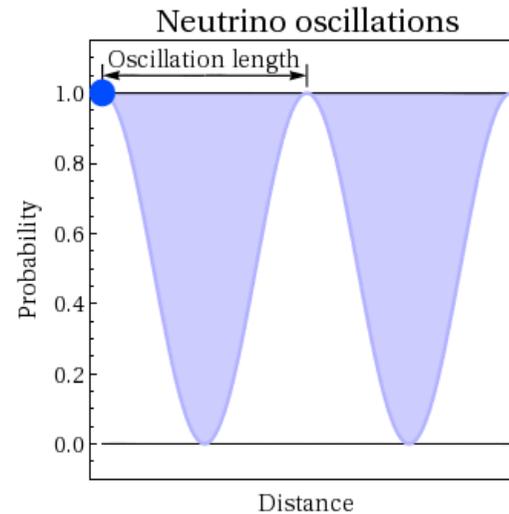


Perspectives for neutrino oscillations.



Walter Winter

DESY, 10./11.06.2014

Contents

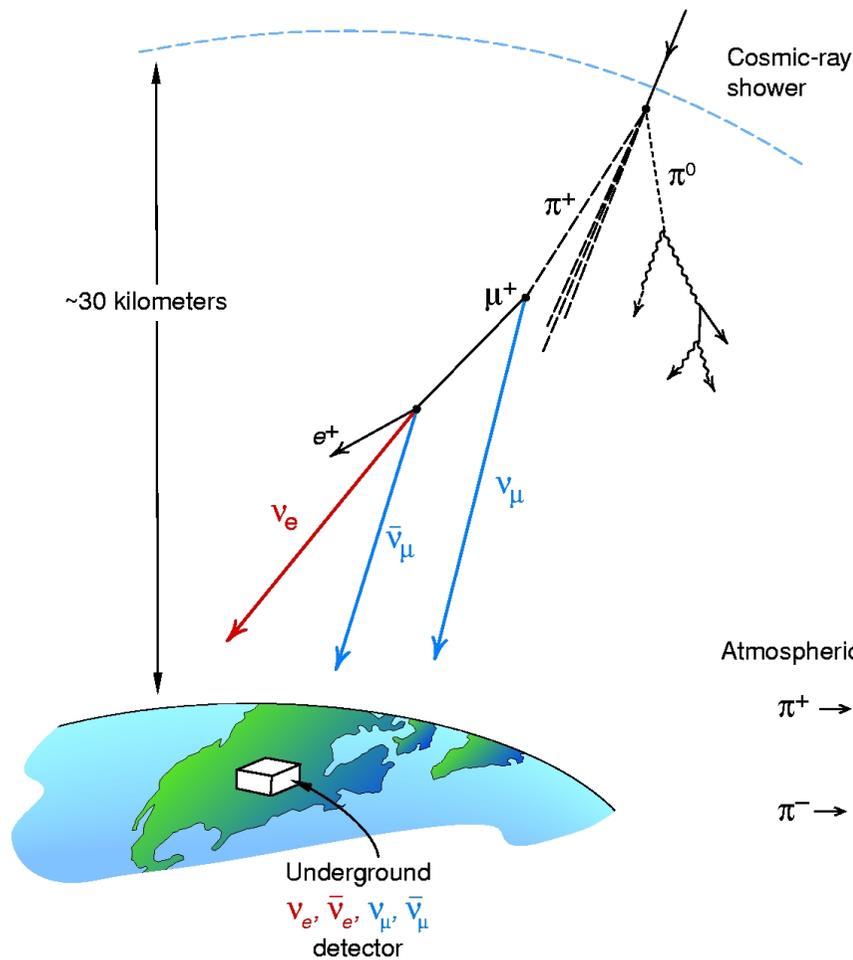
- > Introduction to neutrino oscillations
- > Current knowledge of neutrino oscillations
- > The future:
 - Measurement of δ_{CP}
 - Determination of the neutrino mass ordering
- > Summary and conclusions



Introduction to neutrino oscillations



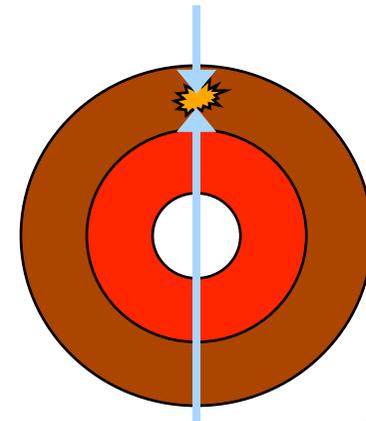
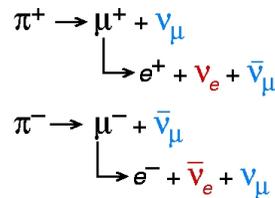
Neutrinos from the atmosphere



- > The rate of neutrinos should be the same from below and above
- > But: About 50% missing from below
- > Neutrino change their flavor on the path from production to detection: Neutrino oscillations
- > Neutrinos are massive!

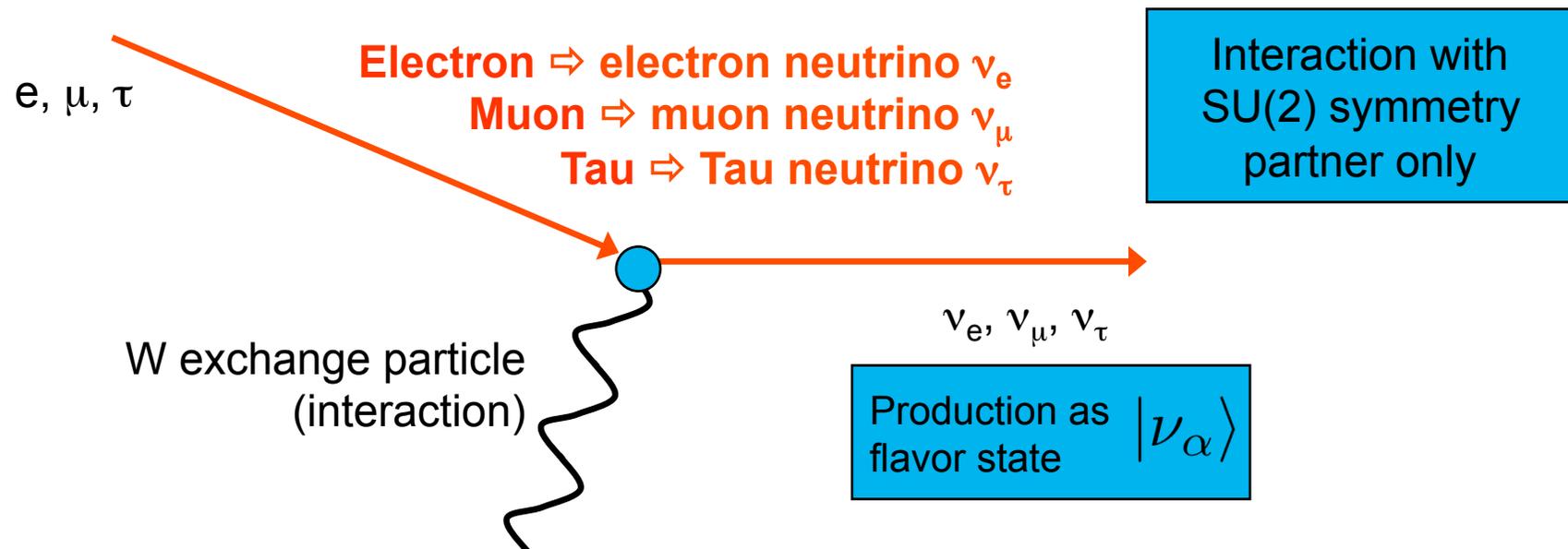
(Super-Kamiokande: “Evidence for oscillations of atmospheric neutrinos”, 1998)

Atmospheric neutrino source



Neutrino production/detection

- > Neutrinos are only produced and detected by the weak interaction:



- > The dilemma: One cannot assign a mass to the flavor states ν_e, ν_μ, ν_τ !

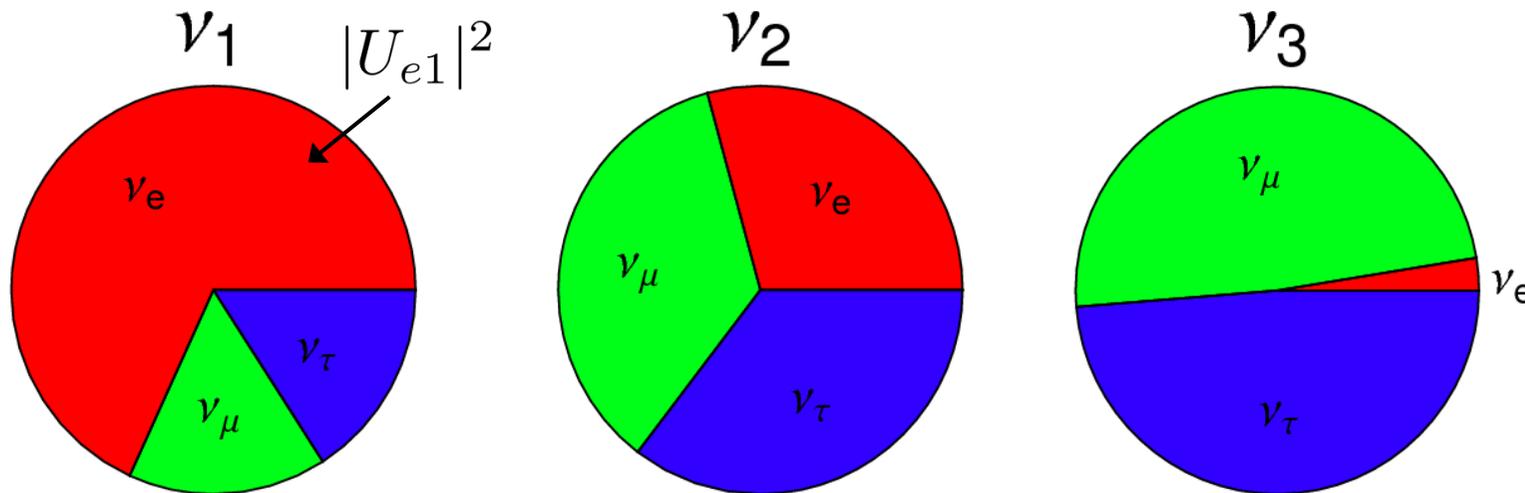
Which mass do the neutrinos have?

- > There is a set of neutrinos ν_1, ν_2, ν_3 , for which a mass can be assigned.

$$|\nu_i\rangle$$

- > **Mixture** of flavor states:

$$|\nu_\alpha\rangle = \sum_{k=1}^3 U_{\alpha k}^* |\nu_k\rangle$$



$$\sin^2 2\theta_{13} = 0.1, \delta = \pi/2$$

- > Not unusual, know from the Standard Model for quarks
- > However, the mixings of the neutrinos are much larger!



Neutrino oscillation probability

Standard derivation N active, S sterile (not weakly interacting) flavors

- > Mixing of flavor states

$$|\nu_\alpha\rangle = \sum_{k=1}^{N+S} U_{\alpha k}^* |\nu_k\rangle$$

- > Time evolution of mass state

$$|\nu_k(t)\rangle = \exp(-iE_k t) |\nu_k\rangle$$

- > Transition amplitude

$$A_{\nu_\alpha \rightarrow \nu_\beta} \equiv A_{\alpha\beta} = \langle \nu_\beta | \nu_\alpha(t) \rangle = \sum_{k=1}^{N+S} U_{\alpha k}^* U_{\beta k} \exp(-iE_k t)$$

- > Transition probability

$$P_{\alpha\beta} = A_{\alpha\beta}^* A_{\alpha\beta} = \sum_{k,j=1}^{N+S} \underbrace{U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^*}_{\equiv J_{kj}^{\alpha\beta}} \exp(-i(E_k - E_j)t)$$

“quartic re-phasing invariant”



Further simplifications

- > Ultrarelativistic approximations:

$$E_k = \sqrt{\vec{p}^2 + m_k^2} \simeq E + \frac{m_k^2}{2E}, \quad t \simeq L$$

L: baseline (distance source-detector)

- > Plus some manipulations: “Master formula”

$$P_{\alpha\beta} = \delta_{\alpha\beta} - \underbrace{4 \sum_{k>j} \text{Re} J_{kj}^{\alpha\beta} \sin^2 \left(\frac{\Delta m_{kj}^2 L}{4E} \right)}_{\text{CP conserving}} + \underbrace{2 \sum_{k>j} \text{Im} J_{kj}^{\alpha\beta} \sin \left(\frac{\Delta m_{kj}^2 L}{2E} \right)}_{\text{CP violating}}$$

$\Delta m_{kj}^2 \equiv m_k^2 - m_j^2$ “mass squared difference”

$F(L,E)=L/E$ “spectral dependence”

- > For antineutrinos: $U \Leftrightarrow U^*$



Two flavor limit: N=2, S=0

- > Only two parameters:

$$U = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix}$$

- > From the master formula:
Disappearance or **survival** probability

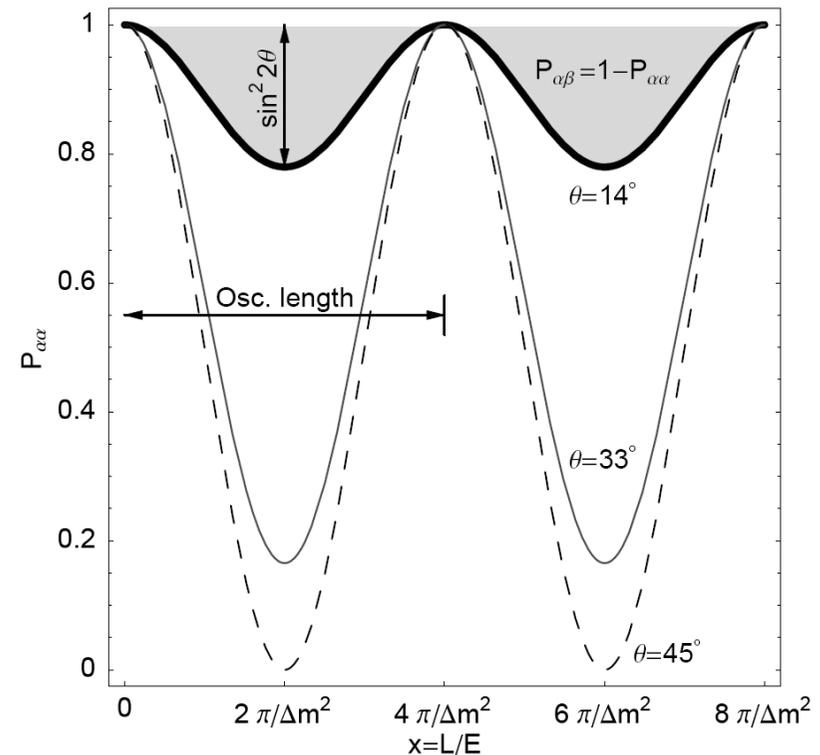
$$P_{\alpha\alpha} = 1 - \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E} \right)$$

Appearance probability

$$P_{\alpha\beta} = \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E} \right)$$

Lower limit for neutrino mass!

$$\Delta m^2 \equiv m_2^2 - m_1^2$$



Three flavors: Mixings

- > Use same parameterization as for CKM matrix

Potential CP violation $\sim \theta_{13}$

$$U_{\text{PMNS}} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13} e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13} e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}.$$

$(s_{ij} = \sin \theta_{ij} \quad c_{ij} = \cos \theta_{ij})$

$$= \left(\text{Image 1} \right) \times \left(\text{Image 2} \right) \times \left(\text{Image 3} \right)$$

Pontecorvo-Maki-Nakagawa-Sakata matrix

- > Neutrinos \Leftrightarrow Anti-neutrinos: $\mathbf{U} \Leftrightarrow \mathbf{U}^*$ (neutrino oscillations)
- > If neutrinos are their own anti-particles (Majorana neutrinos):
 $\mathbf{U} \Leftrightarrow \mathbf{U} \text{diag}(1, e^{i\alpha}, e^{i\beta})$ - do enter $0\nu\beta\beta$, but not neutrino oscillations

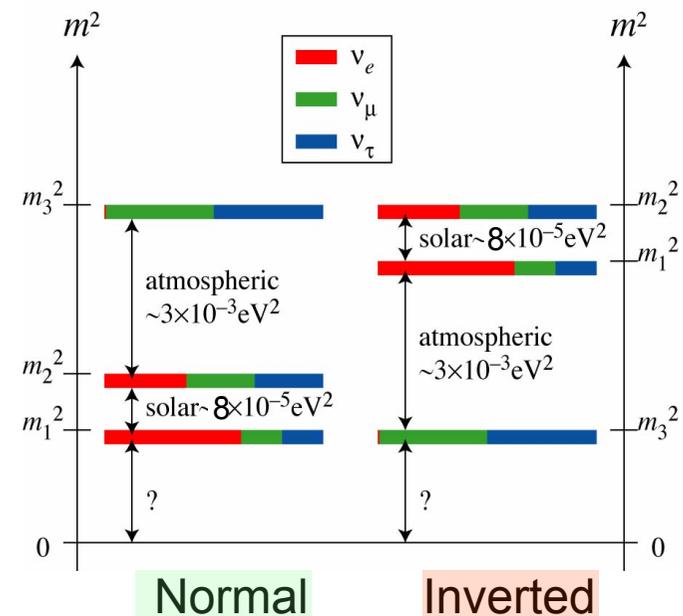


Three active flavors: Masses

- > Two independent mass squared splittings, typically Δm_{21}^2 (solar) Δm_{31}^2 (atmospheric)

Will be relevant for neutrino oscillations!

- > The third is given by $\Delta m_{32}^2 = \Delta m_{31}^2 - \Delta m_{21}^2$
- > The (atmospheric) mass **ordering** (hierarchy) is unknown (normal or inverted)
- > The absolute neutrino mass scale is unknown ($< eV$)



Current knowledge of neutrino oscillations



Three flavors: Simplified

> What we know (qualitatively):

- Hierarchy of mass splittings $\Delta m_{21}^2 \ll |\Delta m_{31}^2| \simeq |\Delta m_{32}^2|$

- Two mixing angles large, one (θ_{13}) small ~ 0 ?

$$U_{\text{PMNS}}^{\theta_{13} \rightarrow 0} = \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} c_{23} & c_{12} c_{23} & s_{23} \\ s_{12} s_{23} & -c_{12} s_{23} & c_{23} \end{pmatrix}$$

> From the “master formula“, we have

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \left(J_{31}^{\alpha\beta} + J_{32}^{\alpha\beta} \right) \sin^2 \Delta_{31} - 4 J_{21}^{\alpha\beta} \sin^2 \Delta_{21}$$

$$J_{kj}^{\alpha\beta} = U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^* \quad \Delta_{ij} \equiv \Delta m_{ij}^2 L / (4E)$$



Two flavor limits

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \left(J_{31}^{\alpha\beta} + J_{32}^{\alpha\beta} \right) \sin^2 \Delta_{31} - 4 J_{21}^{\alpha\beta} \sin^2 \Delta_{21}$$

$$\Delta_{ij} \equiv \Delta m_{ij}^2 L / (4E)$$

Two flavor limits by selection of frequency:

- Atmospheric frequency: $\Delta_{31} \sim \pi/2 \Rightarrow \Delta_{21} \ll 1$

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \left(J_{31}^{\alpha\beta} + J_{32}^{\alpha\beta} \right) \sin^2 \Delta_{31} - 4 J_{21}^{\alpha\beta} \sin^2 \Delta_{21}$$

- Solar frequency: $\Delta_{21} \sim \pi/2 \Rightarrow \Delta_{31} \gg 1$

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \left(J_{31}^{\alpha\beta} + J_{32}^{\alpha\beta} \right) \sin^2 \Delta_{31} - 4 J_{21}^{\alpha\beta} \sin^2 \Delta_{21}$$

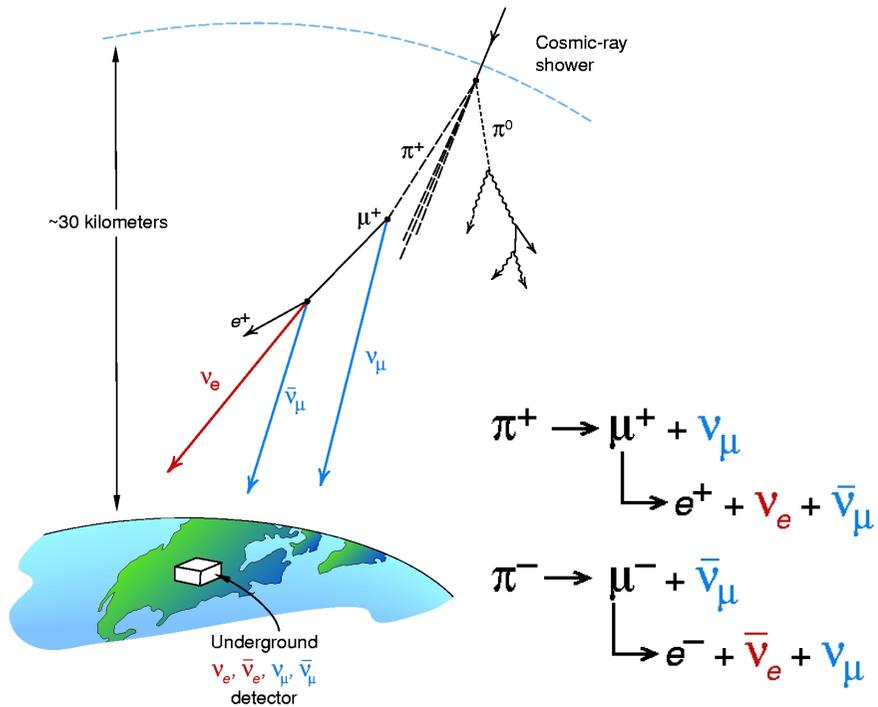
Select sensitive term
by choice of L/E!

$\sin^2 \Delta_{31}$
averages
out

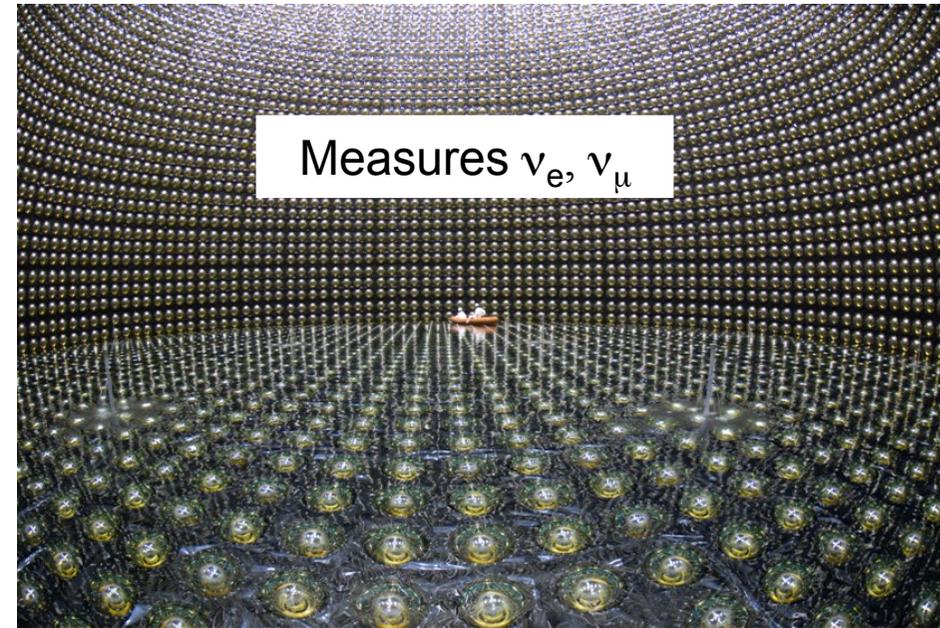
0.5



Atmospheric neutrinos



Super-Kamiokande



From $P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \left(J_{31}^{\alpha\beta} + J_{32}^{\alpha\beta} \right) \sin^2 \Delta_{31}$ and θ_{13} small
 we have: $P_{ee} \sim 1$, $P_{e\mu} \sim P_{\mu e} \sim 0$ and

$$P_{\mu\mu} \simeq 1 - \sin^2(2\theta_{23}) \sin^2 \Delta_{31}$$

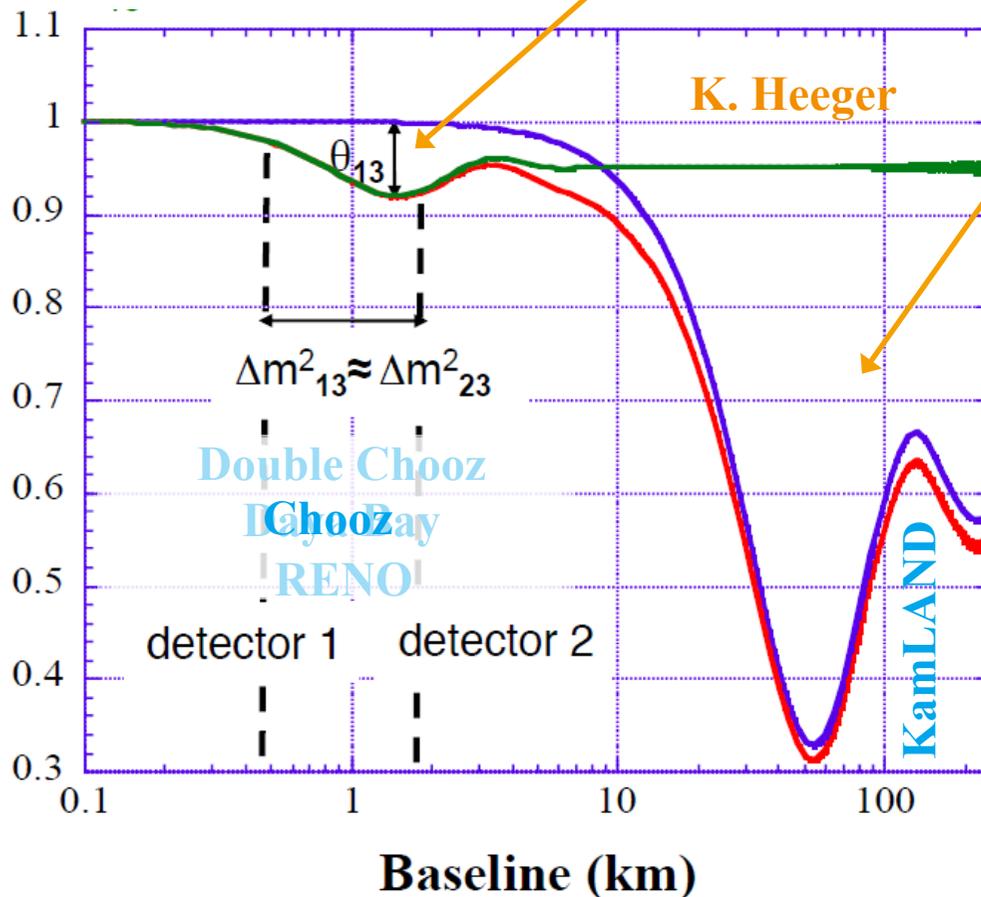
⇒ Two flavor limit with particular parameters $\theta_{23}, \Delta m_{31}^2$



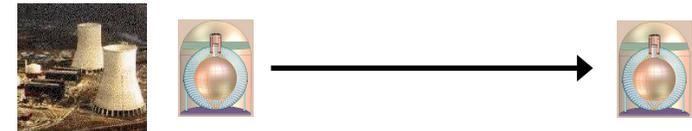
Reactor neutrinos

> In the presence of θ_{13} :

$$P_{\bar{e}\bar{e}} \simeq 1 - \underbrace{\sin^2(2\theta_{13}) \sin^2 \Delta_{31}}_{\text{atmospheric}} - \underbrace{\cos^4 \theta_{13} \sin^2(2\theta_{12}) \sin^2 \Delta_{21}}_{\text{solar}}$$



New idea:



Identical detectors, $L \sim 1-2$ km
to control systematics

(Minakata, Sugiyama, Yasuda,
Inoue, Suekane, 2003;
Huber, Lindner, Schwetz,
Winter, 2003)



$$P_{\bar{e}\bar{e}} \simeq 1 - \sin^2(2\theta_{13}) \sin^2 \Delta_{31}$$

(short distance)

Kirk T McDonald
Princeton U
(April 24, 2012)

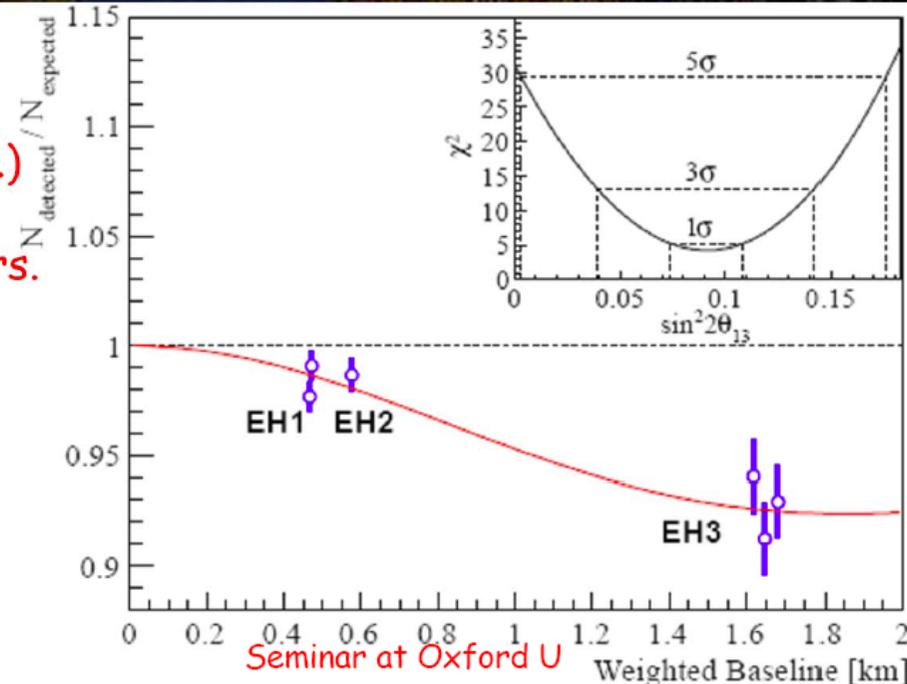
on behalf of the Daya Bay Collaboration



We observe that
 $\sin^2 2\theta_{13} = 0.092 \pm 0.016$ (stat.) ± 0.005 (syst.)
after 55 days of operation with 6 detectors
at 3 sites close to 3 pairs of ~ 3 GW reactors.

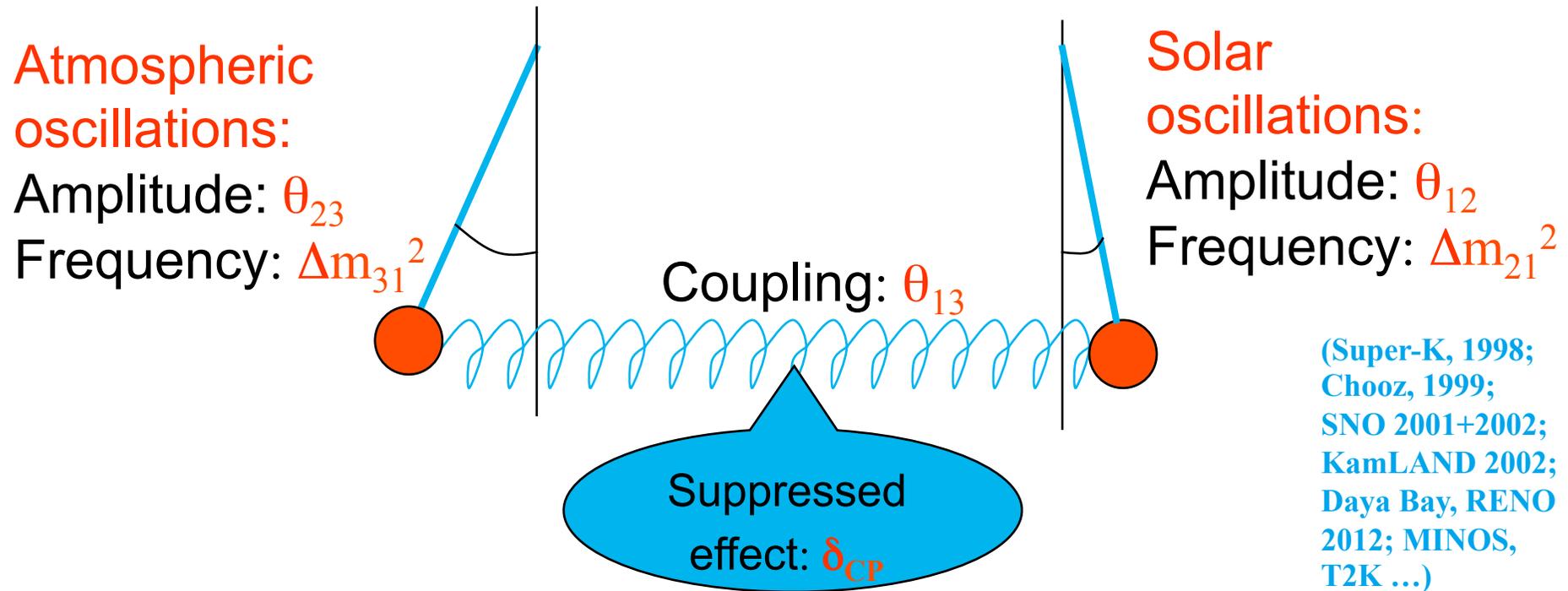
F.P. Ahn *et al.*
Phys. Rev. Lett. **108**, 171803 (2012).

(also: T2K, Double Chooz, RENO)



Three flavors: Summary

- > Three flavors: 6 params (3 angles, one phase; 2 x Δm^2)



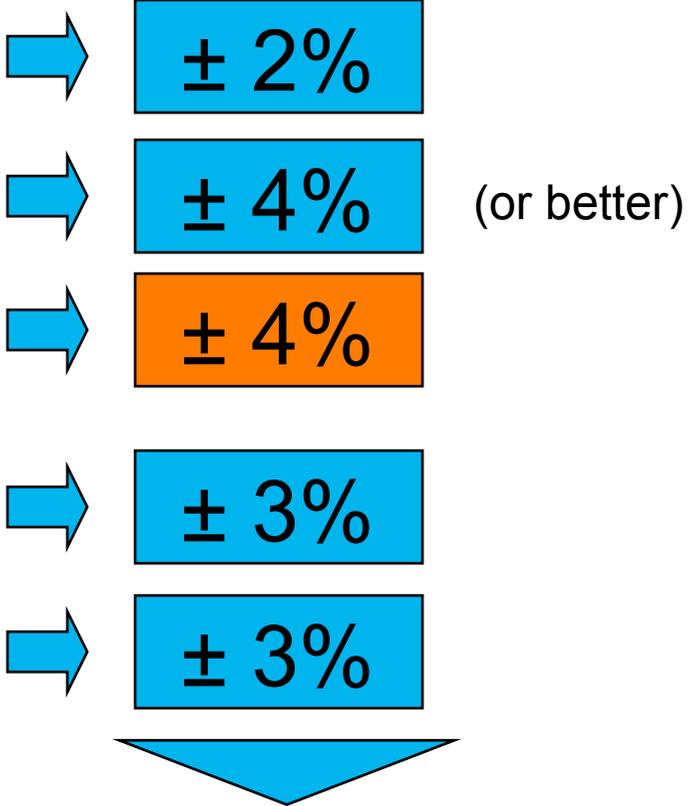
- > Describes solar and atmospheric neutrino anomalies, as well as reactor antineutrino disappearance!



Precision of parameters?

	bfp $\pm 1\sigma$	3σ range
$\sin^2 \theta_{12}$	$0.306^{+0.012}_{-0.012}$	0.271 \rightarrow 0.346
$\theta_{12}/^\circ$	$33.57^{+0.77}_{-0.75}$	31.38 \rightarrow 36.01
$\sin^2 \theta_{23}$	$0.446^{+0.007}_{-0.007} \oplus 0.587^{+0.032}_{-0.037}$	0.366 \rightarrow 0.663
$\theta_{23}/^\circ$	$41.9^{+0.4}_{-0.4} \oplus 50.0^{+1.9}_{-2.2}$	37.2 \rightarrow 54.5
$\sin^2 \theta_{13}$	$0.0229^{+0.0020}_{-0.0019}$	0.0170 \rightarrow 0.0288
$\theta_{13}/^\circ$	$8.71^{+0.37}_{-0.38}$	7.50 \rightarrow 9.78
$\delta_{CP}/^\circ$	265^{+56}_{-61}	0 \rightarrow 360
$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$	$7.45^{+0.19}_{-0.16}$	6.98 \rightarrow 8.05
$\frac{\Delta m_{31}^2}{10^{-3} \text{ eV}^2}$ (N)	$+2.417^{+0.013}_{-0.013}$	+2.247 \rightarrow +2.623
$\frac{\Delta m_{32}^2}{10^{-3} \text{ eV}^2}$ (I)	$-2.410^{+0.062}_{-0.062}$	-2.602 \rightarrow -2.226

NuFIT 1.2 (2013)



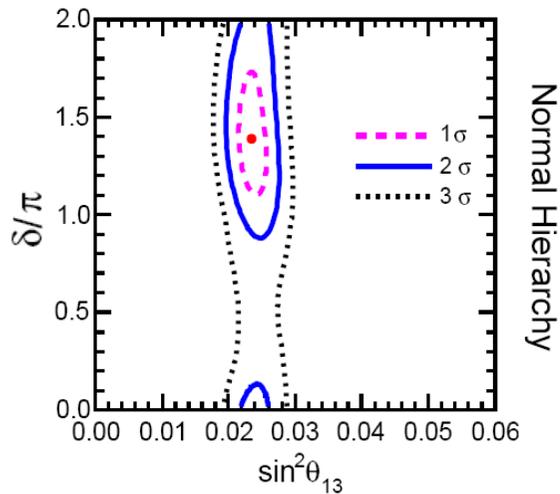
Open issues:

- Degeneracies (mass ordering, octant)
- CP phase



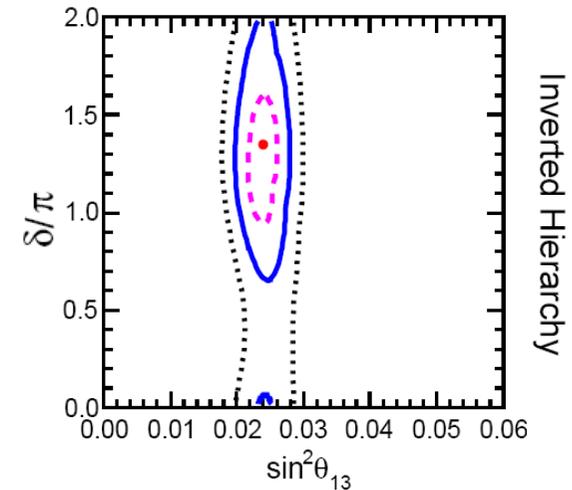
Current status and perspectives for existing equipment

- Indication for δ_{CP} , no evidence for mass hierarchy



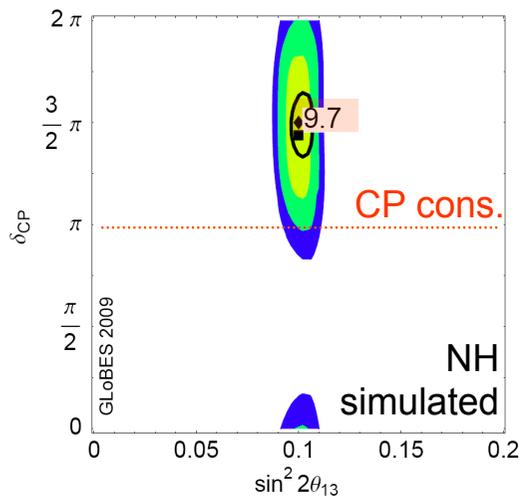
LBL Acc + Solar + KL
+ SBL Reactors
+ SK Atm

Capozzi, Fogli,
Lisi, Marrone,
Montanino, Palazzo,
arXiv:1312.2878



Main difference: NOvA data!

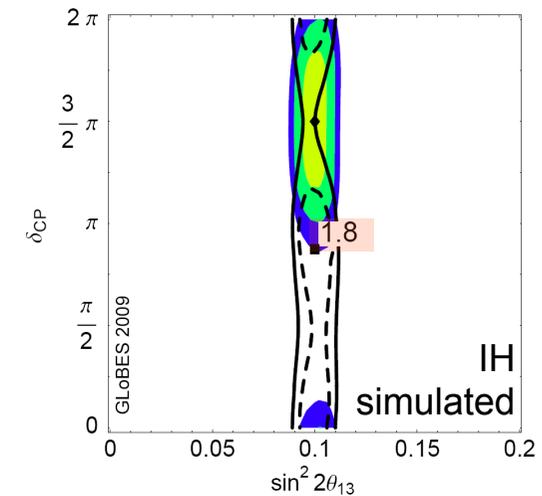
- Potential of existing equipment



T2K, NOvA,
Double Chooz, Daya Bay;
5 years each

High CL determination
requires new equipment

Huber, Lindner, Schwetz, Winter,
JHEP 0911 (2009) 044



NEUTRINO 2014

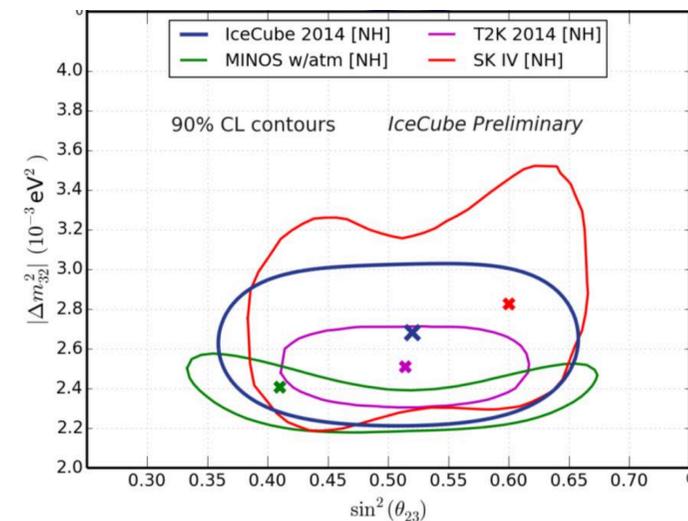
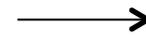
XXVI International Conference on Neutrino Physics and Astrophysics

June 2-7, 2014, Boston, U.S.A.

> Largest, bi-annual conference of the neutrino community, 550 participants

> Highlights:

- Discovery of cosmic neutrinos, by IceCube
- First direct high-CL evidence for ν_μ to ν_e flavor transitions, by T2K
- First atmospheric measurement of leading atmospheric parameters comparable with long-baseline experiments, by IceCube-DeepCore **(analysis done by DESY-Zeuthen group)**



Juan-Pablo Yanez @ Neutrino 2014

- θ_{13} central value shifts down (Daya Bay), tension with T2K increased; may strengthen hint for maximal leptonic CP violation $\delta_{CP} \sim -\pi/2$
- P5 prioritisation of Long Baseline Neutrino Oscillation Experiment at Fermilab
- DESY with three plenary talks one of the strongest represented institutions in the program (after INFN, before Fermilab)



The future: measurement of δ_{CP}



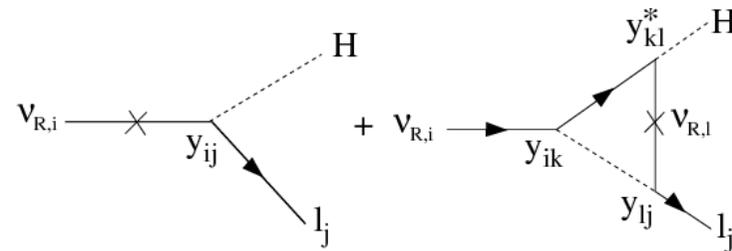
Why is δ_{CP} interesting?

> CP violation $\sin\delta$

Necessary condition for successful baryogenesis

(dynamical mechanism to create matter-antimatter asymmetry of the universe)

⇒ thermal leptogenesis by decay of heavy see-saw partner?



> Model building $\cos\delta$

$$U_{PMNS} = U_{\ell}^{\dagger} U_{\nu}$$

CKM-like correction
leading to non-zero θ_{13} ?

Symmetry
e.g. TBM, BM, ...?

e.g. TBM sum rule: $\theta_{12} = 35^{\circ} + \theta_{13} \cos\delta$ (Antusch, King; Masina ...)

> Need performance which is equally good for all δ_{CP}



Necessary conditions for the observation of CP violation

> Since

$$\left\langle \sin \left(\frac{\Delta m_{kj}^2 L}{2E} \right) \right\rangle_{L/E} = 0$$

⇒ need spectral info!

> Since for $\alpha=\beta$

$$J_{kj}^{\alpha\alpha} = |U_{\alpha k}|^2 |U_{\alpha j}|^2$$

⇒ need to observe flavor transitions

> Need (at least) three flavors
(actually conclusion in quark sector by
[Kobayashi, Maskawa, Nobel Prize 2008](#))

⇒ No CP violation in two flavor subspaces!

⇒ Need to be sensitive to (at least) two mass squared splittings at the same time!

$$P_{\alpha\beta}^{\text{CPV}} \simeq$$

~ Jarlskog invariant

$$-2 \sum_{k>j} \text{Im} J_{kj}^{\alpha\beta} \sin \left(\frac{\Delta m_{kj}^2 L}{2E} \right)$$

CP violating

$$J_{kj}^{\alpha\beta} = U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^*$$



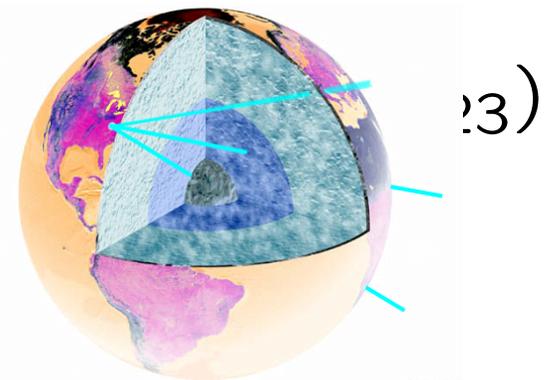
Electron-muon neutrino flavor transitions

$$\begin{aligned}
 P_{e\mu} &\simeq \sin^2 2\theta_{13} \frac{\cos^2 \theta_{23}}{\sin^2 \theta_{23}} \frac{\sin^2[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})^2} & \alpha &\equiv \frac{\Delta m_{21}^2}{\Delta m_{31}^2}, \Delta \equiv \frac{\Delta m_{31}^2 L}{4E}, \hat{A} \equiv \frac{2\sqrt{2}G_F n_e E}{\Delta m_{31}^2} \\
 &+ \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{\text{CP}} \sin(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})} \\
 &+ \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{\text{CP}} \cos(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 \pm \hat{A})\Delta]}{(1 \pm \hat{A})} \\
 &+ \alpha^2 \frac{\sin^2 \theta_{23}}{\cos^2 \theta_{23}} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}
 \end{aligned}$$

> Antineutrinos: $P_{\bar{e}\bar{\mu}} = P_{e\mu}(\delta_{\text{CP}}, \rightarrow -\delta_{\text{CP}}, \hat{A} \rightarrow -\hat{A})$

> Silver: $P_{e\tau} = P_{e\mu}(s_{23}^2 \leftrightarrow c_{23}^2, \sin 2\theta_{13})$

> Platinum, T-inv.: $P_{\mu e} = P_{e\mu}(\delta_{\text{CP}}, \rightarrow -\delta_{\text{CP}}, \hat{A} \rightarrow -\hat{A})$



(Cervera et al. 2000; Freund, Huber, Lindner, 2000; Akhmedov et al, 2004)

Possible experimental setups (future)

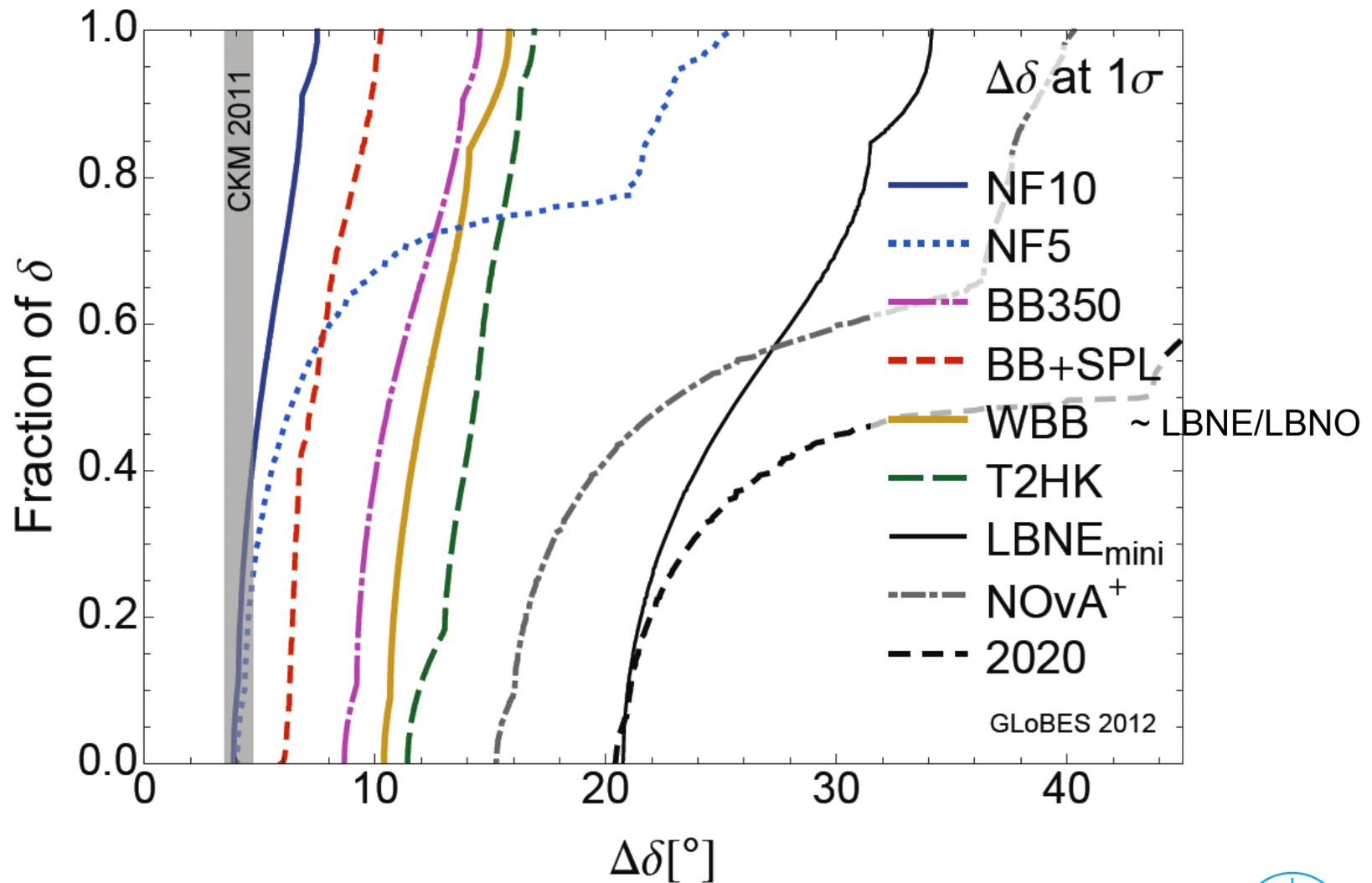
	Setup	E_ν^{peak}	L	OA	Detector	kt	MW	Decays/yr	$(t_\nu, t_{\bar{\nu}})$
→ Benchmark	BB350	1.2	650	–	WC	500	–	$1.1(2.8) \times 10^{18}$	(5,5)
	NF10	5.0	2 000	–	MIND	100	–	7×10^{20}	(10,10)
	WBB	4.5	2 300	–	LAr	100	0.8	–	(5,5)
	T2HK	0.6	295	2.5°	WC	560	1.66	–	(1.5,3.5)
→ Alternative	BB100	0.3	130	–	WC	500	–	$1.1(2.8) \times 10^{18}$	(5,5)
	+ SPL			–			4		–
	NF5	2.5	1 290	–	MIND	100	–	7×10^{20}	(10,10)
	LBNE _{mini}	4.0	1 290	–	LAr	10	0.7	–	(5,5)
	NO ν A ⁺	2.0	810	0.8°	LAr	30	0.7	–	(5,5)
2020	T2K	0.6	295	2.5°	WC	22.5	0.75	–	(5,5) + Daya Bay
	NO ν A	2.0	810	0.8°	TASD	15	0.7	–	(4,4)

(Coloma, Huber, Kopp, Winter, arXiv:1209.5973)



Performance

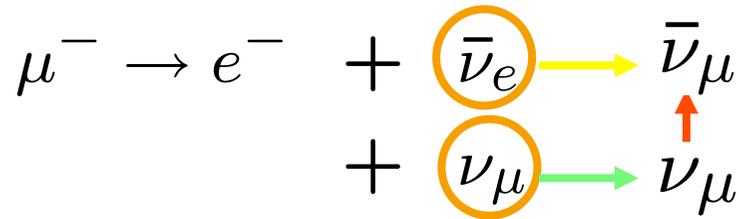
➤ Comparison at default systematics:



(Coloma, Huber, Kopp,
Winter, arXiv:1209.5973)



THE INTERNATIONAL DESIGN STUDY FOR THE NEUTRINO FACTORY

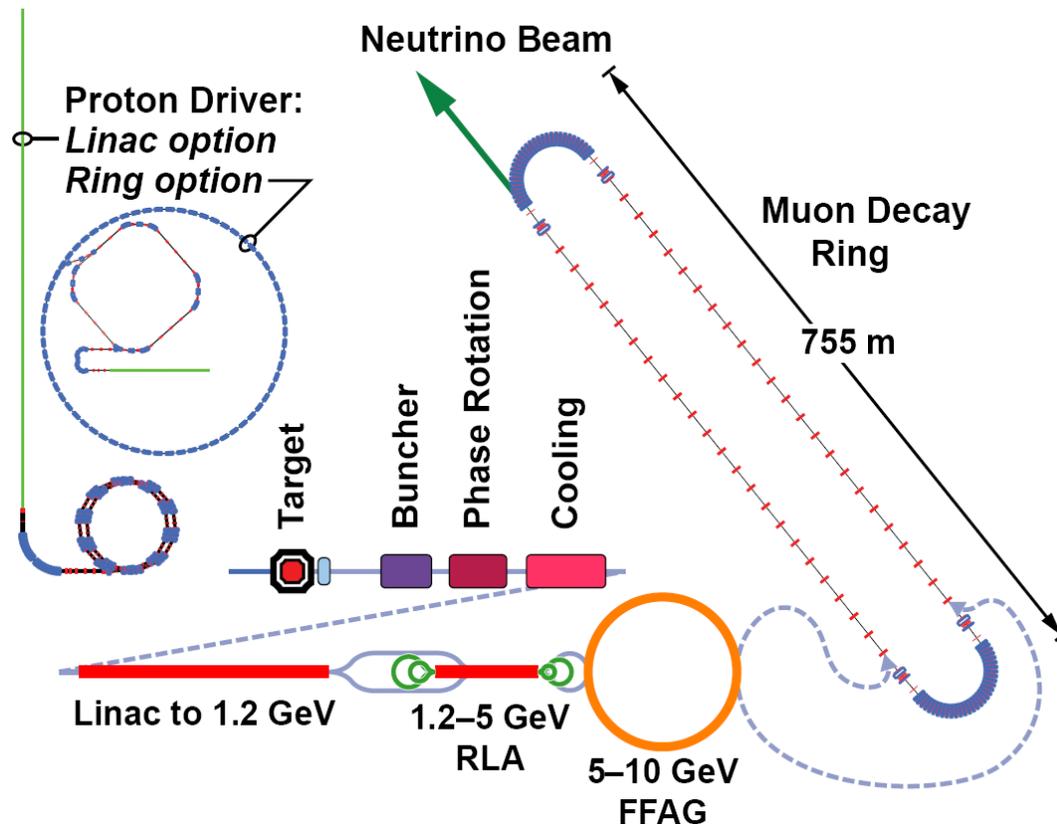


(Geer, 1997; de Rujula, Gavela, Hernandez, 1998; Cervera et al, 2000)

→ Signal prop. $\sin^2 2\theta_{13}$

→ Contamination \Rightarrow magnetized detector!

Muons decay in straight sections of a storage ring



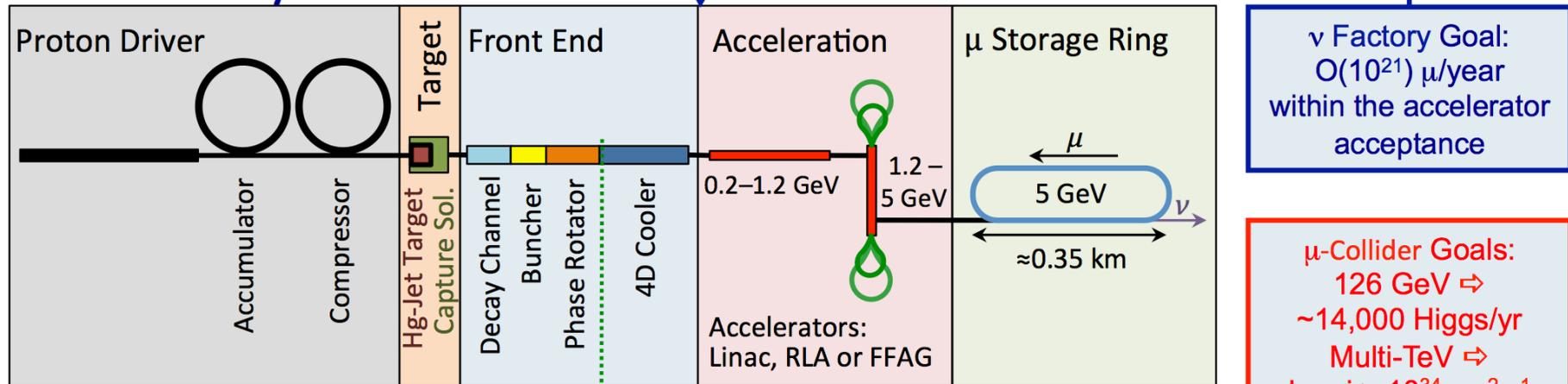
> IDS-NF:
Initiative from
~ 2007-2014 to
present a design
report, schedule, cost
estimate, risk
assessment for a
neutrino factory



Muon Accelerator Staging Programme

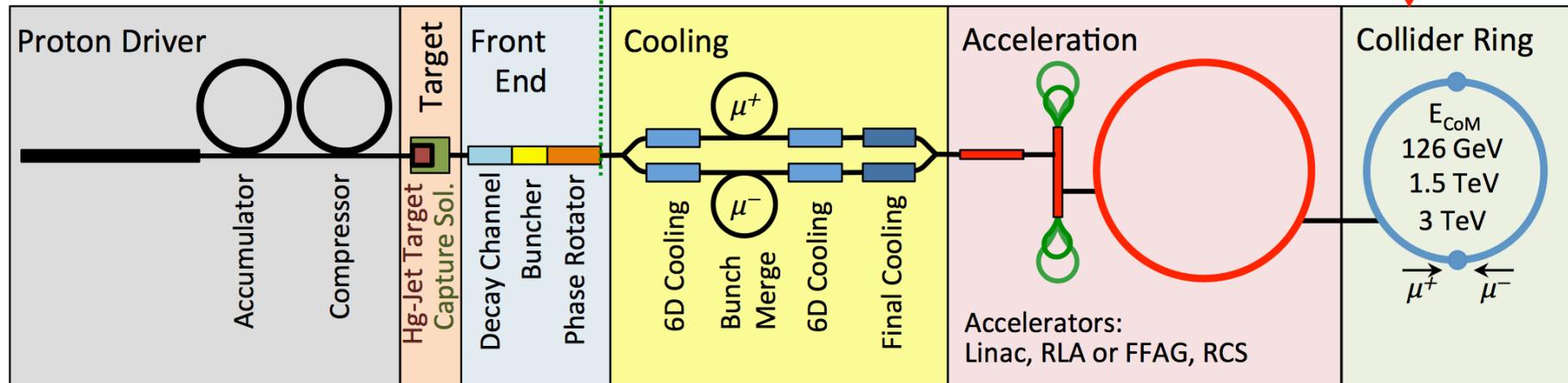
Synergies between NuMAX and Muon Collider components
Muon Accelerator Staging Study (MASS)

Neutrino Factory

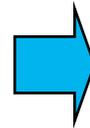
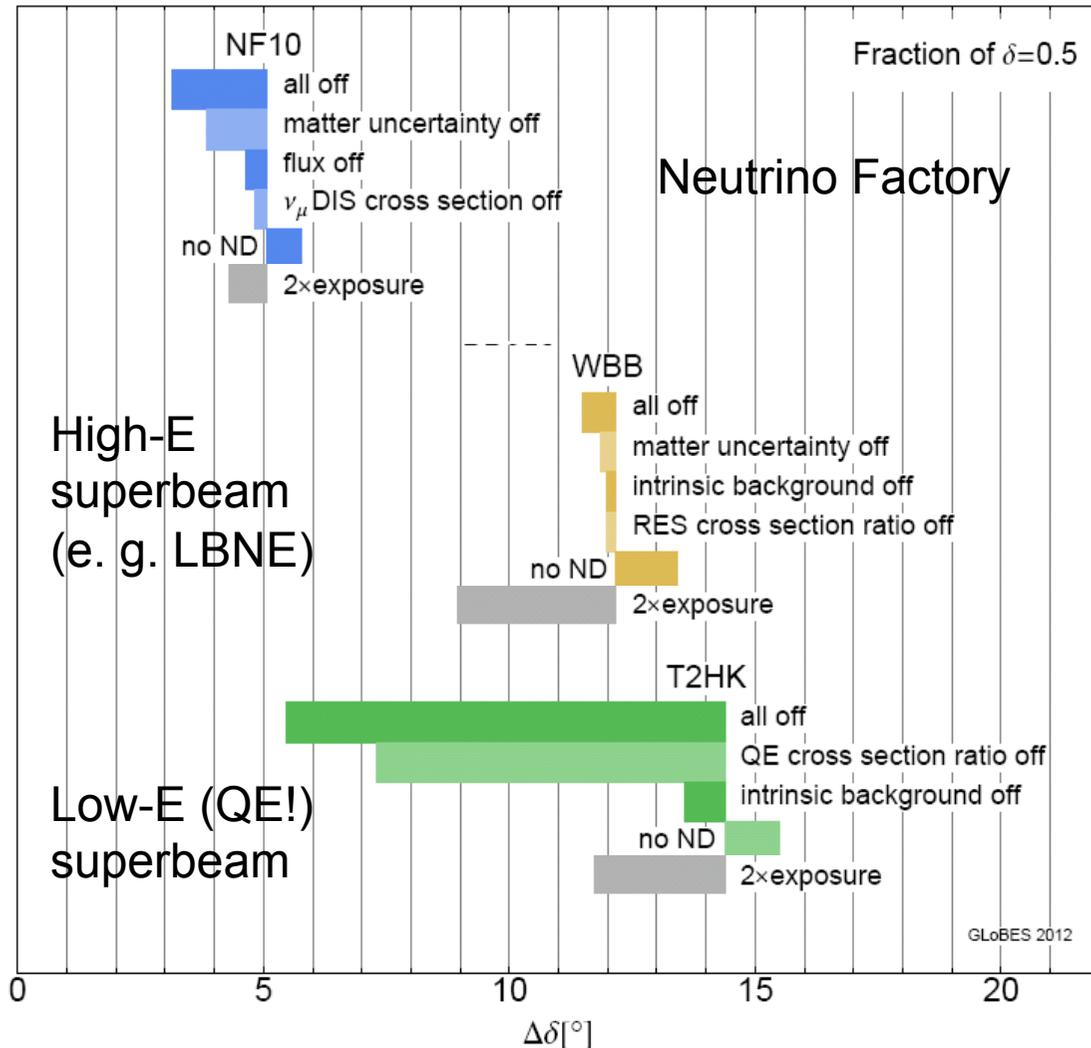


← Share same complex

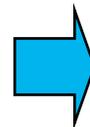
Muon Collider



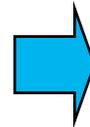
Systematics: Main challenges for δ_{CP}



Robust wrt systematics
Main impact:
Matter density uncertainty



Operate in statistics-limited regime
Exposure more important than near detector



QE ν_e X-sec critical: cannot be measured in near detector
Theory: ν_e/ν_μ ratio?
Experiment: 
 $\mu^- \rightarrow e^- + \bar{\nu}_e + \nu_\mu$

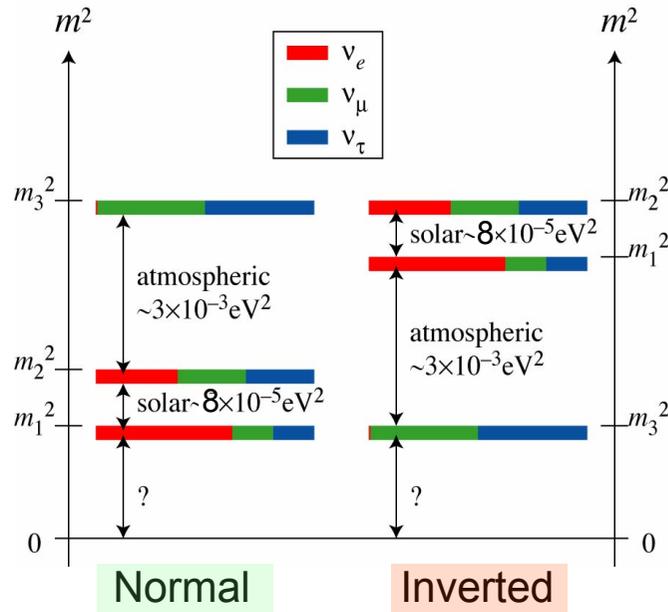
(Coloma, Huber, Kopp, Winter, arXiv:1209.5973)



Mass hierarchy determination



Why would one like to measure the mass ordering?



➤ Specific models typically come together with specific MH prediction (e.g. textures are very different)

➤ Good model discriminator \longrightarrow

(Albright, Chen, [hep-ph/0608137](https://arxiv.org/abs/hep-ph/0608137))

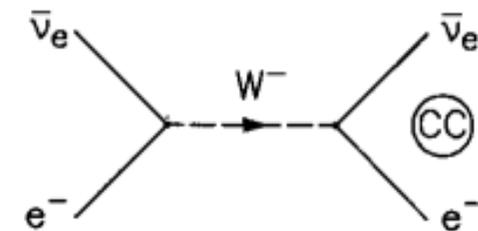
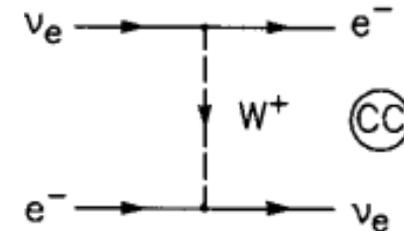
TABLE I: Mixing Angles for Models with Lepton Flavor Symmetry.

Reference	Hierarchy	$\sin^2 2\theta_{23}$	$\tan^2 \theta_{12}$	$\sin^2 \theta_{13}$
Anarchy Model:				
dGM [18]	Either			≥ 0.011 @ 2σ
$L_e - L_\mu - L_\tau$ Models:				
BM [35]	Inverted			0.00029
BCM [36]	Inverted			0.00063
GMN1 [37]	Inverted		≥ 0.52	≤ 0.01
GL [38]	Inverted			0
PR [39]	Inverted		≤ 0.58	≥ 0.007
S_3 and S_4 Models:				
CFM [40]	Normal			0.00006 - 0.001
HLM [41]	Normal	1.0	0.43	0.0044
	Normal	1.0	0.44	0.0034
KMM [42]	Inverted	1.0		0.000012
MN [43]	Normal			0.0024
MNY [44]	Normal			0.000004 - 0.000036
MPR [45]	Normal			0.006 - 0.01
RS [46]	Inverted	$\theta_{23} \geq 45^\circ$		≤ 0.02
	Normal	$\theta_{23} \leq 45^\circ$		0
TY [47]	Inverted	0.93	0.43	0.0025
T [48]	Normal			0.0016 - 0.0036
A_4 Tetrahedral Models:				
ABGMP [49]	Normal	0.997 - 1.0	0.365 - 0.438	0.00069 - 0.0037
AKKL [50]	Normal			0.006 - 0.04
Ma [51]	Normal	1.0	0.45	0
SO(3) Models:				
M [52]	Normal	0.87 - 1.0	0.46	0.00005
Texture Zero Models:				
CPP [53]	Normal			0.007 - 0.008
	Inverted			≥ 0.00005
	Inverted			≥ 0.032
WY [54]	Either			0.0006 - 0.003
	Either			0.002 - 0.02
	Either			0.02 - 0.15

Method 1: Matter effects in neutrino oscillations

- > Ordinary matter: electrons, but no μ, τ
- > Coherent forward scattering in matter: Net effect on electron flavor
- > Hamiltonian in matter (matrix form, flavor space):

(Wolfenstein, 1978; Mikheyev, Smirnov, 1985)



$$\mathcal{H}(n_e) = U \begin{pmatrix} 0 & 0 & 0 \\ 0 & \frac{\Delta m_{21}^2}{2E} & 0 \\ 0 & 0 & \frac{\Delta m_{31}^2}{2E} \end{pmatrix} U^\dagger + \begin{pmatrix} V(n_e) & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$V_\nu = +\sqrt{2}G_F n_e, \quad V_{\bar{\nu}} = -\sqrt{2}G_F n_e, \quad n_e = Y \rho_j / m_N$

Y : electron fraction ~ 0.5
 (electrons per nucleon)



Parameter mapping ... for two flavors, constant matter density

Oscillation probabilities in

vacuum: $P_{\alpha\alpha} = 1 - \sin^2 2\theta \sin^2 \frac{\Delta m^2 L}{4E}$

matter: $P_{\alpha\alpha} = 1 - \sin^2 2\tilde{\theta} \sin^2 \frac{\Delta \tilde{m}^2 L}{4E}$

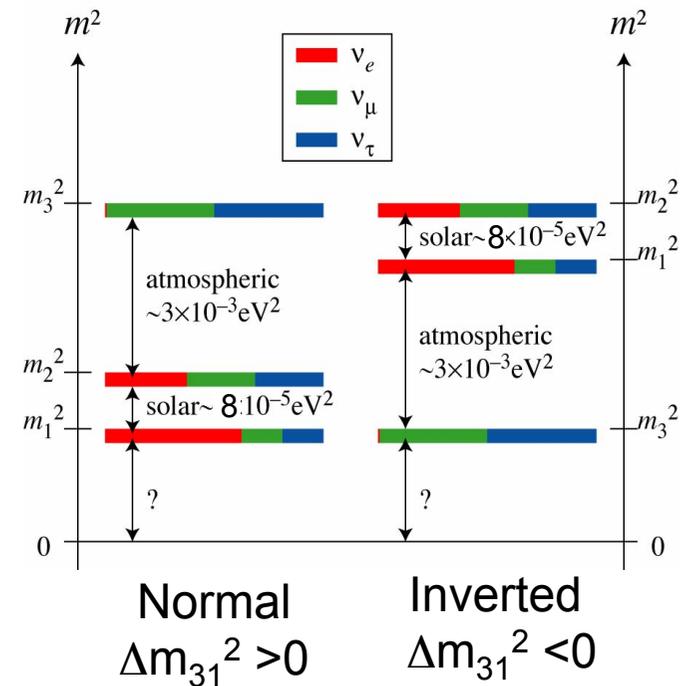
$$\Delta \tilde{m}^2 = \xi \cdot \Delta m^2, \quad \sin 2\tilde{\theta} = \frac{\sin 2\theta}{\xi},$$

$$\xi \equiv \sqrt{\sin^2 2\theta + (\cos 2\theta - \hat{A})^2},$$

$$\hat{A} = \frac{2EV}{\Delta m^2} = \frac{\pm 2\sqrt{2}E G_F n_e}{\Delta m^2}$$

Enhancement condition

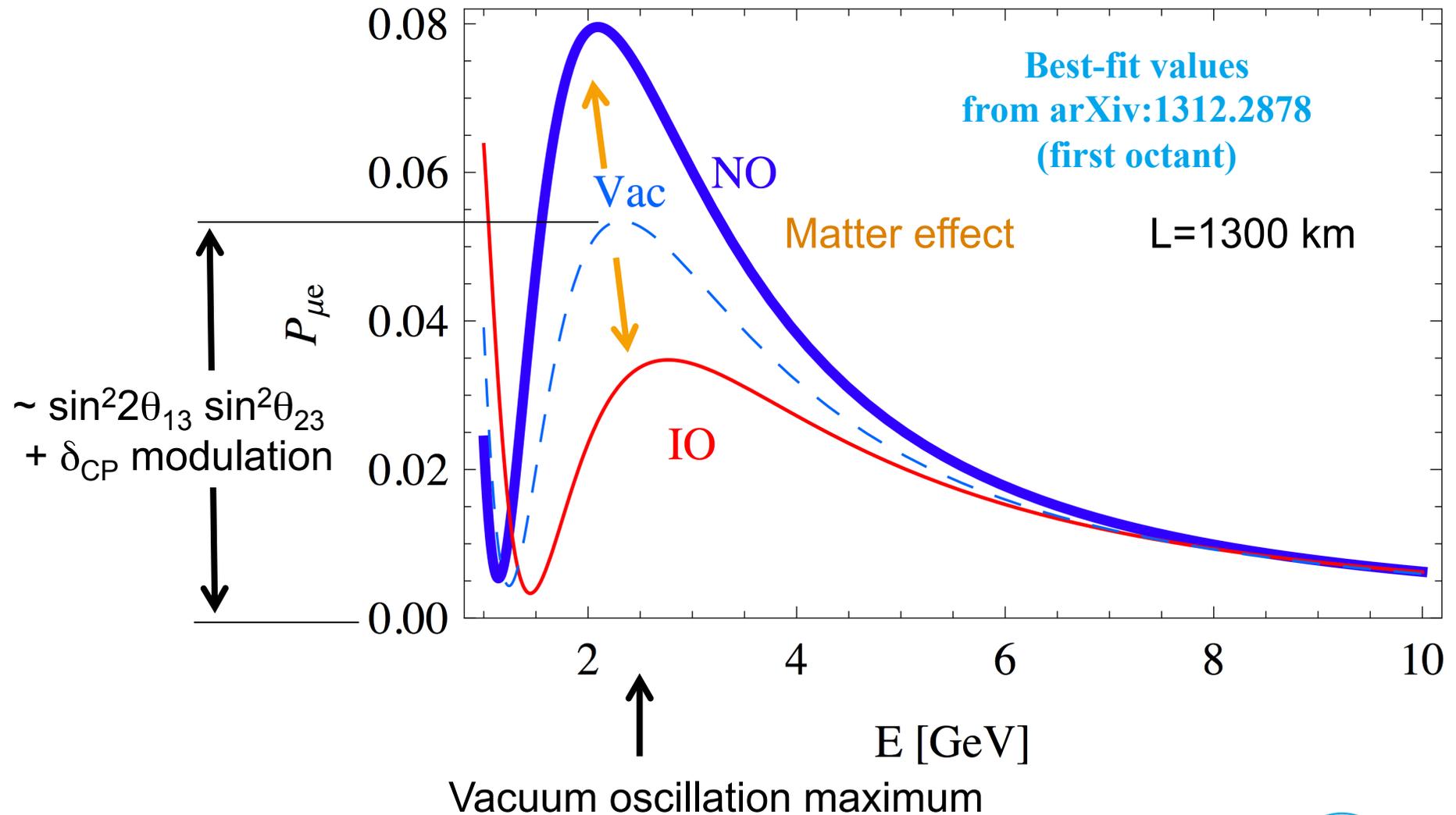
$$\cos 2\theta \rightarrow \hat{A} :$$



	Normal	Inverted
Neutrinos	Resonance	Suppression
Antineutrinos	Suppression	Resonance



Long baseline experiments (up to first vacuum osc. maximum)



Long-Baseline Neutrino Oscillation Experiment (LBNE)



Bob Wilson @
Neutrino 2014

> Particle Physics Project Prioritization Panel (P5) in the US; Report May '14

• The Science Drivers:

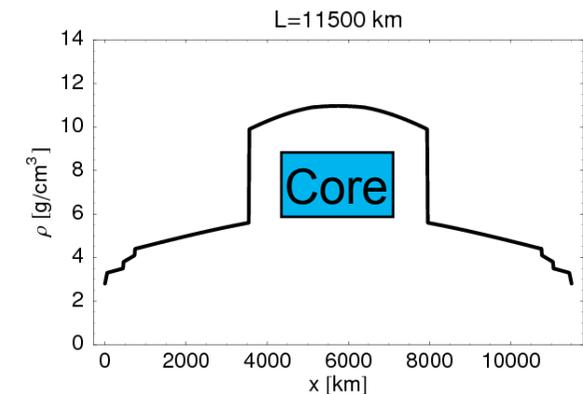
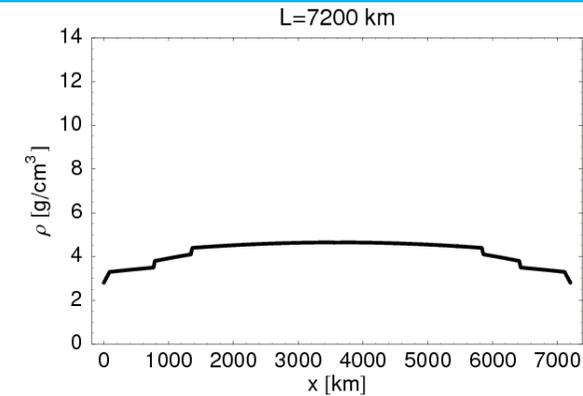
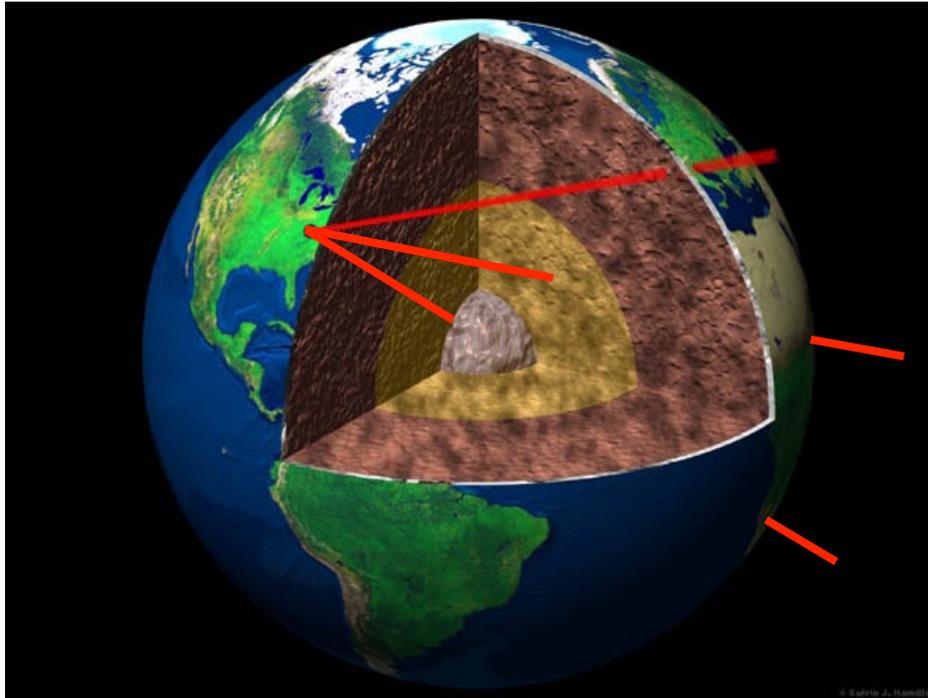
- Use the Higgs boson as a new tool for discovery
- **Pursue the physics associated with neutrino mass**
- Identify the new physics of dark matter
- Understand cosmic acceleration: dark energy and inflation
- Explore the unknown: new particles, interactions, and physical principles



Recommendation 13: Form a new international collaboration to design and execute a highly capable Long-Baseline Neutrino Facility (LBNF) hosted by the U.S. To proceed, a project plan and identified resources must exist to meet the minimum requirements in the text. LBNF is the highest-priority large project in its timeframe.



Matter profile of the Earth ... as seen by a neutrino



(Preliminary Reference Earth Model)

Resonance energy (from $\hat{A} \rightarrow \cos 2\theta$):

$$E_{\text{res}} [\text{GeV}] \sim 13\,200 \cos 2\theta \frac{\Delta m^2 [\text{eV}^2]}{\rho [\text{g}/\text{cm}^3]}$$

For ν_μ appearance, Δm_{31}^2 :

- $\rho \sim 4.7 \text{ g}/\text{cm}^3$ (Earth's mantle): $E_{\text{res}} \sim 6.4 \text{ GeV}$

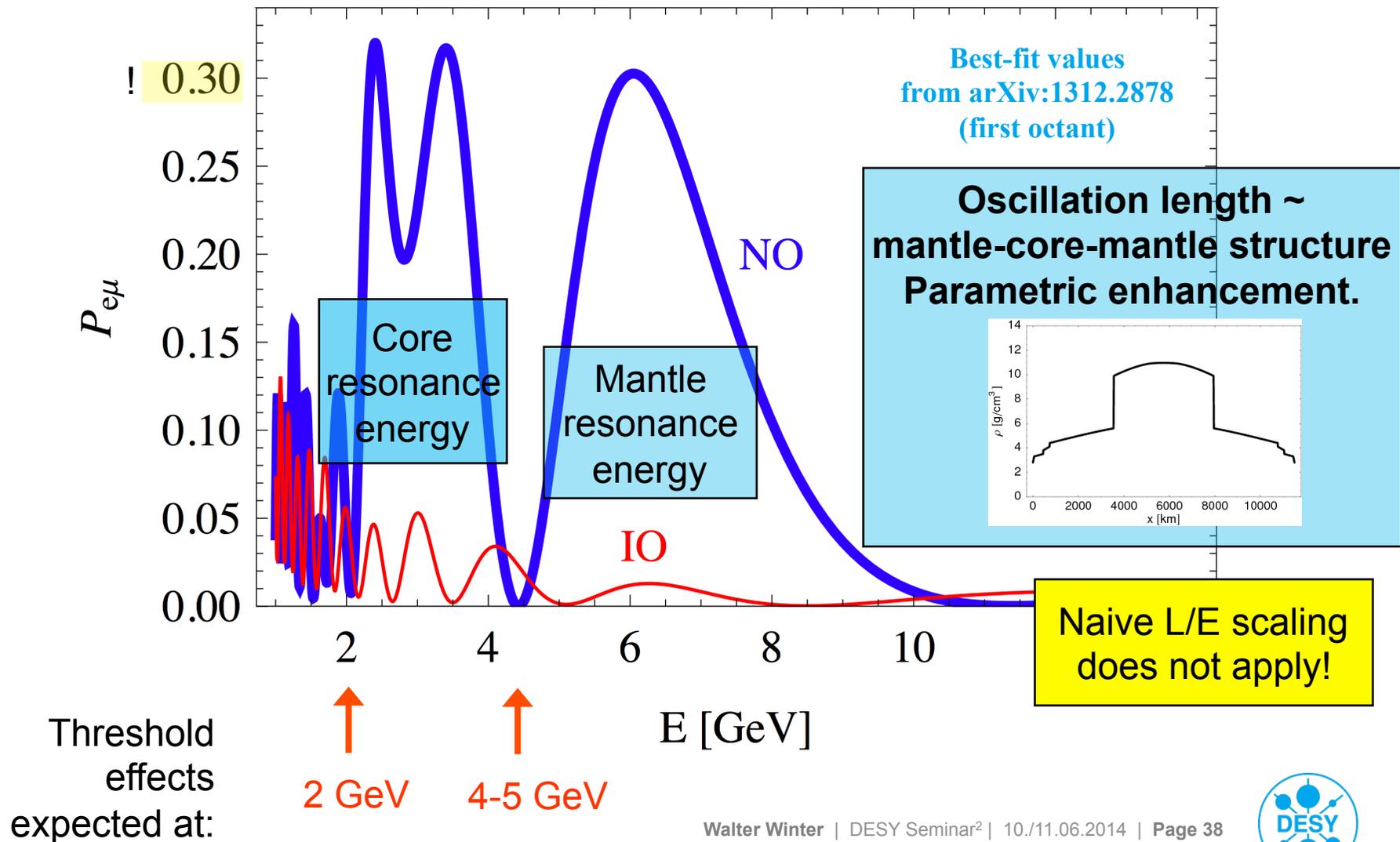
- $\rho \sim 10.8 \text{ g}/\text{cm}^3$ (Earth's outer core): $E_{\text{res}} \sim 2.8 \text{ GeV}$



Mantle-core-mantle profile

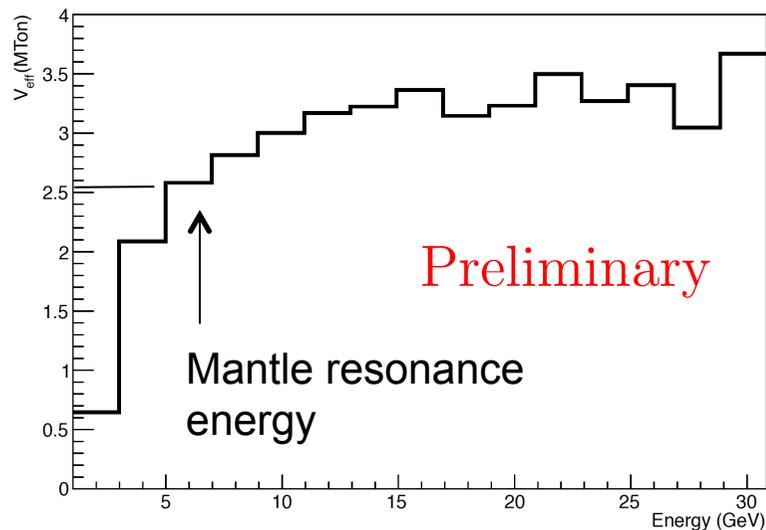
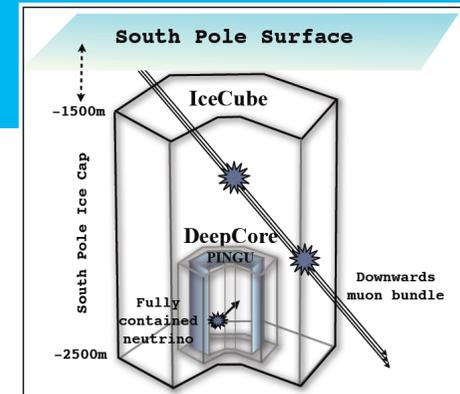
(Parametric enhancement: Akhmedov, 1998; Akhmedov, Lipari, Smirnov, 1998; Petcov, 1998)

> Probability for $L=11810$ km

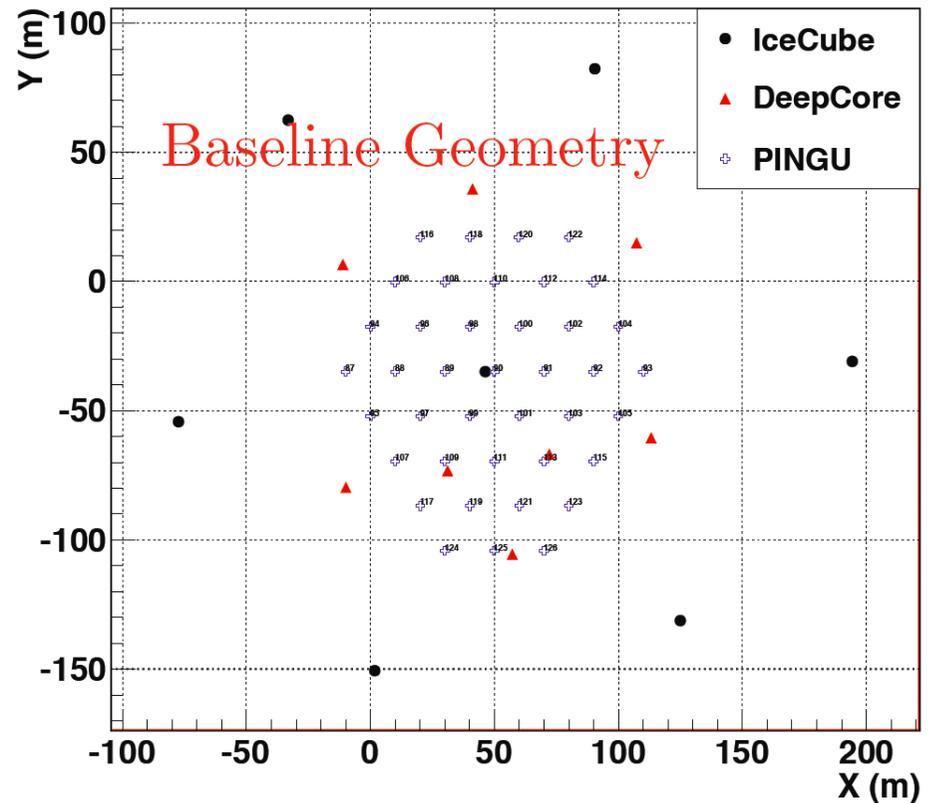


Emerging technologies: Atmospheric ν s

- Example: PINGU
("Precision IceCube Next Generation Upgrade")
- 40 additional strings, 60 optical modules each
- Lower threshold, few Mtons at a few GeV
- ORCA, INO: similar methods



(a) $V_{\text{eff}}(\nu_\mu)$



(PINGU LOI, arXiv:1401.2046)



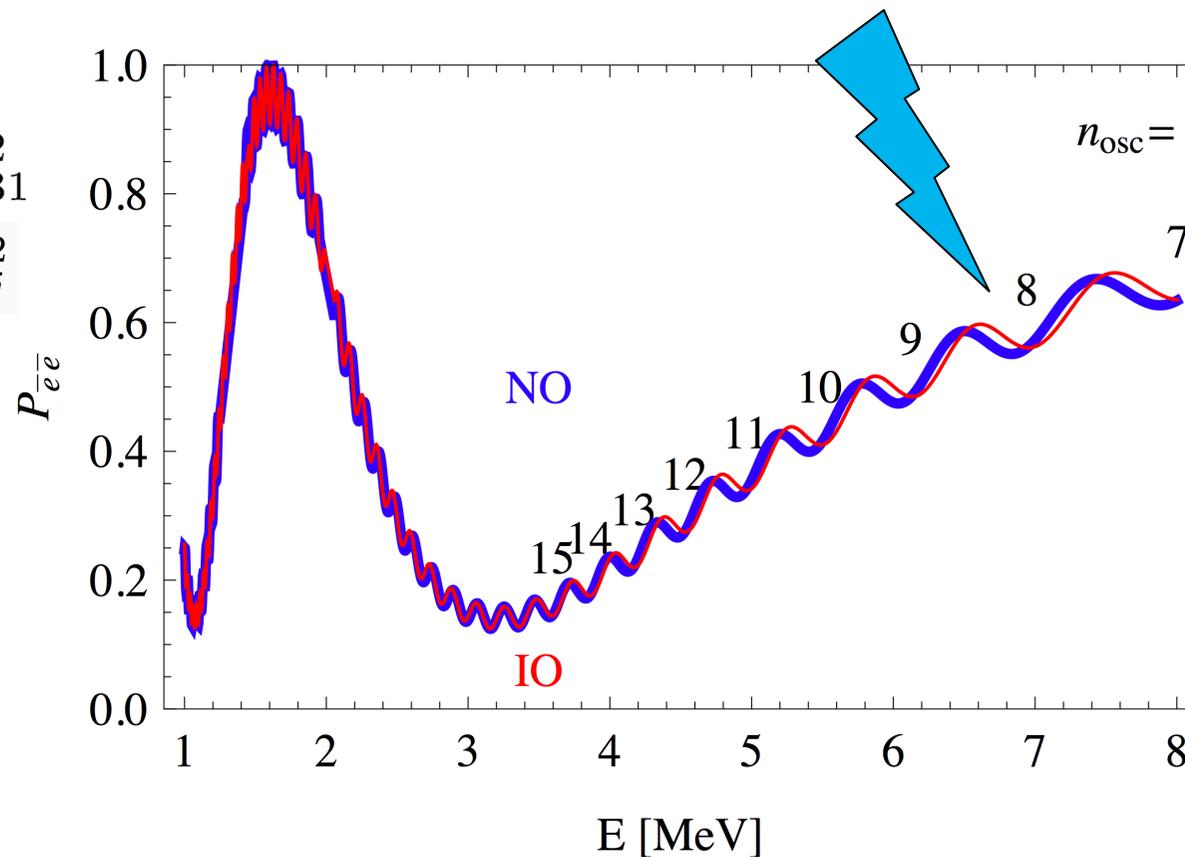
Method 2: Disappearance probabilities

- > Works in vacuum, and even for $\theta_{13}=0$
- > Just flipping the sign of Δm^2 is not sufficient
- > Example: Reactor experiment, $L=53$ km

$$\Delta m_{31}^2 \rightarrow -\Delta m_{31}^2$$

$$\Delta m_{ij}^2 = m_i^2 - m_j^2$$

Probabilities
apparently
different
(unphysical
effect!)



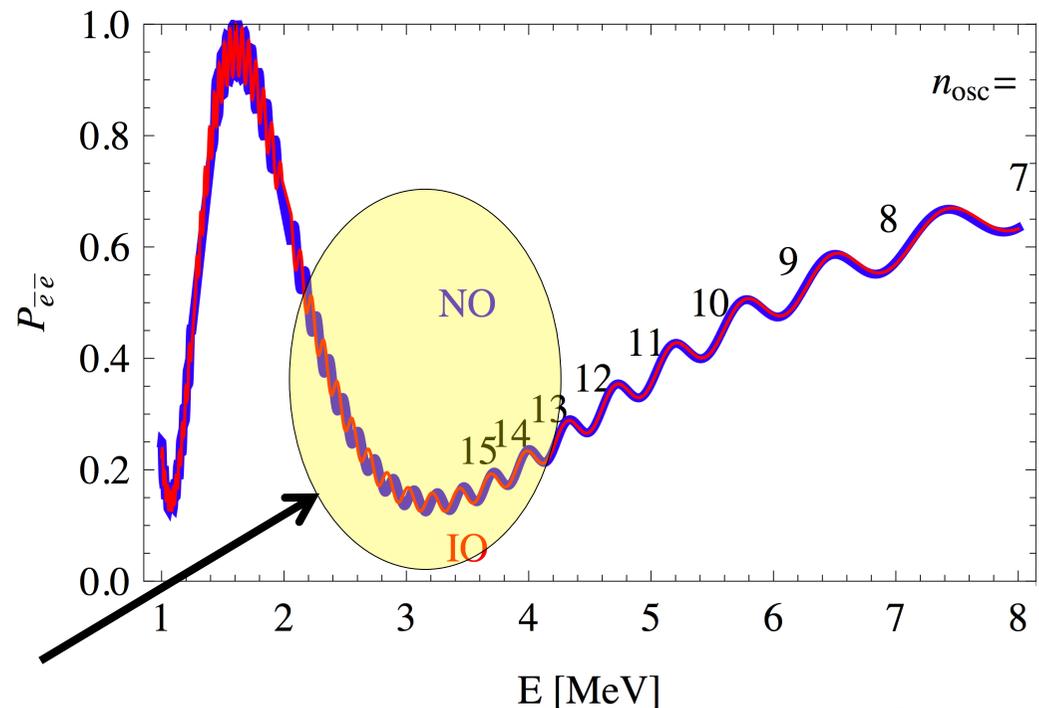
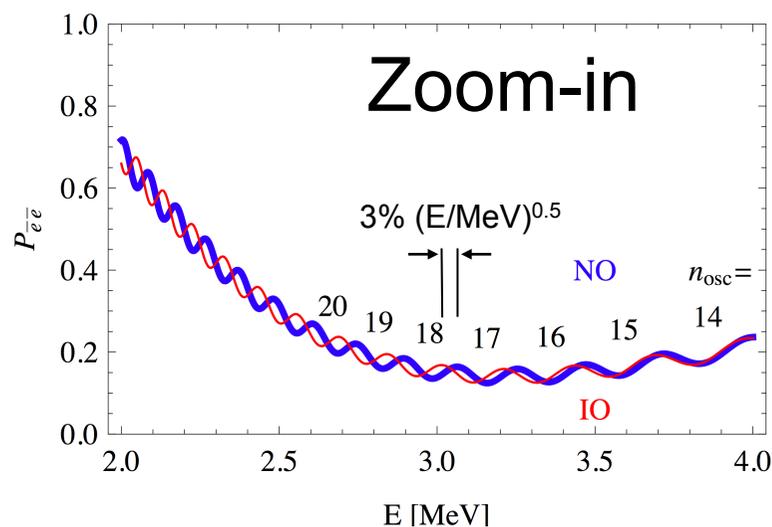
Method 2: Disappearance probabilities

- The disappearance Δm^2 depends on the channel. Consequence e. g.

$$|\delta m_{\text{eff}}^2|_e - |\delta m_{\text{eff}}^2|_\mu = \pm \delta m_{21}^2 (\cos 2\theta_{12} - \cos \delta \sin \theta_{13} \sin 2\theta_{12} \tan \theta_{23})$$

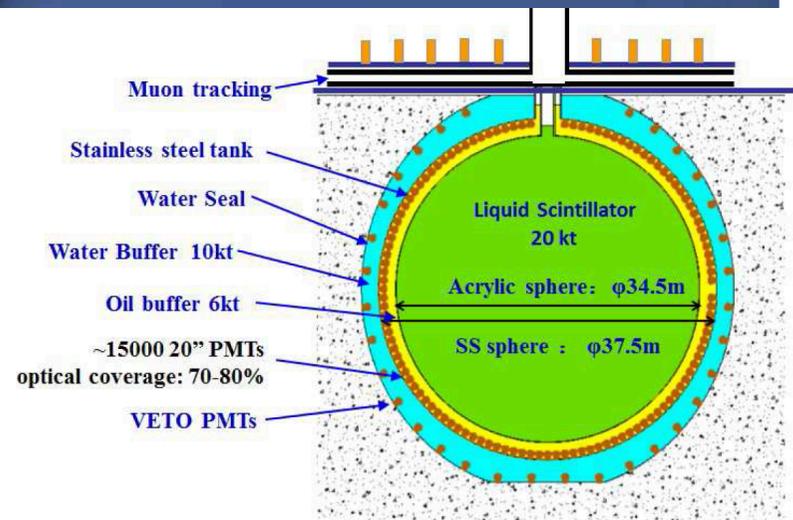
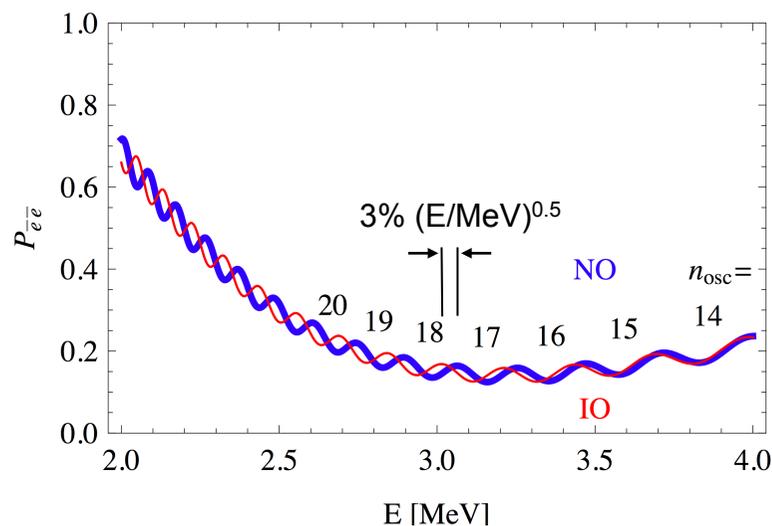
de Gouvea, Jenkins, Kayser, [hep-ph/0503079](#);
 Nunokawa, Parke, Zukanovich, [hep-ph/0503283](#)

- Now first oscillation maxima match. Discrimination by higher osc. Maxima. Need energy resolution!



Emerging technologies: Reactor experiments

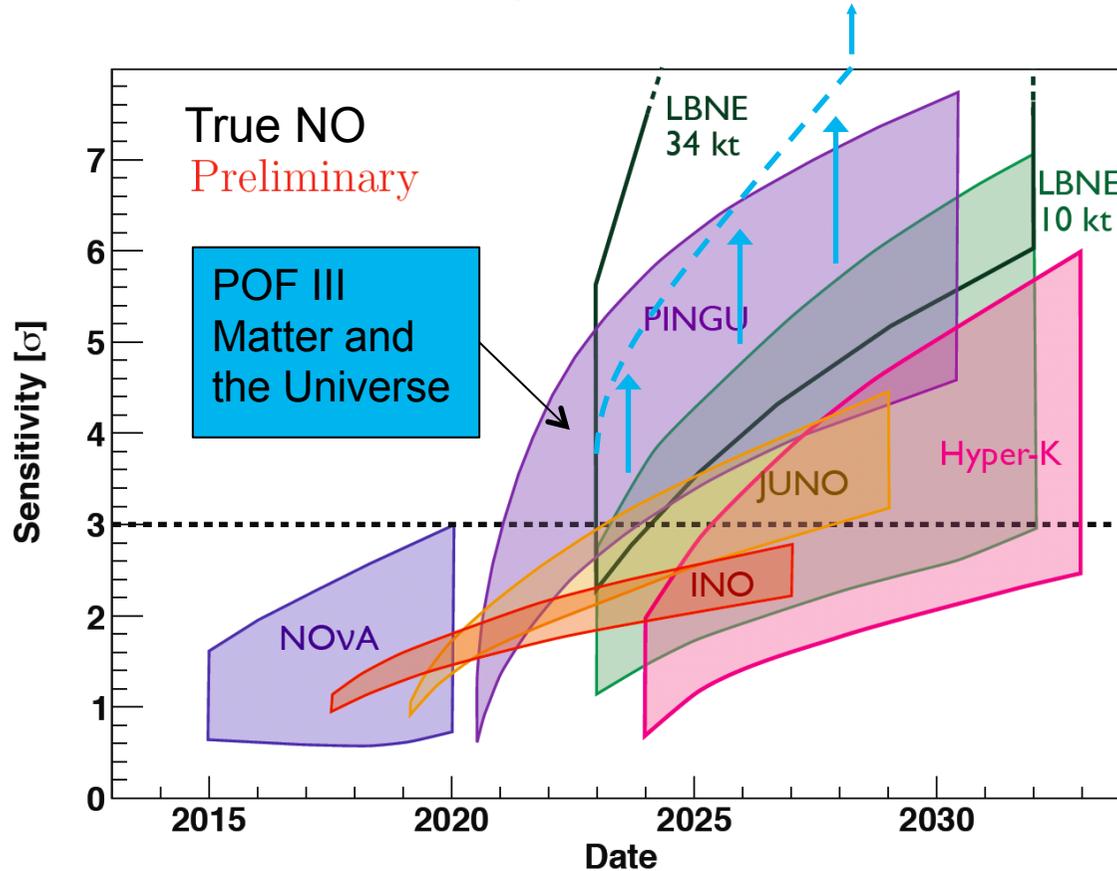
- Jiangmen Underground Neutrino Observatory (JUNO) [formerly Daya Bay-II]
- $L=53$ km
- Excellent energy resolution ($3\% (E/\text{MeV})^{0.5}$) requires $O(100\%)$ PMT coverage



Global context

LBNE 10kt if θ_{23} varied as well

⇒ Fig. 9 in arXiv:1305.5539; see also arXiv:1311.1822v2



- > Bands: risk wrt θ_{23} (PINGU, INO), δ_{CP} (NOvA, LBNE), energy resolution (JUNO)
- > LBNE and sensitivity also scales with θ_{23} !

(version from PINGU LOI, arXiv:1401.2046, based on Blennow, Coloma, Huber, Schwetz, arXiv:1311.1822)



Summary and conclusions

- > Mass hierarchy: may be tested in beginning of 2020s by “emerging technologies“, such as PINGU or JUNO

PINGU has a good chance to be the first experiment to measure the mass hierarchy if timely; DESY involved

- > CP violation: requires a new long-baseline experiment, such as LBNE, T2HK, NuFact;
P5 recognition milestone towards such a program at Fermilab
- > Other issues: θ_{23} maximal? Octant?
Sun and Earth tomography? New physics?
- > Light sterile neutrinos - best candidate for physics B ν SM? Test short-baseline anomalies, measure neutrino X-secs, ...
- > Perspectives for neutrino oscillations? Fantastic!

