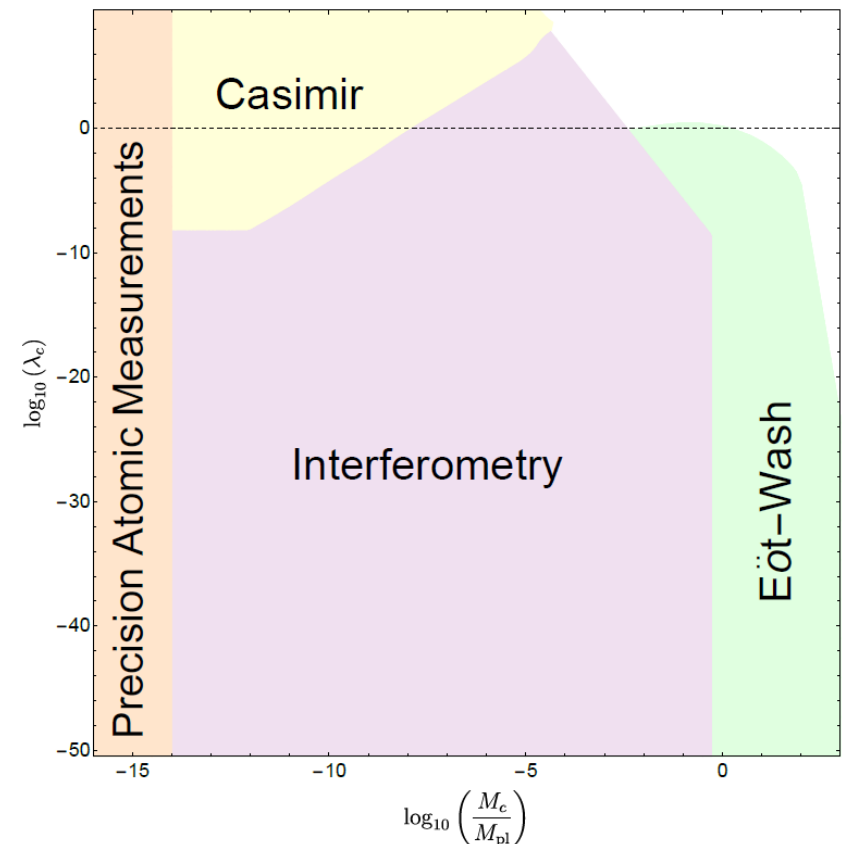
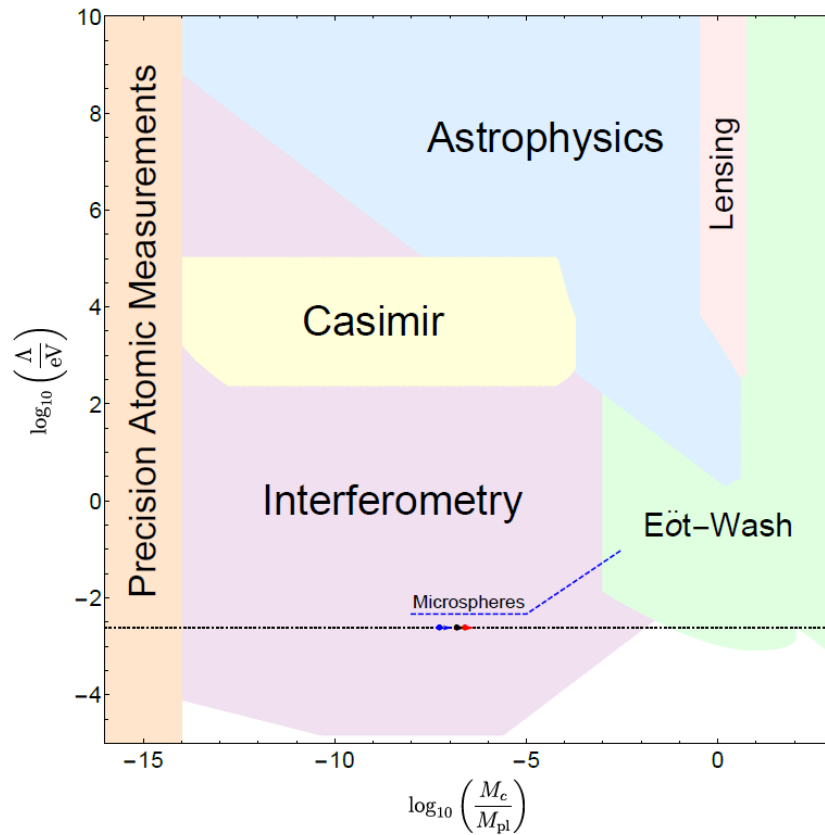


Upadhye (2012). Kapner, Cook, Adelberger, Gundlach, Heckel, Hoyle, Swanson. (2006)
Brax, van de Bruck, Davis, Mota, Shaw (2007)
Elder, Vardanyan, Akrami, Brax, Davis, Decca (2019)
For a review see CB & Sakstein (2017)

Combined Chameleon Constraints

$$V(\phi) = \frac{\Lambda^5}{\phi}$$

$$V(\phi) = \frac{\lambda}{4}\phi^4$$



Summary

Explanations for dark energy typically introduce new scalar fields but the corresponding long range forces are not seen

Screening mechanisms (non-linearities) hide these forces from fifth force searches

- Can still be detected in suitably designed experiments
- Atom interferometry a particularly powerful technique

Complementary to large scale cosmological surveys
e.g. Euclid, LSST

