# Vacuum stability in the Standard Model at two loops

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Introduction	Higgs potential	Vacuum stability	Numerical analysis	Outlook
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#### **Outline**



- Higgs potential
- Vacuum stability
- Mumerical analysis





#### Introduction: experimental status quo



Higgs boson of mass  $m_H = (125.6 \pm 0.3)$  GeV w/ SM properties





- $m_H = 125.6 \text{ GeV}$  agrees w/ EW precision data.
- Triviality bound satisfied.
- How about vacuum stability bound?



- $\frac{g_V}{2} \equiv m_W \quad \rightsquigarrow \quad v = 2^{-1/4} G_F^{-1/2} = 246.220 \text{ GeV}$
- $m_H$  is free parameter.

So far, bare fields and parameters.

# Renormalizaton: RG evolution

Cosmological applications require reliable predictions over very large range of scales:  $v \leq \mu \leq M_P$ Use  $\overline{\text{MS}}$  renromalization scheme: running couplings  $\lambda(\mu^2), y_t(\mu^2), g_s(\mu^2), \dots$ Two-step procedure: 1. RG evolution:

$$\mu^{2} \frac{d\lambda(\mu^{2})}{d\mu^{2}} = \beta_{\lambda} = \frac{1}{16\pi^{2}} (12\lambda^{2} + 6\lambda y_{t}^{2} - 3y_{t}^{4}) + \cdots$$
$$\mu^{2} \frac{dy_{t}(\mu^{2})}{d\mu^{2}} = \beta_{y_{t}} = \frac{1}{16\pi^{2}} y_{t} \left(\frac{9}{4}y_{t}^{2} - 4g_{s}^{2}\right) + \cdots$$
$$\mu^{2} \frac{dg_{s}(\mu^{2})}{d\mu^{2}} = \beta_{g_{s}} = \frac{1}{16\pi^{2}} g_{s}^{3} \left(-\frac{11}{2} + \frac{n_{f}}{3}\right) + \cdots$$

 $eta_{\lambda}^{(3)},eta_{y_t}^{(3)}$ 

 $\beta_{\cdots}^{(3)}, \beta_{g_s, y_t}^{(3)}$ 

Chetyrkin, Zoller, JHEP06(2012)033; 04(2013)091 Bednyakov *et al.*, PLB722(2013)336; arXiv:1303.4364 Mihaila *et al.*, PRL108(2012)151602; PRD86(2012)096008 Tarasov *et al.*, PLB93(1980)429

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#### **Threshold corrections**

2. Matching at  $\mu_0 = \mathscr{O}(v)$ :

 $\delta_{H}^{(lpha lpha_{s})}, \delta_{t}^{(lpha lpha_{s})} \ \delta_{H}^{(y_{t}^{4})}, \delta_{t}^{(y_{t}^{4})}$ 

Bezrukov, Kalmykov, BK, Shaposhnikov, JHEP10(2012)140 Degrassi *et al.*, JHEP08(2012)098

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#### Triviality and vaccum stability in a nutshell

Recall

$$\mu^{2} \frac{d\lambda(\mu^{2})}{d\mu^{2}} = \frac{1}{16\pi^{2}} (12\lambda^{2} + 6\lambda y_{t}^{2} - 3y_{t}^{4}) + \cdots$$
$$\mu^{2} \frac{dy_{t}(\mu^{2})}{d\mu^{2}} = \frac{1}{16\pi^{2}} y_{t} \left(\frac{9}{4} y_{t}^{2} - 4g_{s}^{2}\right) + \cdots$$
$$\lambda(m_{H}^{2}) = 0.130 \times \left(\frac{m_{H}}{125.6 \text{ GeV}}\right)^{2}, \quad y_{t}(m_{t}^{2}) = 0.993 \times \frac{m_{t}}{172.9 \text{ GeV}}, \quad g_{s}(m_{Z}^{2}) = 1.220$$

- Triviality bound: Maiani *et al.*, NPB136(1979)115 If  $m_H > M_{max}$ , then  $\lambda(\mu^2) \to \infty$  for  $\mu \to \mu_{Landau}$ . (For  $m_t \ll m_H$ ,  $\mu_{Landau} \approx m_H \exp \frac{2\pi^2}{3\lambda(m_H^2)} = 1.2 \times 10^{24} \text{ GeV}$ )  $\to \lambda(\mu_0^2) = 0 \text{ trivial}$
- Vacuum stability bound: Lindner, ZPC31(1986)295 If  $m_H < M_{min}$ , then  $\lambda(\mu^2) < 0$  for  $\mu > \mu_{stab}$ .  $\rightarrow$  Decay of universe

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#### Vaccum stability condition



Determine  $M_{\min}$  so that for  $m_H = M_{\min}$ 

$$\lambda(\mu_{\text{stab}}) = 0 = \beta_{\lambda}(\lambda(\mu_{\text{stab}}))$$

at some given  $\mu_{stab} \gg v$ , *e.g.*  $\mu_{stab} = M_P$ . Caveat:  $M_{min}$  is (slightly) scheme dependent.  $\rightsquigarrow$  theoretical uncertainty

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#### Effective potential



Determine  $\widetilde{M}_{\min}$  so that for  $m_H = \widetilde{M}_{\min}$ 

 $V(\Phi_{\rm SM}) = V(\Phi_1), \qquad V'(\Phi_{\rm SM}) = V'(\Phi_1)$ 

at some given  $\Phi_1 \gg \Phi_{SM}$ , *e.g.*  $\Phi_1 = \Phi_P$ . NB: Numerically,  $\widetilde{M}_{\min} - M_{\min} = \mathcal{O}(0.1 \text{ GeV})$ , *i.e.* well within theoretical uncertainty.

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Numerical analysis Bezrukov, Kalmykov, BK, Shaposhnikov, JHEP10(2012)140

3-loop evolution / 2-loop matching yields:

$$M_{\rm min} = \left[ 128.95 + \frac{M_t - 172.9 \,\,{\rm GeV}}{1.1 \,\,{\rm GeV}} \times 2.2 - \frac{\alpha_s - 0.1184}{0.0007} \times 0.56 \right] \,\,{\rm GeV}$$

Source of uncertainty	Nature of estimate	$\Delta_{\text{theor}} M_{\text{min}} \text{ [GeV]}$
3-loop matching $\lambda$	sensitivity to $\mu_0$	1.0
3-loop matching y <sub>t</sub>	sensitivity to $\mu_0$	0.2
4-loop $\alpha_s$ to $y_t$	educated guess [Kataev, Kim]	0.4
confinement, y <sub>t</sub>	educated guess $\sim \Lambda_{QCD}$	0.5
4-loop running $M_W \rightarrow M_P$	educated guess	< 0.2
total uncertainty	sum of squares	1.2
total uncertainty	linear sum	2.3

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#### Anatomy



Matching scale  $\mu_0$ , GeV

Contribution	$\Delta M_{\min}$ [GeV]
3-loop beta functions	-0.23
$\delta y_t \propto O(\alpha_s^3)$	-1.15
$\delta y_t \propto O(\alpha \alpha_s)$	-0.13
$\delta\lambda \propto O(\alpha\alpha_{\rm s})$	0.62
$\delta y_t, \delta \lambda \propto O(y_t^4)$	0.2

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## $\mathcal{O}(\alpha \alpha_s)$ threshold corrections



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#### Reduction of fundamental scales



- $\mu_{\text{stab}} = 2.9 \times 10^{18} \text{ GeV}$  stable w.r.t. variations of  $m_t = (172.9 \pm 1.1) \text{ GeV}, \ \alpha_s^{(5)}(m_Z^2) = 0.1184 \pm 0.0007 \text{ (dashed)},$ and  $m_Z < \mu_0 < m_t$  (yellow).
- $\mu_{\rm stab} \approx M_P = 2.44 \times 10^{18} {\rm GeV}$
- Electroweak scale is determined by Planck scale physics!



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## LC as top and Higgs factory



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#### Vacuum metastability



Degrassi et al., JHEP08(2012)098

# EW vacuum is metastable / unstable, if its lifetime overshoots / undershoots that of the universe.

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# Outlook: pole mass m<sub>t</sub>

- PDG value  $M_X(t \rightarrow X) = (172.9 \pm 1.1)$  GeV is not pole mass  $m_t$ , but just parameter in MC programs w/o RC to partonic cross sections.
- Rigorous determination of  $\overline{\text{MS}}$  mass  $\overline{m}_t(\mu^2)$  from  $\sigma_{\text{tot}}(p\bar{p}, pp \rightarrow t\bar{t} + X) \rightsquigarrow \text{Moch's talk}$
- $\overline{m}_t(\mu^2) m_t$  receives large EW RC from tadpole contributions.



Jegerlehner, Kalmykov, BK, PLB722(2013)123

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BSM physics				

- Depending on future precision measurements of  $m_H, m_t, \alpha_s$  and higher-loop RC calculations, SM may be stable way up to  $M_P$ .
- Reduction of  $m_t$  by 1.6 GeV [cf.  $m_t m_{\tilde{t}} = (-1.4 \pm 2.0)$  GeV,  $\Gamma_t = (2.0^{+0.7}_{-0.6})$  GeV]  $\rightsquigarrow M_{\min} = m_H = 125.6$  GeV
- BSM physics still necessary to solve open problems, *e.g.* smallness of neutrino masses, strong CP problem, dark matter, baryon asymmetry of universe, *etc.* → Westphal's talk