# Inflationary Ramifications of a Meta-Stable Higgs Vacuum

arXiv: 1210.6987 Oleg Lebedev & AW DESY Hamburg the Higgs scalar potential ... if the coupling runs negative!





### geometry of tunneling - bubble nucleation II



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### geometry of tunneling - bubble nucleation II



## geometry of tunneling - bubble nucleation III



## how do we avoid that fate ...

we either hope to be in one of the few Universes, where the Higgs does not jump over the barrier ... [Espinosa, Giudice & Riotto '08]

#### or

Peccei-Quinn axion for theta-angle ...) we have a tiny coupling to a heavy singlet scalar: threshold effect stabilizes the Higgs potential [Lebedev; Elias-Miro, Espinosa, Giudice, Lee & Strumia '12]

(cf. see-saw for neutrino masses;

#### or

[Lebedev & Lee '12] we use inflation in our past .....[Lebedev & AW '12]



### Cosmic Microwave Background: PLANCK cosmology results 2013!







# what we know so far ...

- The initial fluctuations are plasma sound waves seeded by perturbations in energy density with very peculiar properties:
  - near uniform distribution of power ('scale invariance')
  - coherence (have the same initial phase)
  - correlations which have been out of causal contact at 370.000 yrs, if the Universe expanded always driven by matter or radiation

mounting evidence for a very early phase of hyperfast exponential expansion. of the Universe - Inflation.







equations of motion:

$$\ddot{\phi} + 3H\dot{\phi} = -V' =: \frac{dV}{d\phi}$$

scalar field equation

slow-roll:

scale factor: size of the Universe  $\frac{1}{2}\dot{\phi}^2 + V = H^2 = \frac{\dot{a}^2}{a^2}$ 

Friedmann (Einstein) equation

$$\begin{split} |\ddot{\phi}| \ll 3H|\dot{\phi}| & \text{and} \quad \dot{\phi}^2 \ll V \quad \Rightarrow \quad H \simeq const. \\ \Rightarrow \quad a = e^{Ht} \text{ for many e-folds } N = H \,\Delta t \quad , \quad \text{need: } N > 60 \\ & \text{exponential expansion:} \\ & inflation. \\ \text{slow-roll conditions on the scalar potential:} \\ \Rightarrow \quad \epsilon \equiv -\frac{\dot{H}}{H^2} \simeq \frac{1}{2} \left(\frac{V'}{V}\right)^2 \ll 1 \quad , \quad \eta \equiv \frac{\dot{\epsilon}}{\epsilon H} \simeq \frac{V''}{V} \ll 1 \end{split}$$





## instability insurance: couple the Higgs to the inflaton ...

[Lebedev & AW;12]

 $\mathbf{\Omega}$ 

<u>a simple example:</u>

$$\mathcal{L} = \frac{1}{2} (\partial \mu h)^2 + \frac{1}{2} (\partial \mu \phi)^2 - V \quad \text{with: } V = V(h) + \frac{1}{2} m^2 \phi^2 + \frac{1}{2} \xi \phi^2 h^2$$

radiative stability  
of inflation:
$$\xi \lesssim 10^{-6}$$
Higgs is stable  
for: $\phi_0 \gtrsim 20 M_P$ 

<u>can start inflaton & Higgs at:</u>  $\phi_0 \gtrsim 20 M_P$  ,  $h_0 \simeq 0.1 M_P$ 

and Higgs will run exponentially fast to  $b = v = 10^{-16} M_P$  in about 20 e-folds !

see also SFB lectures by: Oleg Lebedev

## evolution of the Higgs during inflation ... [Lebedev & AVV ; 12]



## Thanks!